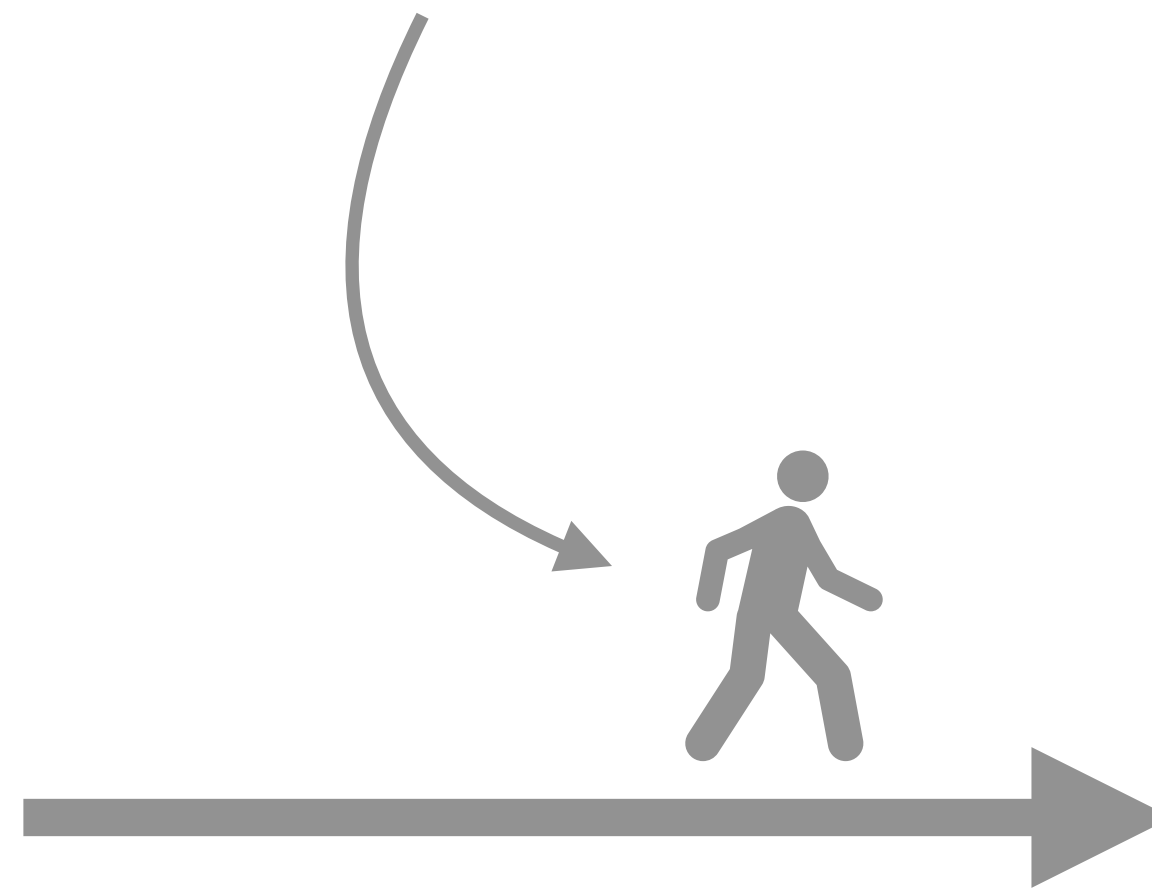


Nonreciprocal active matter

— Irreversibility, heat flows, and persistent currents

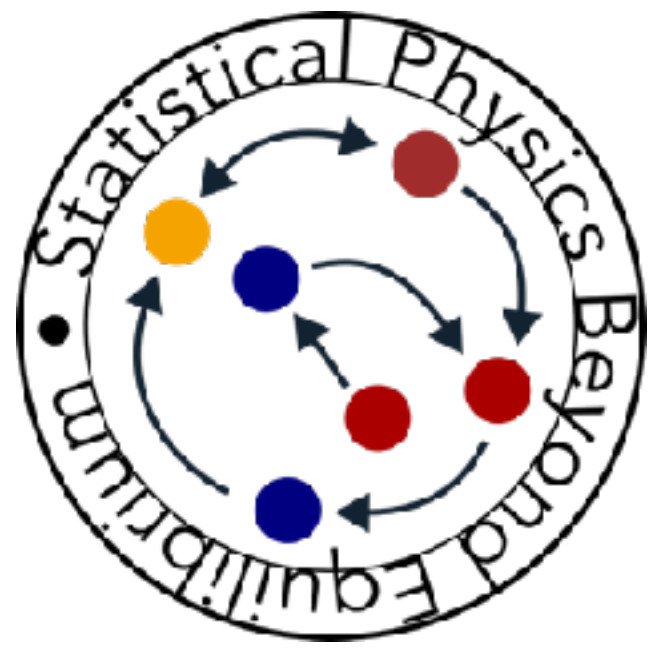
Sarah A. M. Loos



MAX PLANCK INSTITUTE
FOR DYNAMICS AND
SELF-ORGANIZATION



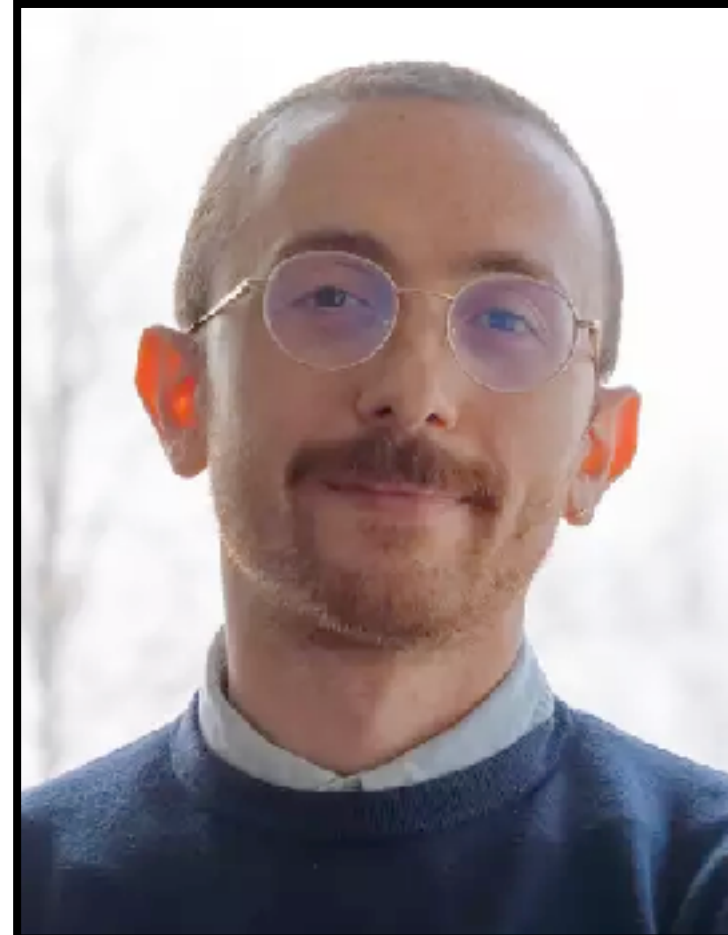
“Statistical Physics Beyond Equilibrium” - group



