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A statistical physics approach to compressed sensing or y=Ax revisited

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Compressed sensing or y=Ax revisited

- What is compressed sensing?
- What is the link between statistical physics and compressed sensing?
- How can one use statistical physics to improve on compressed sensing technics?

Compressed sensing or y=Ax revisited

• What is compressed sensing?

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From 10⁶ wavelet coefficients, keep 25.000

Most signal of interest are <u>sparse</u> in an <u>appropriated basis</u> \Rightarrow Exploited for data compression (JPEG2000).



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Why do we record a huge amount of data, and then keep only the important bits?

Couldn't we record <u>only</u> the relevant information directly?



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Couldn't we record <u>only</u> the relevant information directly?

Compressed Sensing

Record <u>directly</u> in compressed form (gain of time and storage)
Reconstruct the original signal afterwards



Why do we record a huge amount of data, and then keep

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Compressed sensing

DL Donoho - Information Theory, IEEE Transactions on, 2006 - ieeexplore.ieee.org Abstract Suppose x is an unknown vector in Ropf m (a digital image or signal); we plan to measure n general linear functionals of x and then reconstruct. If x is known to be compressible by transform coding with a known transform, and we reconstruct via the ... <u>Cité 4384 fois</u> - <u>Autres articles</u> - <u>findit@espci</u> - <u>Les 46 versions</u>



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An example from magnetic resonance imaging





Left: image acquired with CS Acceleration by a factor 2.5

Lustig, Donoho, Pauly '07

Possible applications

- Rapid Magnetic Resonance Imaging
- Image acquisition (single-pixel camera)
- DNA microarrays
- Group testing
- Fast data compression
- Herschel spacial telescope
- Compressed Sensing Microscopes
- Sparse Principal Component Analysis
- Compressed quantum state tomography



M measurements

M linear operations on the vector





vector of size M



M×N matrix



M×N matrix



Problem: you know y and G, how to reconstruct I?





If M<N representations of equations



M×N matrix



<u>Compressed sensing input:</u> The signal is sparse in an appropriate basis









 $M \times N$ matrix

The problem to

solve is now

- $\vec{y} = F\vec{x}$
- with $F=G\psi$

F=M×N matrix

- Need to find a <u>sparse solution</u> of an <u>under-constrained</u> set of linear equations
- Ideally works as long as M>R
- Robust to noise

The reconstruction problem: Inverting an underconstrained linear system

Consider a system of linear measurements



 $F = M \times N$ matrix

Generically: • if M = N

Unique solution obtained by inversion $x = F^{-1}y$

Generically: • if M = NUnique solution obtained by inversion $x = F^{-1}y$ • if M > N solution obtained from the inversion of a $N \times N$ submatrix of F with full rank



NB: too many equations, redundant system, **<u>but</u>** consistent because the y measurements are obtained as y = Fx

Generically: • if M < N



Not enough measurements to determine the signal x from its linear transform y

Generically: • if M < N



Not enough measurements to determine the signal x from its linear transform y

To invert, you need as many measurements (M) as number of unknowns (N)

• if M < N but x is <u>sparse</u> (only R of its components are $\neq 0$)



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CLAIM:To invert, you need as many measurements (M) as number of unknown (R)

• if M < N but x is <u>sparse</u> (only R of its components are $\neq 0$) $\begin{bmatrix} y \end{bmatrix} = \begin{bmatrix} F \end{bmatrix}$

CLAIM:To invert, you need as many measurements (M) as number of unknown (R)

If R < M < N : the reconstruction of the signal x from the measurement y is possible

• if M < N but x is <u>sparse</u> (only R of its components are $\neq 0$) $\begin{bmatrix} y \\ \end{bmatrix} = \begin{bmatrix} F \\ F \end{bmatrix} \begin{bmatrix} x \\ x \\ \end{bmatrix} \begin{bmatrix} x \\ R \text{ non zero} \\ N-R \text{ zero} \end{bmatrix}$

A 'simple' solution: guess the positions where $x_i \neq 0$

• if M < N but x is <u>sparse</u> (only R of its components are $\neq 0$) $\begin{bmatrix} y \\ y \end{bmatrix} = \begin{bmatrix} G \\ F \end{bmatrix} \begin{bmatrix} x \\ x_1, \dots, x_R \neq 0 \end{bmatrix}$