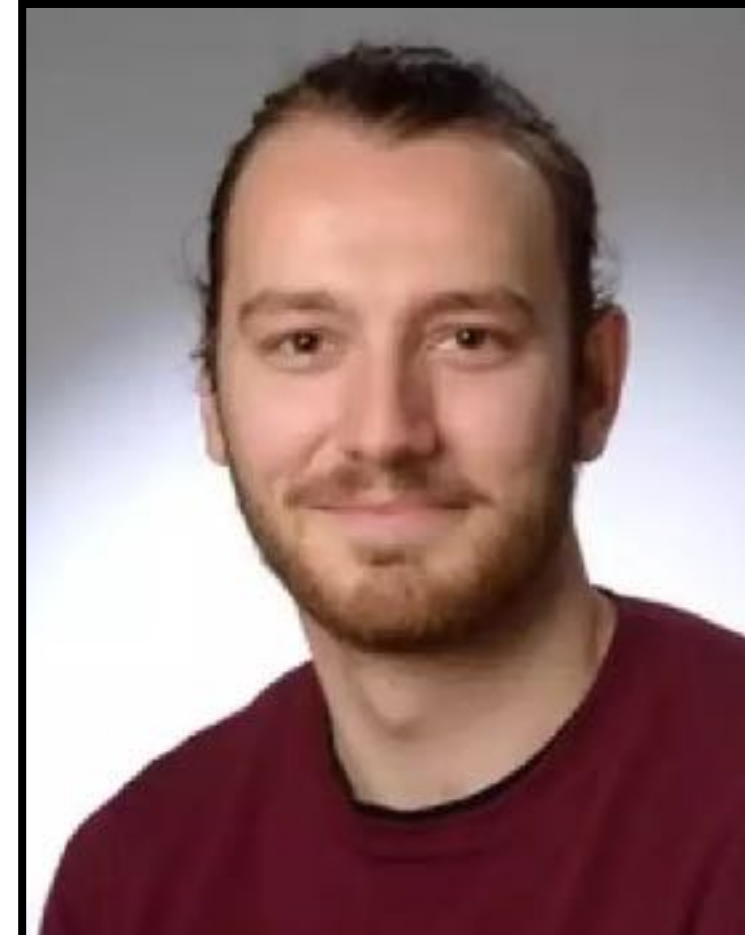
MAX PLANCK INSTITUTE
FOR DYNAMICS AND
SELF-ORGANIZATION



Sarah Loos



Gennaro Tucci
Postdoc



Lars Stutzer
PhD student



Thomas Suchanek
PhD student

*Joint supervision
with Klaus Kroy*



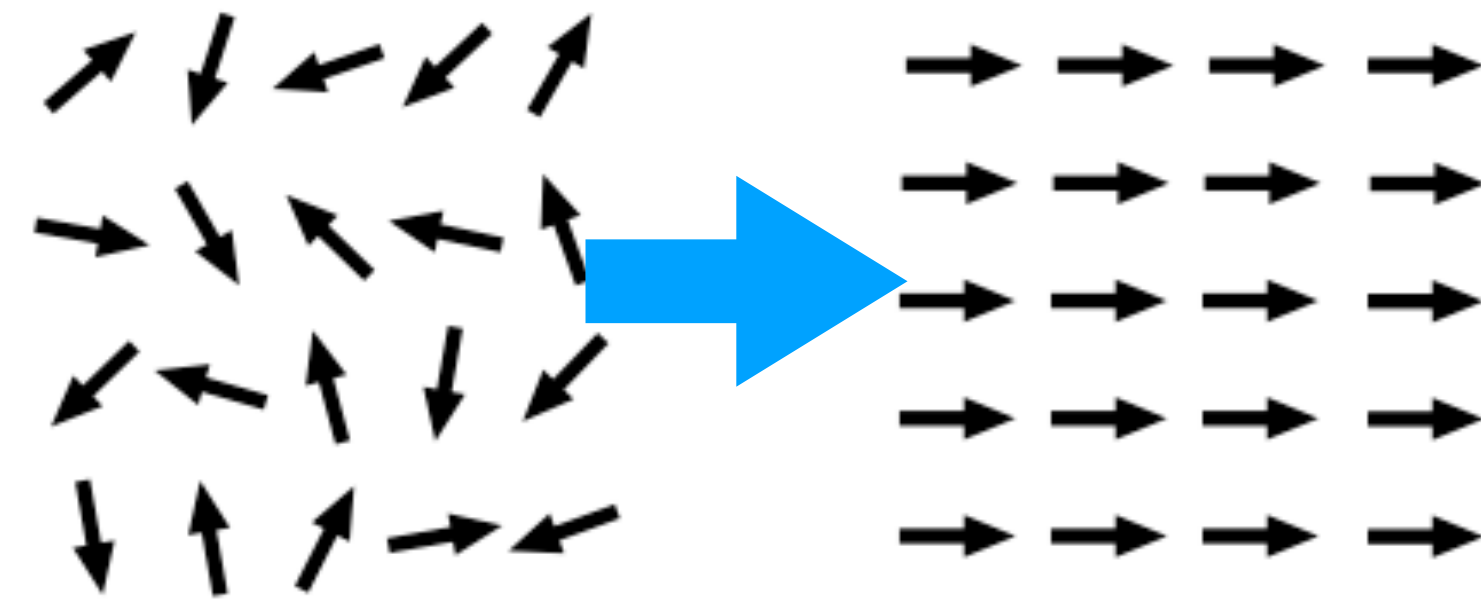
UNIVERSITÄT
LEIPZIG



UNIVERSITY OF
CAMBRIDGE

- **Goal:** Extend concepts from statical physics to far from equilibrium systems (including biological systems, synthetic systems under feedback control, microrobotic systems...)
- Focus: Noise, Memory, Nonreciprocity
- Incorporate ideas from **information theory** and **stochastic thermodynamics**

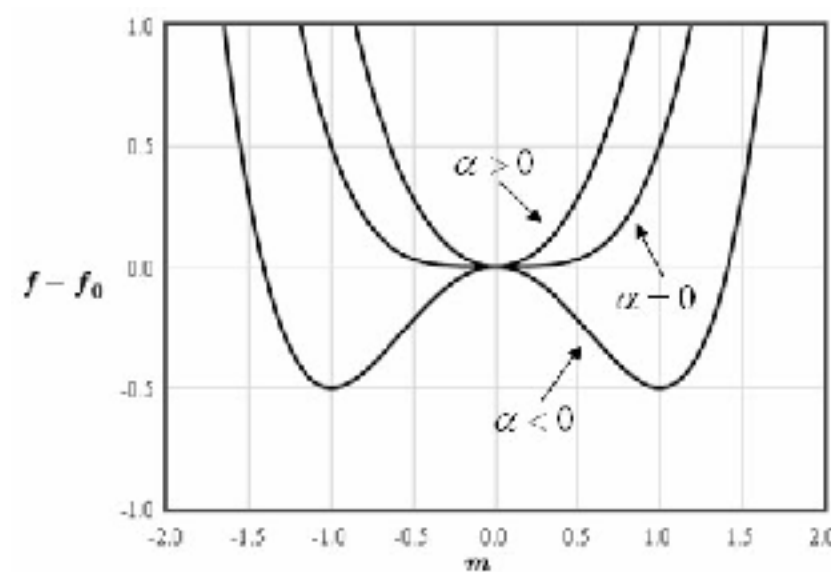
“Passive matter”



- Conservative Interactions:

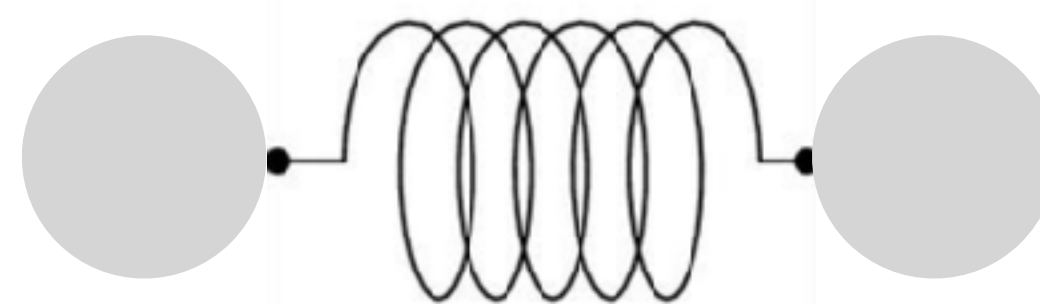
Free energy

(Ginsburg-Landau)



$$f(T) = f_0(T) + \alpha m^2 + \frac{1}{2} \beta m^4$$

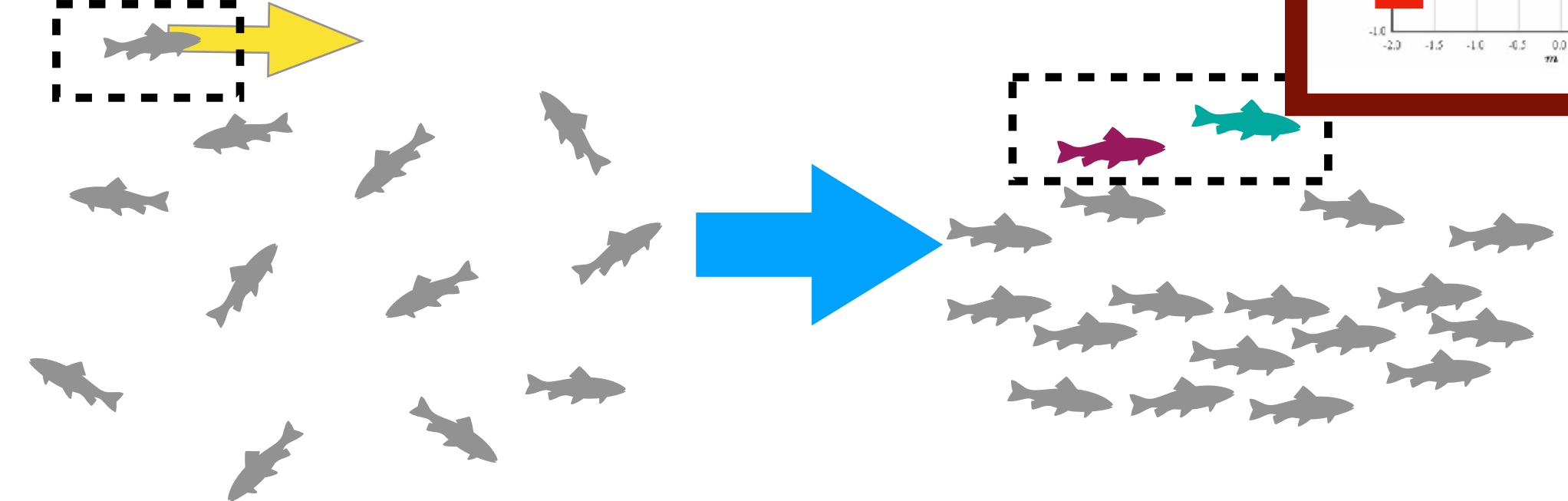
Reciprocity



Action = Reaction

vs. “Active matter”