A 'simple' solution: guess the positions where $x_i \neq 0$

Solve:
$$y^{\mu} = \sum_{i=1}^{R} G^{\mu i} x_i \qquad \mu = 1, \dots, M$$

 $R < M \implies$ too many equations

generically inconsistent (no solution), except if the guess of locations of $x_i \neq 0$ was correct
The problem: y = Fx , find x

• if M < N but x is <u>sparse</u> (only R of its components are $\neq 0$) $\begin{bmatrix} y \\ y \end{bmatrix} = \begin{bmatrix} G \\ F \end{bmatrix} \begin{bmatrix} x \\ x_1, \dots, x_R \neq 0 \end{bmatrix}$ $\binom{N}{R}$ possible guesses Long, but finite time... R < Iexcept if the guess of locations of $x_i \neq 0$ was correct

One can reconstruct a N-dimensional sparse signal with R non-zero components from N>M>R measurements



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Less measurements (gain of time and precision)

One can reconstruct a N-dimensional sparse signal with R non-zero components from N>M>R measurements



- Less measurements (gain of time and precision)
- Data already compressed (gain of memory storage)

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- Price to pay: a reconstruction algorithm is needed

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M×N matrix

- Less measurements (gain of time and precision)
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- Price to pay: a reconstruction algorithm is needed

The "simple" algorithm we have presented is too slow! (need to try exponentially many cases)

One can reconstruct a N-dimensional sparse signal with R non-zero components from N>M>R measurements



M×N matrix

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The goal of CS theory:

Determine a <u>sensing matrix F</u> and a <u>reconstruction algorithm</u> such that the reconstruction is possible <u>in practice</u>

A phase diagram





- Incoherent samplings (i.e. a random matrix F)
- Reconstruction by minimizing the L_I norm $||\vec{x}||_{L_1} = \sum_i |x_i|$

Candès & Tao (2005) Donoho and Tanner (2005)

State of the art in CS



For a signal with $(1-\rho)N$ zeros R= ρN non zeros

and a Gaussian random matrix with $M = \alpha N$

State of the art in CS



For a signal with $(1-\rho)N$ zeros R= ρN non zeros

and a Gaussian random matrix with $M = \alpha N$

Reconstruction limited by the Donoho-Tanner transition for the $L_{\rm I}$ norm minimization

One measurement (scaling product with a random pattern)



· Each measurement touches every part of the underlying signal/image

Many measurements (scaling product with many random patterns)





signal

• Take K = 96000 incoherent measurements y = GI

From 10⁶ points, but only, 25.000 non zero

Solve

min $\|\mathbf{x}\|_{\ell_1}$ subject to $\mathbf{G}\Psi\mathbf{x} = y$

 Ψ = wavelet transform



original (25k wavelets)



perfect recovery

Compressed Sensing: A (short!)review of the present litterature:

- Record data already in a compressed form
- Less measurements (faster, more precise)...
- ... but need for a reconstruction algorithm!
- State of the art: L_1 -minimization and random measurements

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 $P(\vec{x}|\vec{y}) = \frac{1}{Z} \prod_{i=1}^{N} P(x_i) \prod_{\mu=1}^{M} \delta\left(y_{\mu} - \sum_{i=1}^{N} F_{\mu i} x_i\right) \text{ with } P(x_i) = (1-\rho)\delta(x_i) + \rho\phi(x_i)$



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A statistical-physics approach to compressed sensing



Estimating the probability of each value of x is equivalent to solving a mean-field disordered statistical physics problem

$$P(\vec{x}|\vec{y}) = \frac{1}{Z} e^{-\sum_{i=1}^{N} \log P(x_i) - \frac{1}{2\Delta} \sum_{\mu=1}^{M} (y_{\mu} - \sum_{i=1}^{N} F_{\mu i} x_i)^2}$$



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<u>Theorem</u>: sampling from P(x|y) gives the correct solution as long as $\alpha > \rho_0$ if: a) $\Phi(x) > 0 \forall x$ and b) $1 > \rho > 0$

The probabilistic approach is optimal, even if we do not know the correct $\Phi(x)$! In practice, we use a <u>Gaussian distribution</u>

A sketch of the proof

Consider the system constrained to be at distances larger than D with respect to the solution

$$Y(D,\epsilon) = \int \prod_{i=1}^{N} \left(dx_i \left[(1-\rho)\delta(x_i) + \rho\phi(x_i) \right] \right) \prod_{\mu=1}^{M} \delta_\epsilon \left(\sum_i F_{\mu i}(x_i - s_i) \right) \mathbb{I}\left(\sum_{i=1}^{N} (x_i - s_i)^2 > ND \right)$$

I) Y(0) is infinite if $\alpha > \rho_0$ (equivalently if M>R) (just count the delta functions! N-R+M deltas versus N integrals...)