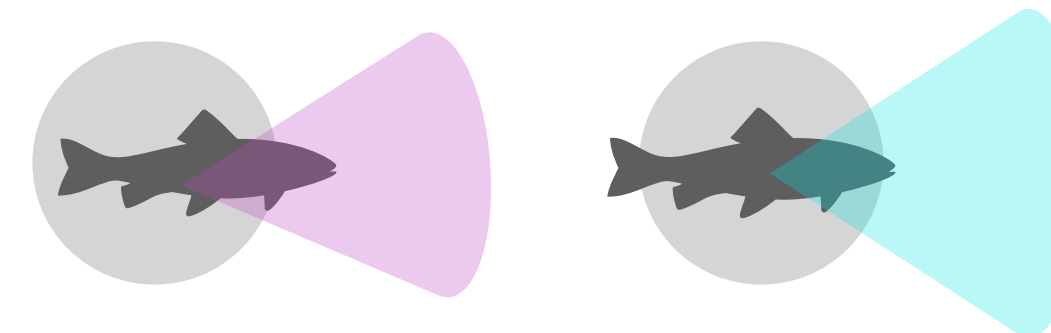
Self-propulsion



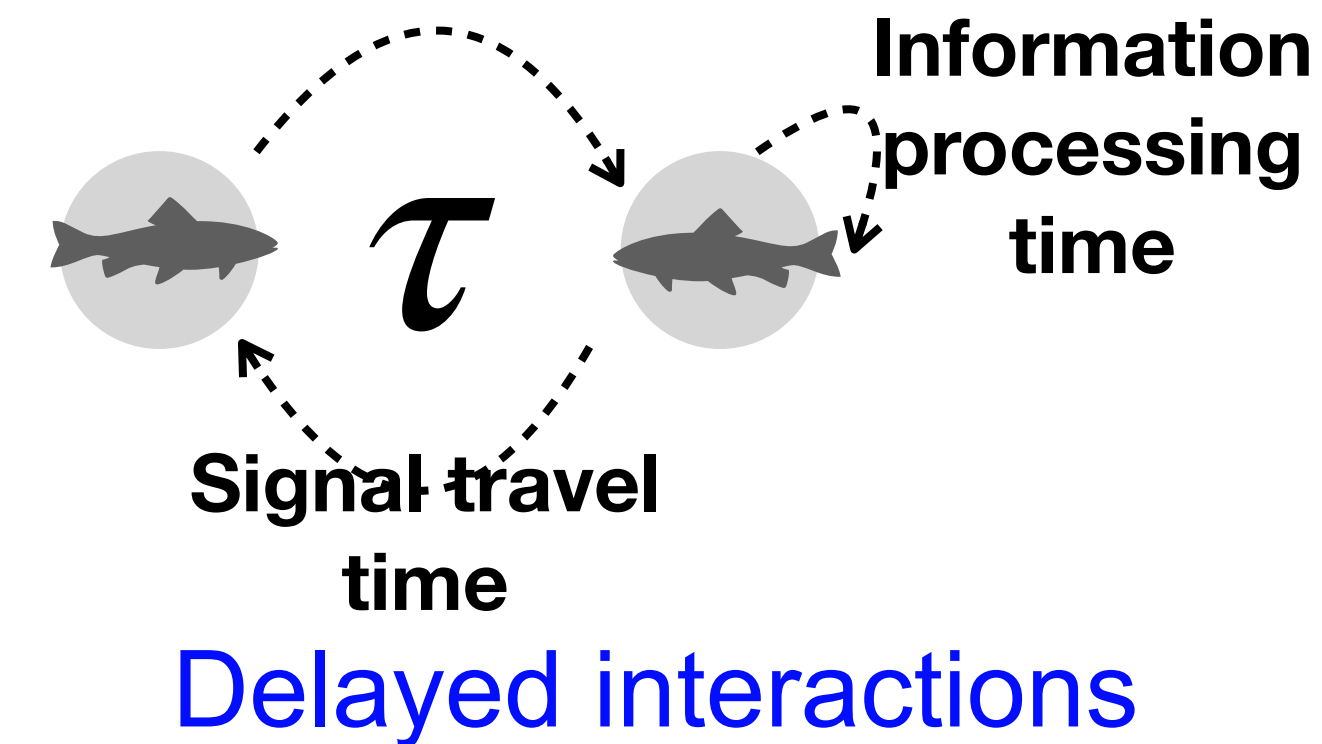
- Energy input** on level of individual “particles” $\sim X(t)$
(Not associated with explicit global symmetry breaking)

- Non-conservative* (effective) interactions:

Nonreciprocity (NR)



Action $\neq$ Reaction



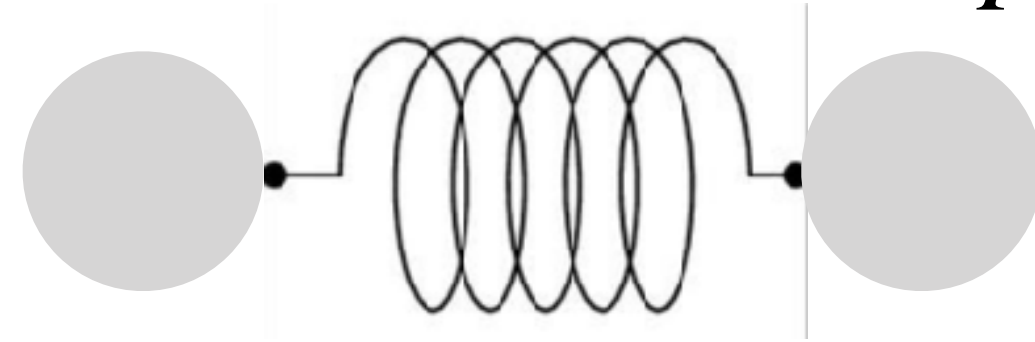
*Interaction forces not representable as ∇V

NR interactions in Nonequilibrium systems

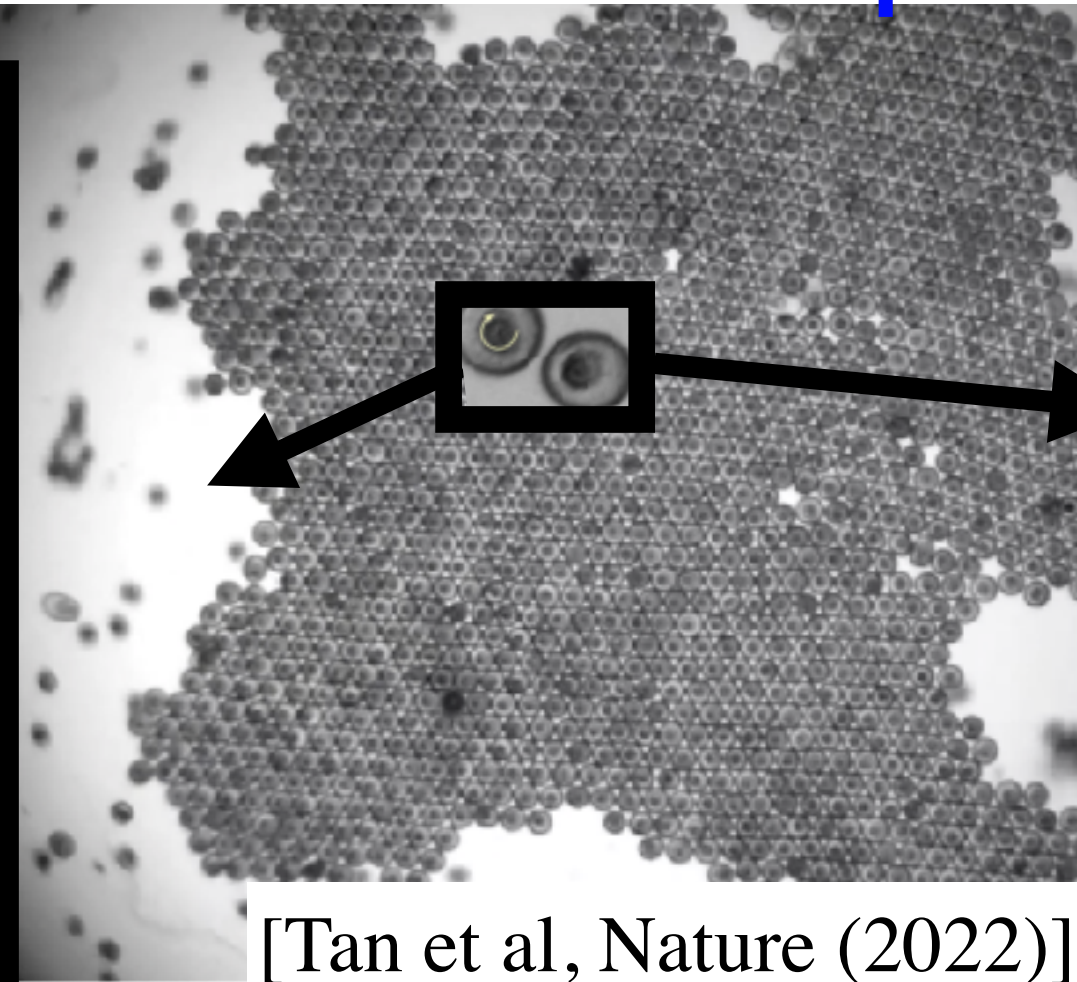
Equilibrium:

Reciprocity

$$F_{12} = -F_{21}$$

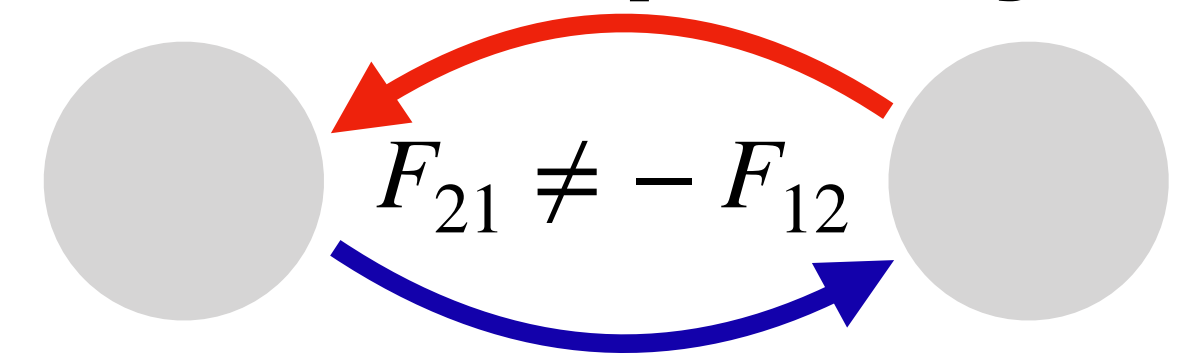


Action = Reaction



Non-Equilibrium:

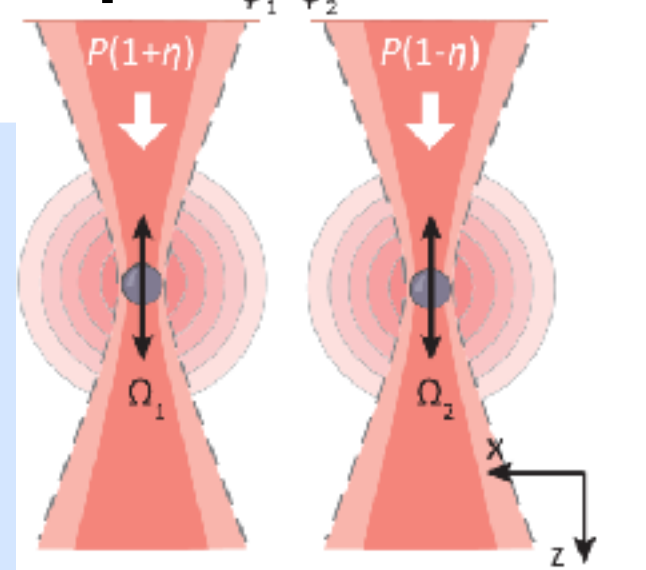
Non-reciprocity



Action $\neq$ Reaction

*Nonequilibrium complex systems with multiple scales
effective interactions often NR!*

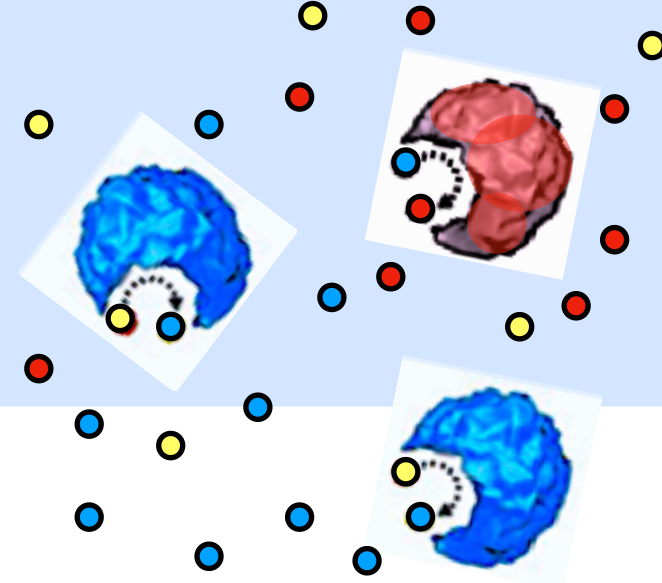
Optomechanics¹



nm

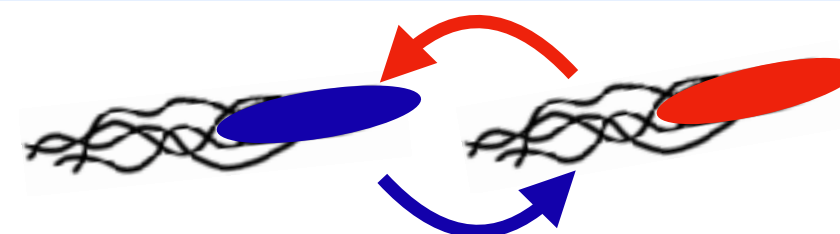
Chemical systems²

(enzymes or catalytically-active colloids)



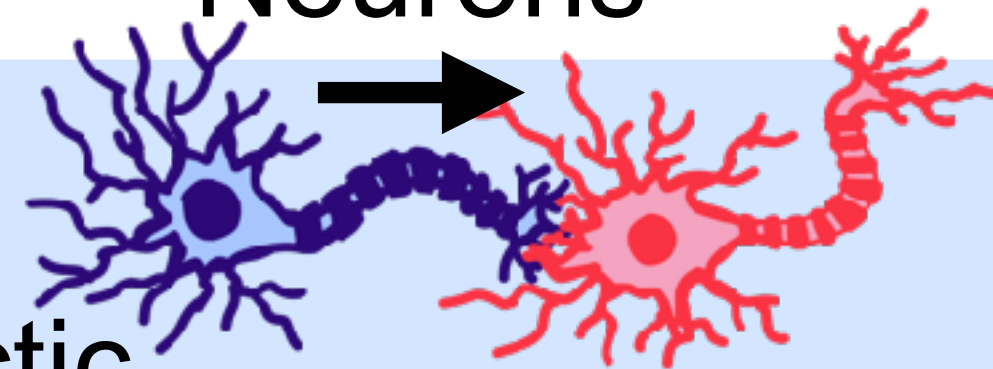
μm

Chemotactic bacteria³

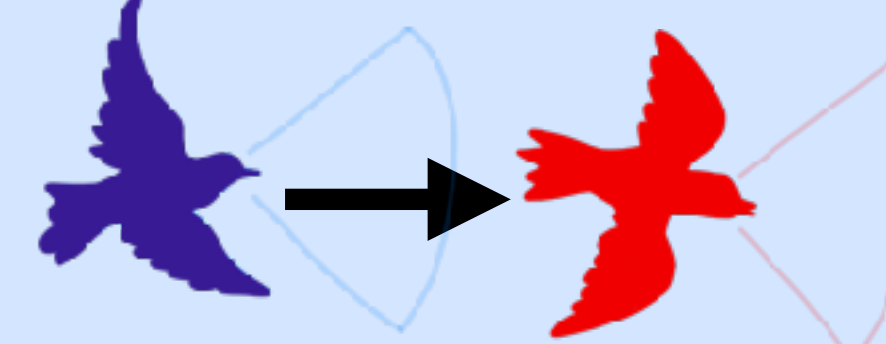


mm

Neurons⁴

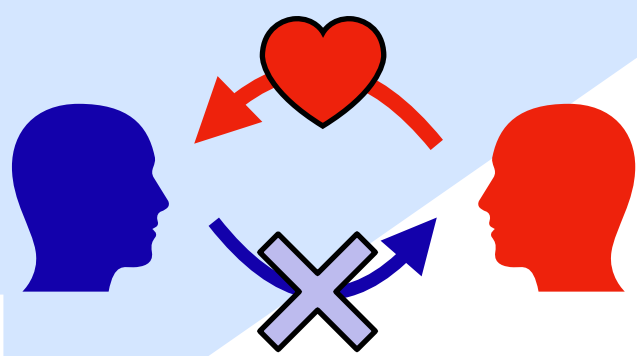


Visual interactions⁵



m

Social interactions



NR interactions in Nonequilibrium systems

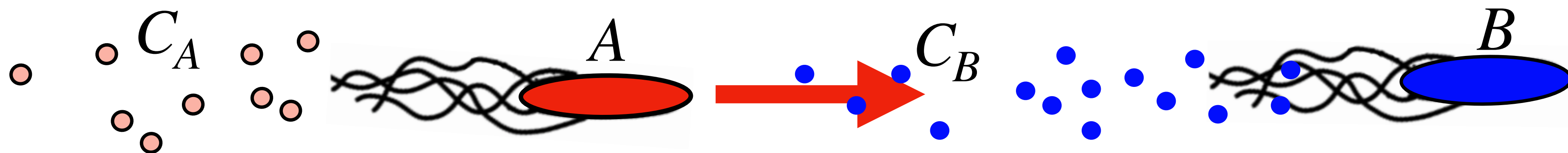
- Mixtures of Enzymes / Colloidal particles / Bacteria:

“Catalytic” activity
(consume/produce chemicals)

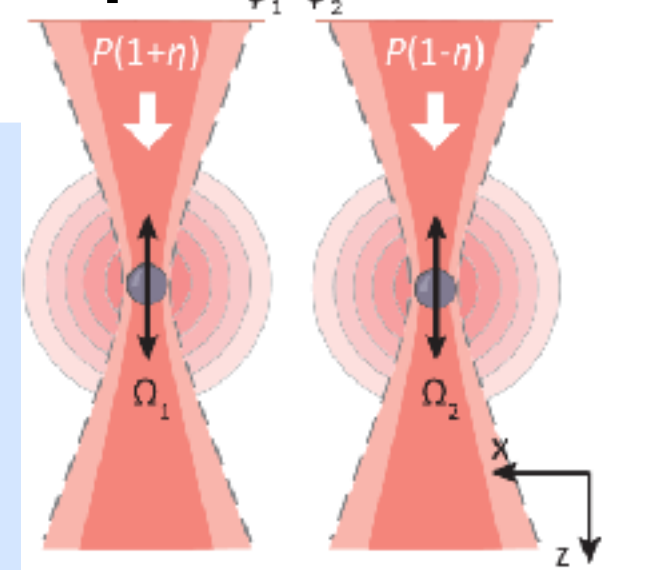
+

Chemophoresis / Chemotaxis
(move along chemical gradients)

Species A is chemotactically attracted to C_B (but not vice versa)