2) Y(D) is finite for any D>0 (bound by a first moment method, or "annealed" computation)
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If $\alpha > \rho_0$, the measure is always dominated by the solution

A sketch of the proof

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A statistical physics approach to compressed sensing

One can use statistical physics tools for

I) Computing phase transitions analytically (reconstruction/non reconstruction, etc...) Tools: Replica method from spin glass theory, etc...

II) Develop new algorithms, and design new matrices to improve on the L_I state-of-the art.

Tools: Replica and Cavity method from spin glass theory, Mean field methods from stat-phys, Physics intuition, etc.... A statistical physics approach to compressed sensing

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Statistical physics of compressed sensing

Model with N infinite-range 1d interacting particles with positions \mathbf{x}_i

What is the phase diagram of the system? $Z(y) = \int \prod_{i=1}^{N} dx_i P(x|y) \qquad \qquad F(\vec{y}) = -\log Z(\vec{y})$

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Averaging over disorder:

 $F_{\mu i} \quad \text{iid Gaussian, variance} \quad 1/N$ $y_{\mu} = \sum_{i=1}^{N} F_{\mu i} x_{i}^{0} \text{ where } x_{i}^{0} \text{ are iid distributed from } (1 - \rho_{0}) \delta(x_{i}^{0}) + \rho_{0} \phi_{0}(x_{i})$

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Analytic study: cavity equations, density evolution, replicas

$$E(Z^n) = \max_{Q, q, m, \hat{Q}, \hat{q}, \hat{m}} e^{Nn\Phi(Q, q, m, \hat{Q}, \hat{q}, \hat{m})}$$

$$\Phi(Q,q,m,\hat{Q},\hat{q},\hat{m}) = -\frac{1}{2N} \sum_{\mu} \frac{q-2m+\rho+\Delta_{\mu}}{\Delta_{\mu}+Q-q} - \frac{1}{2N} \sum_{\mu} \log\left(\Delta_{\mu}+Q-q\right) + \frac{Q\hat{Q}}{2} - m\hat{m} + \frac{q\hat{q}}{2} + \int \mathcal{D}z \int \mathrm{d}x_0 \left[(1-\rho_0)\delta(x_0) + \rho_0\phi_0(x_0)\right] \log\left\{\int \mathrm{d}x \, e^{-\frac{\hat{Q}+\hat{q}}{2}x^2 + \hat{m}xx_0 + z\sqrt{\hat{q}}x} \left[(1-\rho)\delta(x) + \rho\phi(x)\right]\right\}$$

Order parameters:

$$Q = \frac{1}{N} \sum_{i} \langle x_i^2 \rangle \qquad \qquad q = \frac{1}{N} \sum_{i} \langle x_i \rangle^2 \qquad \qquad m = \frac{1}{N} \sum_{i} x_i^0 \langle x_i \rangle$$

Mean square error:

$$E = \frac{1}{N} \sum_{i} \left(\langle x_i \rangle - x_i^0 \right)^2 = q - 2m + \langle (x_i^0)^2 \rangle_0$$

Example with $\rho_0=0.4$, and Φ_0 a Gaussian distribution with zero mean and unit variance



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• Maximum is at E=0 (as long as $\alpha > \rho 0$): Equilibrium behavior dominated by the original signal

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- Maximum is at E=0 (as long as $\alpha > \rho 0$): Equilibrium behavior dominated by the original signal
- For α < 0.58, a secondary maximum appears (meta-stable state): spinodal point

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- A steepest ascent dynamics starting from large E would reach the signal for α >0.58, but would stay block in the meta-stable state for α <0.58, even if the true equilibrium is at E=0.

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- A steepest ascent dynamics starting from large E would reach the signal for α >0.58, but would stay block in the meta-stable state for α <0.58, even if the true equilibrium is at E=0.
- Similarity with metastable phase in first-order transition (supercooled liquids)

Computing the Phase Diagram



Computing the Phase Diagram



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 II) Develop new algorithms, and design new matrices to improve on the L₁ state-of-the art.
 Tools: Bethe-Peirls method/Belief propagation, Mean field methods from stat-phys, Physics intuition, etc....

The Belief-Propagation algorithm (a sketchy description)

• NO averaging: work on a given problem

•Compute $f(\{\mathcal{P}_i(x_i)\}) = \log Z(\{\mathcal{P}_i(x_i)\})$ the potential with constrained local probabilities (marginals) for each variable.

•Derive the recursion equation for by steepest ascent/descent:

 $\mathcal{P}_i^{t+1} = \nabla_{\mathcal{P}_i^t} f(\{\mathcal{P}_i^t\})$

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 $\mathcal{P}_i^{t+1} = \nabla_{\mathcal{P}_i^t} f(\{\mathcal{P}_i^t\})$

- •This approach has been used :
 - Mean-field, Curie-Weiss, TAP (Thouless-Anderson-Palmer), or Cavity Method in Physics, and can be traced to Bethe-Peierls and Onsager ('35).
 - Belief Propagation in Artificial Intelligence (Pearl, '82)
 - Sum-Product in Error-Correcting-Codes (Gallager, '60)

Gibbs free energy approach: $\log Z = \max_{\{\mathcal{P}(\vec{x})\}} f_{Gibbs} \left(\{\mathcal{P}(\vec{x})\}\right)$ With $f_{Gibbs} \left(\{\mathcal{P}(\vec{x})\}\right) = -\langle \log P(\vec{x}|\vec{y}) \rangle_{\mathcal{P}(\vec{x})} - \int d\vec{x} \mathcal{P}(\vec{x}) \log \mathcal{P}(\vec{x})$

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Mean-Field $\Rightarrow \mathcal{P}(\vec{x}) = \prod_i \mathcal{P}_i(\vec{x}_i)$

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$$\Rightarrow \mathcal{P}(\vec{x}) = \prod_{i} \mathcal{P}_{i}(\vec{x}_{i})$$
 = Convergence problems

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Mean-Field
$$\Rightarrow \mathcal{P}(\vec{x}) = \prod_i \mathcal{P}_i(\vec{x}_i)$$
 = Convergence problems

Belief-Propagation
$$\Rightarrow \mathcal{P}(\vec{x}) = \frac{\prod_{ij} \mathcal{P}_{ij}(\vec{x}_i, \vec{x}_j)}{\prod_i \mathcal{P}_i(\vec{x}_i)^{M-1}}$$

Gibbs free energy approach: $\log Z = \max_{\{\mathcal{P}(\vec{x})\}} f_{Gibbs} \left(\{\mathcal{P}(\vec{x})\}\right)$ With $f_{Gibbs} \left(\{\mathcal{P}(\vec{x})\}\right) = -\langle \log P(\vec{x}|\vec{y}) \rangle_{\mathcal{P}(\vec{x})} - \int d\vec{x} \mathcal{P}(\vec{x}) \log \mathcal{P}(\vec{x})$

Mean-Field
$$\Rightarrow \mathcal{P}(\vec{x}) = \prod_{i} \mathcal{P}_{i}(\vec{x}_{i})$$
 = $\mathcal{N}ot \text{ correct}$
+Convergence problems

Belief-Propagation
$$\Rightarrow \mathcal{P}(\vec{x}) = \frac{\prod_{ij} \mathcal{P}_{ij}(\vec{x}_i, \vec{x}_j)}{\prod_i \mathcal{P}_i(\vec{x}_i)^{M-1}}$$