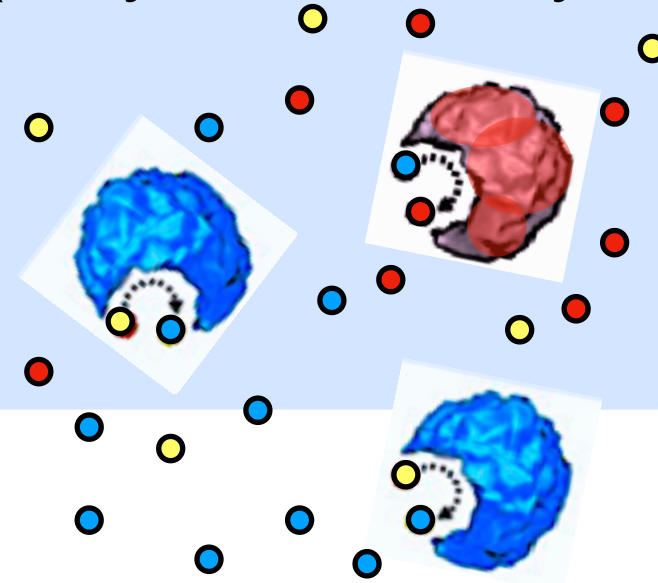
Optomechanics¹



nm

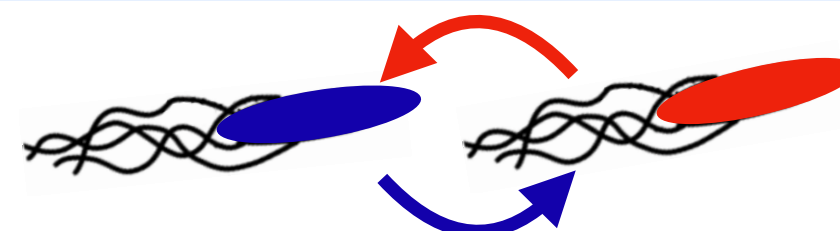
Chemical systems²

(enzymes or catalytically-active colloids)



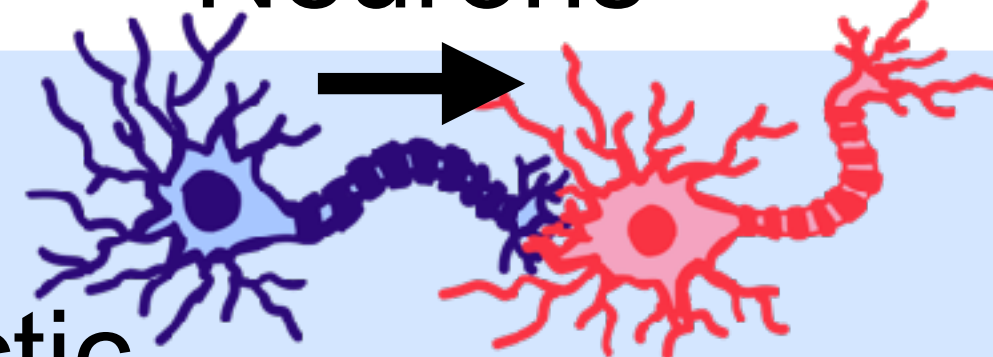
μm

Chemotactic bacteria³

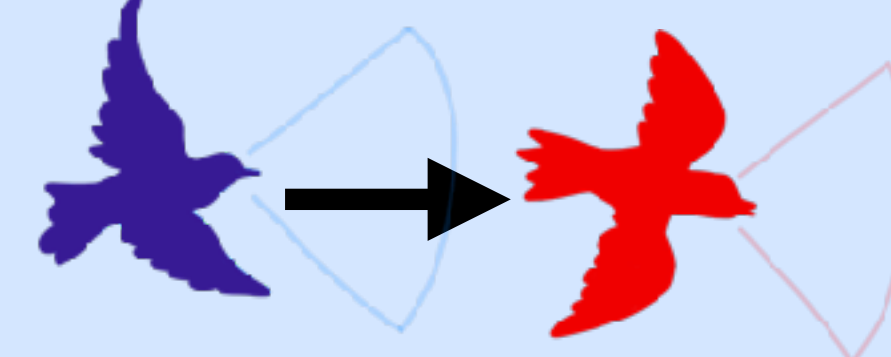


mm

Neurons⁴

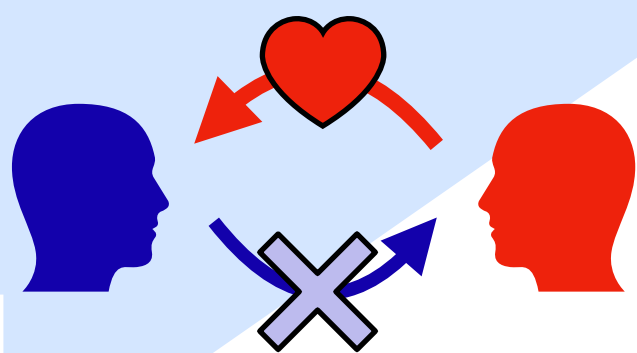


Visual interactions⁵



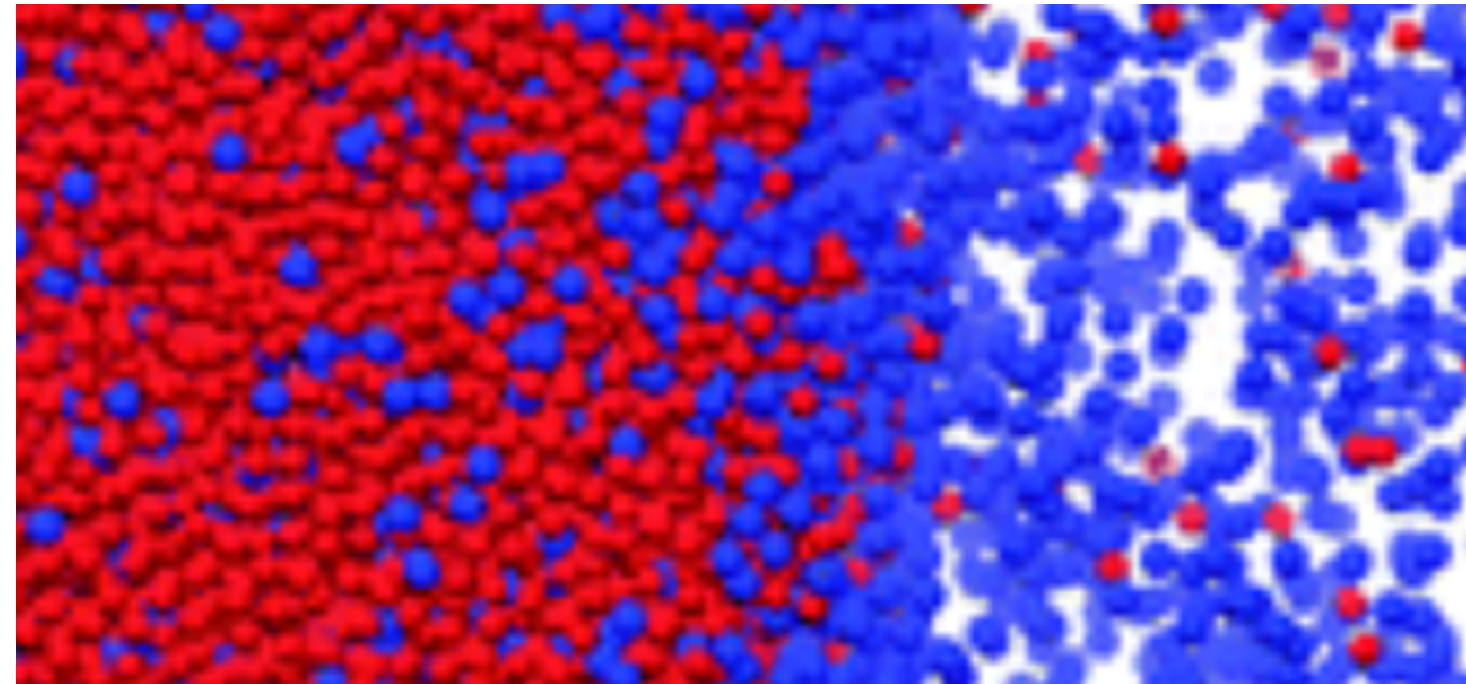
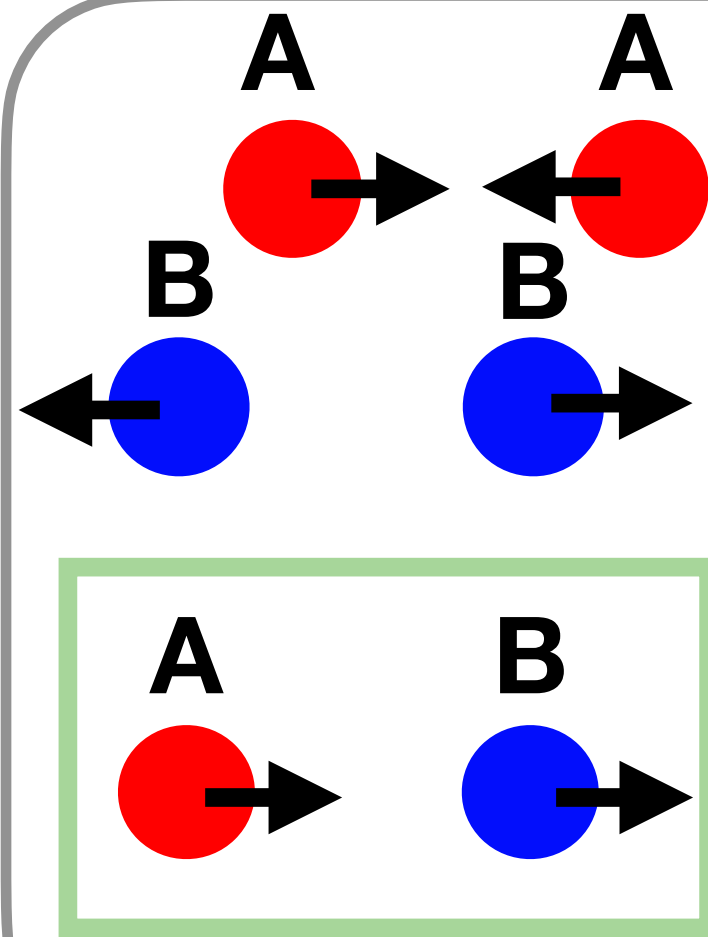
m