(asymptotically) exact in CS with random matrices

<u>Simplification</u> thanks to the dense matrix limit: Projection on first two moments is enough :

$$f\left(\{\mathcal{P}_i(x_i), \mathcal{P}_{ij}(x_i, x_j)\}\right)$$

$$f\left(\left\{\langle x_i \rangle, \langle x_i^2 \rangle\right\}\right)$$

Belief-Propagation equations

$$\left\{ \begin{array}{l} \langle x_i \rangle^{t+1} = \langle x_i \rangle^t + \frac{\partial f}{\partial \langle x_i \rangle} \\ \langle x_i^2 \rangle^{t+1} = \langle x_i^2 \rangle^t + \frac{\partial f}{\partial \langle x_i^2 \rangle} \end{array} \right.$$

$$U_{i}^{(t+1)} = \frac{\alpha}{M} \sum_{\mu} \frac{1}{\Delta_{\mu} + \gamma^{(t)}}$$

$$V_{i}^{(t+1)} = \sum_{\mu} F_{\mu i} \frac{(y_{\mu} - \alpha_{\mu}^{(t)})}{\Delta_{\mu} + \gamma_{\mu}^{(t)}} + f_{a} \left(U_{i}^{(t)}, V_{i}^{(t)}\right) \frac{\alpha}{M} \sum_{\mu} \frac{1}{\Delta_{\mu} + \gamma^{(t)}}$$

$$\alpha_{\mu}^{(t+1)} = \sum_{i} F_{\mu i} f_{a} (U_{i}^{(t+1)}, V_{i}^{(t+1)}) - \frac{(y_{\mu} - \alpha_{\mu}^{(t)})}{\Delta_{\mu} + \gamma^{(t)}} \frac{1}{N} \sum_{i} \frac{\partial f_{a}}{\partial Y} \left(U_{i}^{(t+1)}, V_{i}^{(t+1)}\right)$$

$$\gamma^{(t+1)} = \frac{1}{N} \sum_{i} f_{c} (U_{i}^{(t+1)}, V_{i}^{(t+1)})$$

Using these functions:

$$f_a(X,Y) = \left[\frac{\rho Y}{(1+X)^{3/2}} e^{Y^2/(2(1+X))}\right] \left[1 - \rho + \frac{\rho}{(1+X)^{1/2}} e^{Y^2/(2(1+X))}\right]^{-1}$$
$$f_c(X,Y) = \left[\frac{\rho}{(1+X)^{3/2}} e^{Y^2/(2(1+X))} \left(1 + \frac{Y^2}{1+X}\right)\right] \left[1 - \rho + \frac{\rho}{(1+X)^{1/2}} e^{Y^2/(2(1+X))}\right]^{-1} - f_a(X,Y)^2$$

And finally at the end:

 $\langle x_i \rangle = f_a \left(U_i, V_i \right)$



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And finally at the end: Complexity is O(N²×convergence time)

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Complexity is $O(N^2 \times \text{convergence time})$

http://aspics.krzakala.org

http://kl1p.sourceforge.net/home.html

Steepest ascent of the free entropy



Steepest ascent of the free entropy



 $E = \frac{1}{N} \sum_{i} \left(\langle x_i \rangle - x_i^0 \right)^2$

Thermodynamic potential

BP convergence time



- Maximum is at E=0 (as long as $\alpha > \rho 0$): Equilibrium behavior dominated by the original signal
- For α < 0.58, a secondary maximum appears (meta-stable state): spinodal point
- A steepest ascent dynamics starting from large E reaches the signal for α >0.58, but stay blocked in the meta-stable state for α <0.58, even if the true maximum is at E=0.
- Similarity with the physics of supercooled liquids

Computing the Phase Diagram



Computing the Phase Diagram


BP + probabilistic approach

$$P(\vec{x}|\vec{y}) = \frac{1}{Z} \prod_{i=1}^{N} \left[(1-\rho) \,\delta(x_i) + \rho \phi(x_i) \right] \prod_{\mu=1}^{M} \delta\left(y_\mu - \sum_{i=1}^{N} F_{\mu i} x_i \right)$$

BP + probabilistic approach

- Efficient and fast
- Robust to many type of noises (measurement, matrix coefficients, etc..)
- Very flexible (more information can be put in the prior, correlated variables, etc...)

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This is good, but not good enough

How to pass the spinodal point?