Social interactions



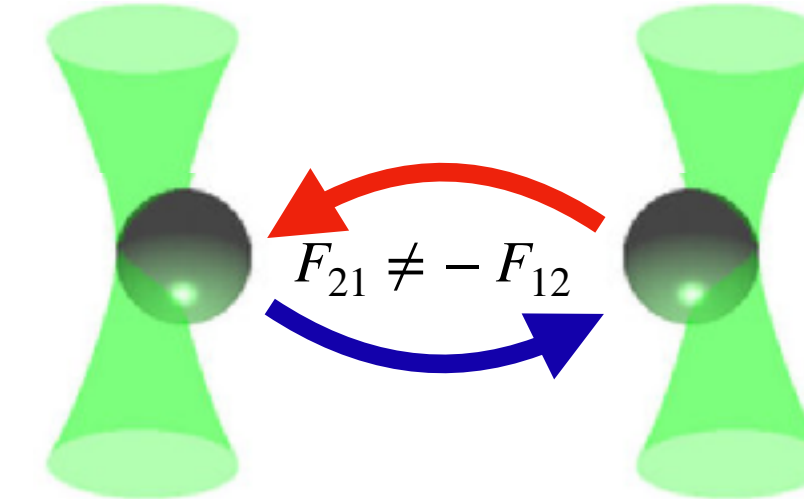
“Stat Phys of NR many-body systems”

- **Binary NR fluids**



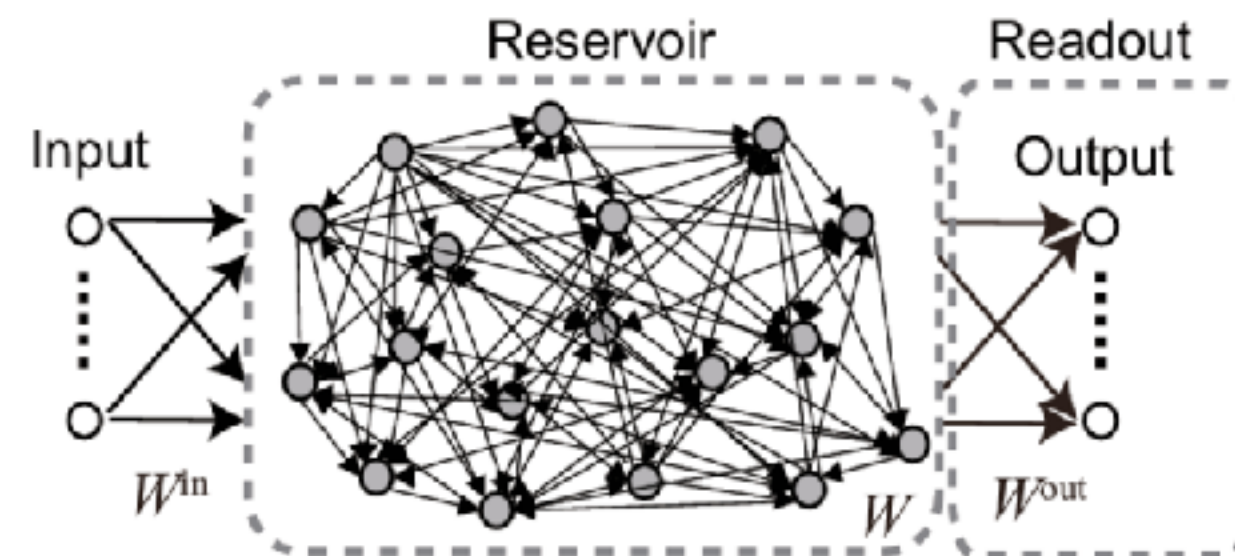
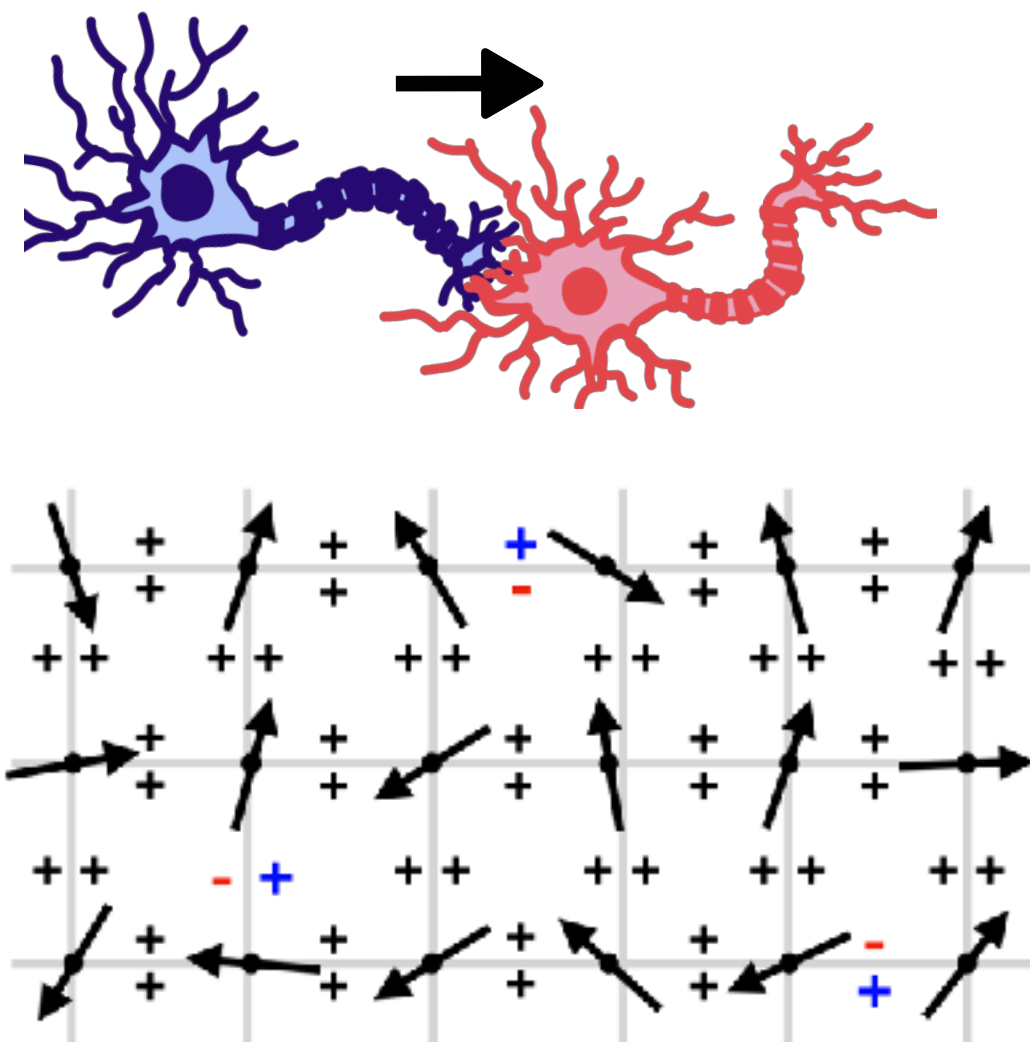
Suchanek, Kroy, **SL**, PRL (2024)

- **NR generated by optical feedback**



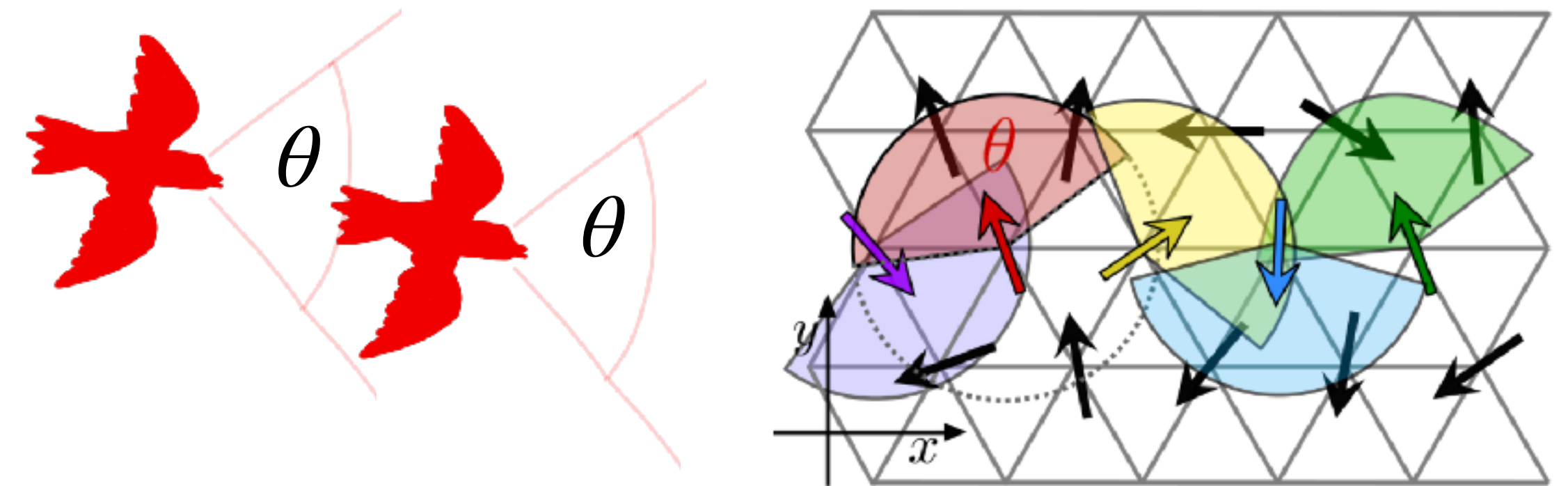
SL, Klapp, NJP (2020)