The dynamics is stuck in a metastable state, just as a liquid cooled too fast remains in a supercooled liquid state instead of crystalizing

This is good, but not good enough



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2) Add a first neighbor coupling



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3) Choose parameters such that the first system is in the region of the phase diagram where there is no metastability



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L = 8

 $N_i = N/L$ $M_i = \alpha_i N/L$



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$$\alpha_{1} > \alpha_{BP}$$

$$\alpha_{j} = \alpha' < \alpha_{BP} \qquad j \ge 2$$

$$\alpha = \frac{1}{L} (\alpha_{1} + (L-1)\alpha')$$



Block I has a large value of M such that the solution arise in this block...

... and then propagate in the whole system!

L = 8

$$N_i = N/L$$
$$M_i = \alpha_i N/L$$

: unit coupling

- : coupling J_1
- \square : coupling J_2
- : no coupling (null elements)

$$\alpha_1 > \alpha_{BP}$$

$$\alpha_j = \alpha' < \alpha_{BP} \qquad j \ge 2$$

$$\alpha = \frac{1}{L} (\alpha_1 + (L-1)\alpha')$$

Replica solution for coupled seeded matrix

The order parameters are now

$$Q_p \equiv \frac{1}{N_p} \sum_{i \in B_p} \langle x_i^2 \rangle \,, \quad q_p \equiv \frac{1}{N_p} \sum_{i \in B_p} \langle x_i \rangle^2 \,, \quad m_p \equiv \frac{1}{N_p} \sum_{i \in B_p} s_i \langle x_i \rangle$$

in each block $p \in \{1, \ldots, L_c\}$. The free entropy analogous to that in Eq. (112) becomes

$$\begin{split} &\Phi(\{Q_p\}_{p=1}^{L_c}, \{q_p\}_{p=1}^{L_c}, \{m_p\}_{p=1}^{L_c}, \{\hat{Q}_p\}_{p=1}^{L_c}, \{\hat{q}_p\}_{p=1}^{L_c}, \{\hat{m}_p\}_{p=1}^{L_c}) = \\ &-\frac{1}{2}\sum_{q=1}^{L_r} n_1 \alpha_{q1} \left[\frac{\tilde{q}_q - 2\tilde{m}_q + \tilde{\rho}_q + \Delta_0}{\tilde{Q}_q - \tilde{q}_q + \Delta} + \log\left(\Delta + \tilde{Q}_q - \tilde{q}_q\right) \right] + \sum_{p=1}^{L_c} n_p \left(\frac{Q_p \hat{Q}_p}{2} - m_p \hat{m}_p + \frac{q_p \hat{q}_p}{2} \right) \\ &+ \sum_{p=1}^{L_c} n_p \int \mathrm{d}s \left[(1 - \rho_0) \delta(s) + \rho_0 \phi_0(s) \right] \int \mathcal{D}z \log \left\{ \int \mathrm{d}x \, e^{-\frac{\hat{Q}_p + \hat{q}_p}{2} x^2 + x(\hat{m}_p s + z\sqrt{\hat{q}_p})} \left[(1 - \rho) \delta(x) + \rho \phi(x) \right] \right\}, \end{split}$$

(after a bit of work...)

Comparing the algorithm and replica theory



Comparing the algorithm and replica theory



Best measurement rates reached!



A combination of Statistical physics technics (Bethe-Peierls, Replica) and concepts (dynamics, nucleation and growth) has allowed to solve a major problem is signal processing theory 0.15

1000

An example



Shepp-Logan phantom, in the Haar-wavelet representation

A more interesting example

 $\alpha = \rho \approx 0.24$



The Lena picture in the Haar-wavelet representation

Conclusions...

- A probabilistic approach to reconstruction
- Analysis of best possible reconstruction for different class of signals
- The Belief Propagation algorithm
- Optimality achieving seeded measurements matrices

... and perspectives:

- More information in the prior (Correlated measurement, wavelets, etc...)
- Other matrices with asymptotic measurements?
- Non-random matrix (e.g. Radon operator in Tomography, Fourier, etc..)
- Additive and multiplicative noise, Quasi-sparsity, etc...?
- Calibration, and matrix/dictionary learning?
- Applications ?

http://leshouches2013.krzakala.org

SPECIAL ANNOUNCEMENTS

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2 Post-doc openings on these topics for 2013

If you work in Statistical physics, Information science, Signal processing, etc...



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<u>COMING SOON:</u> An interdisciplinary school on these topics: Les Houches, October 2013, Organizers F. Krzakala & L. Zdeborová



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