- **Quenched NR Disorder**



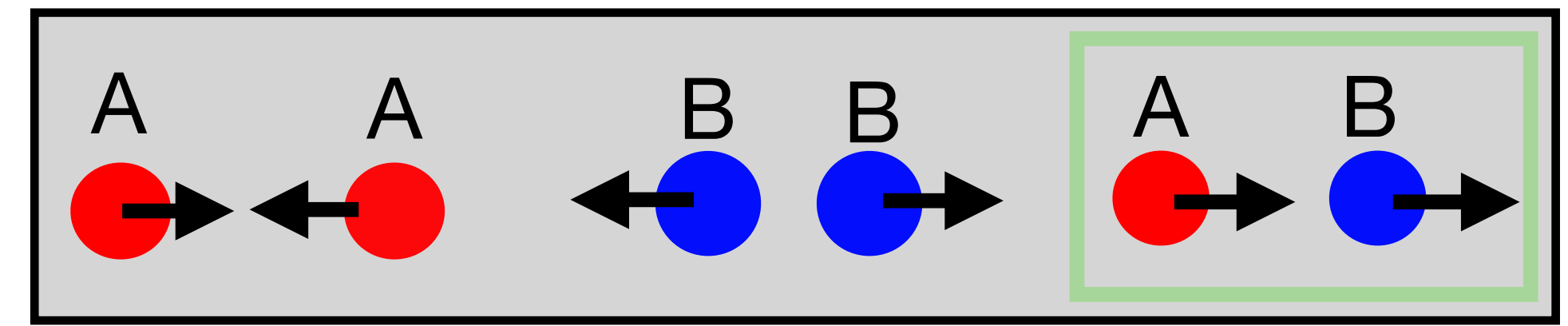
Suchanek, Cates, **SL**,
in preparation

- **Vision cone interactions**



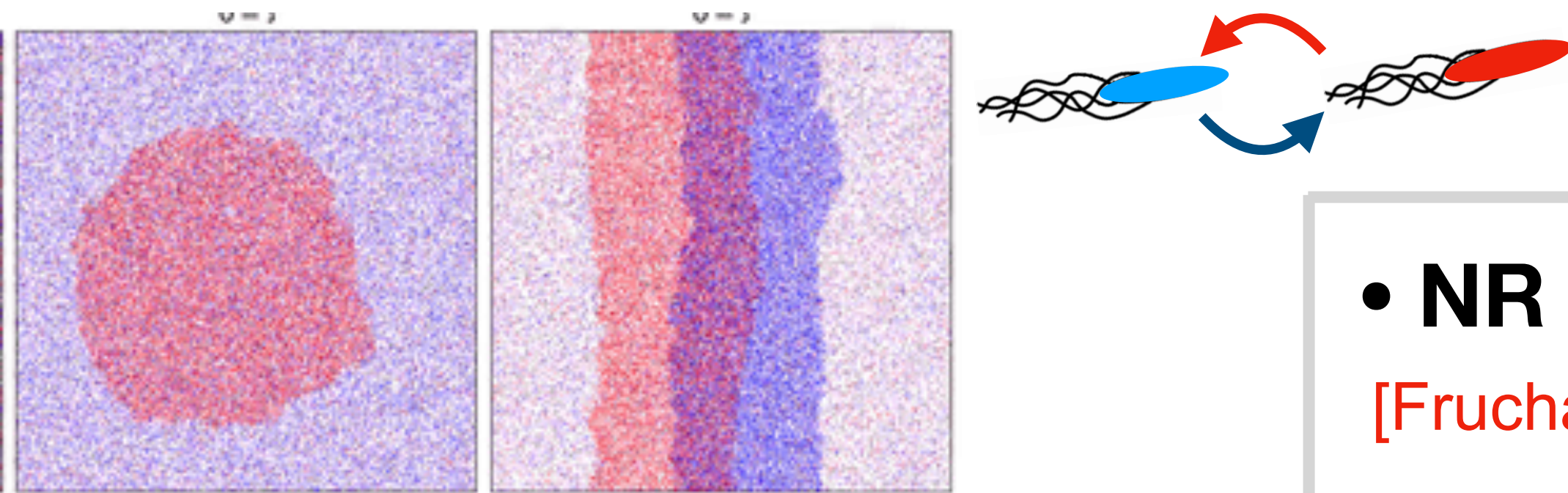
SL, Klapp & Martynek, PRL (2023)
Bandini, Venturelli, **SL**, Jelic, Gambassi,
J. Stat. Mech. (2025)

Inter-species NR interactions

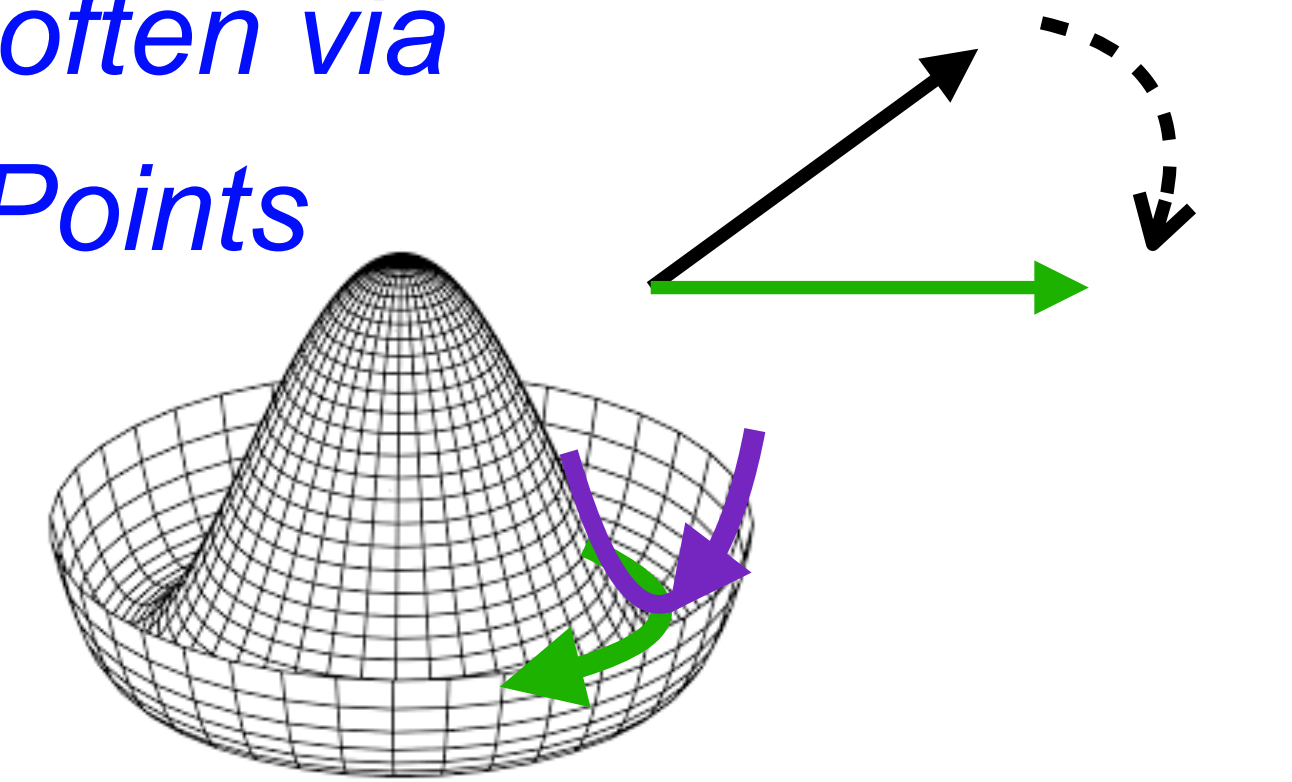


• Quorum-sensing chemotactic bacteria

[Dinelli et al., Nat Comm (2023), Duan et al., PRL (2023)]

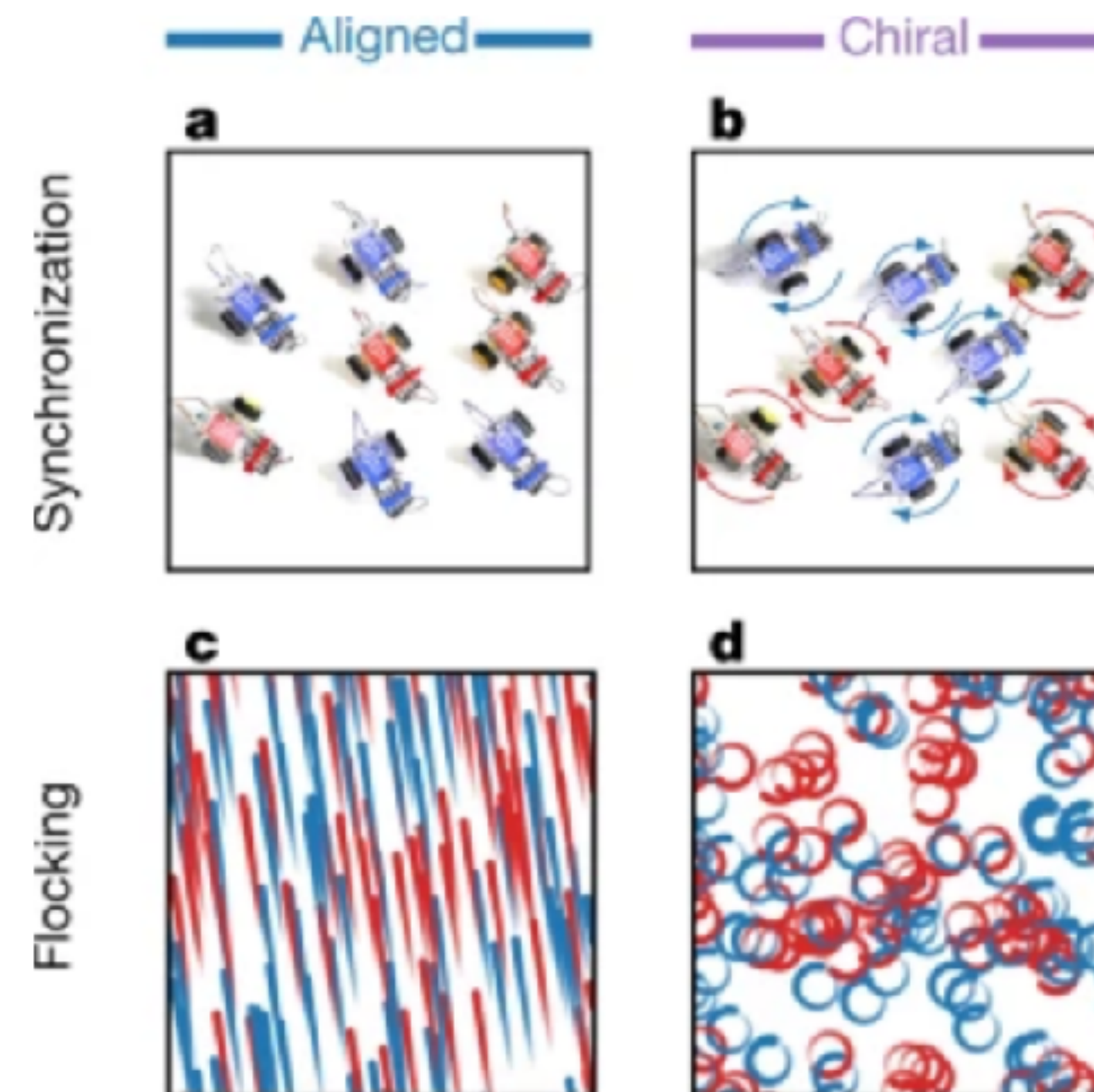


→ *Dynamical phases, often via Critical Exceptional Points*

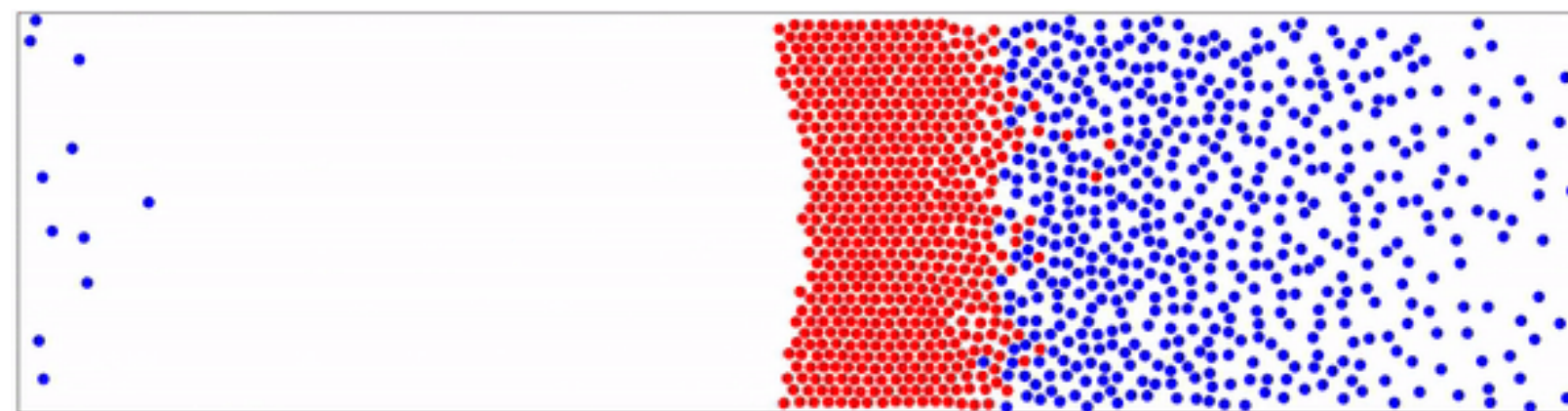


• NR Flocking, ...

[Fruchart et al., Nature (2021)]



• Binary mixtures (Active/Passive)



[You et al., PNAS (2020)]

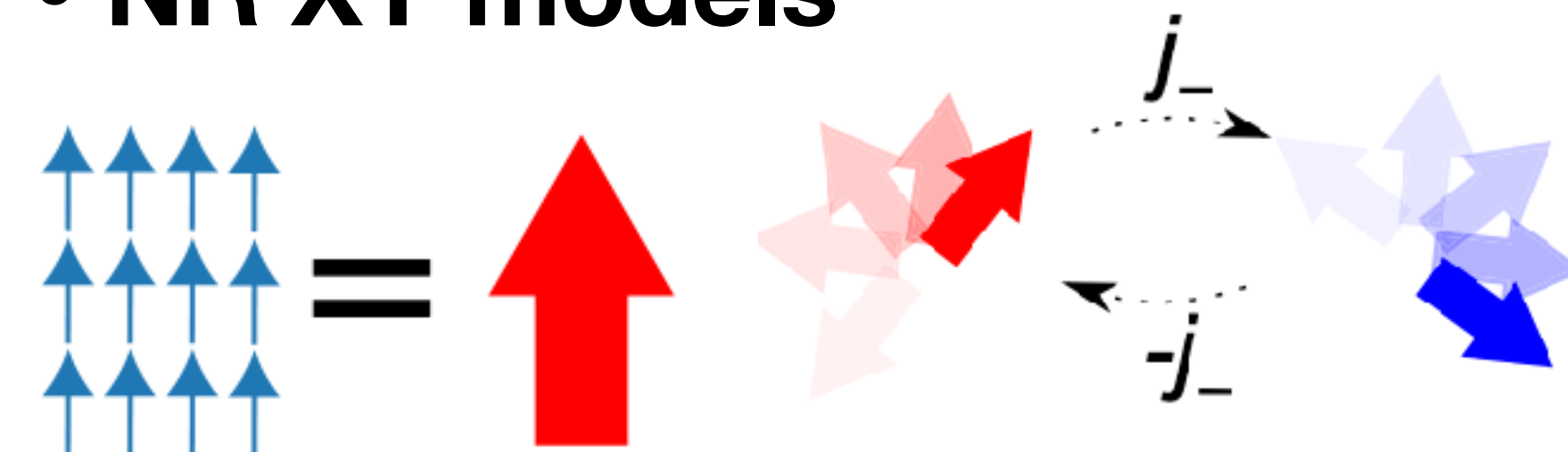
[Saha et al, PRX (2020)]

[Mason et al., arXiv (2024)]

[Mandal et al., PRE (2024)]

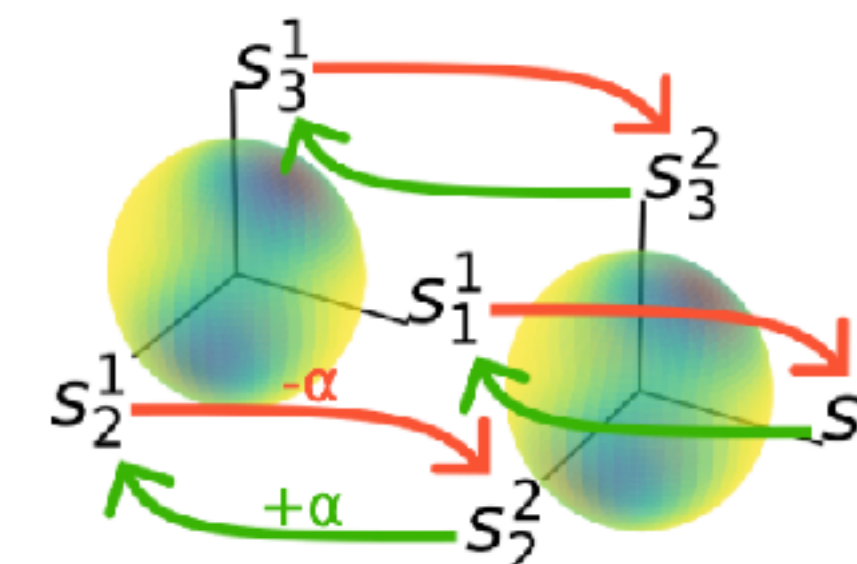
[Brauns et al., PRX (2025)] ...

• NR XY models

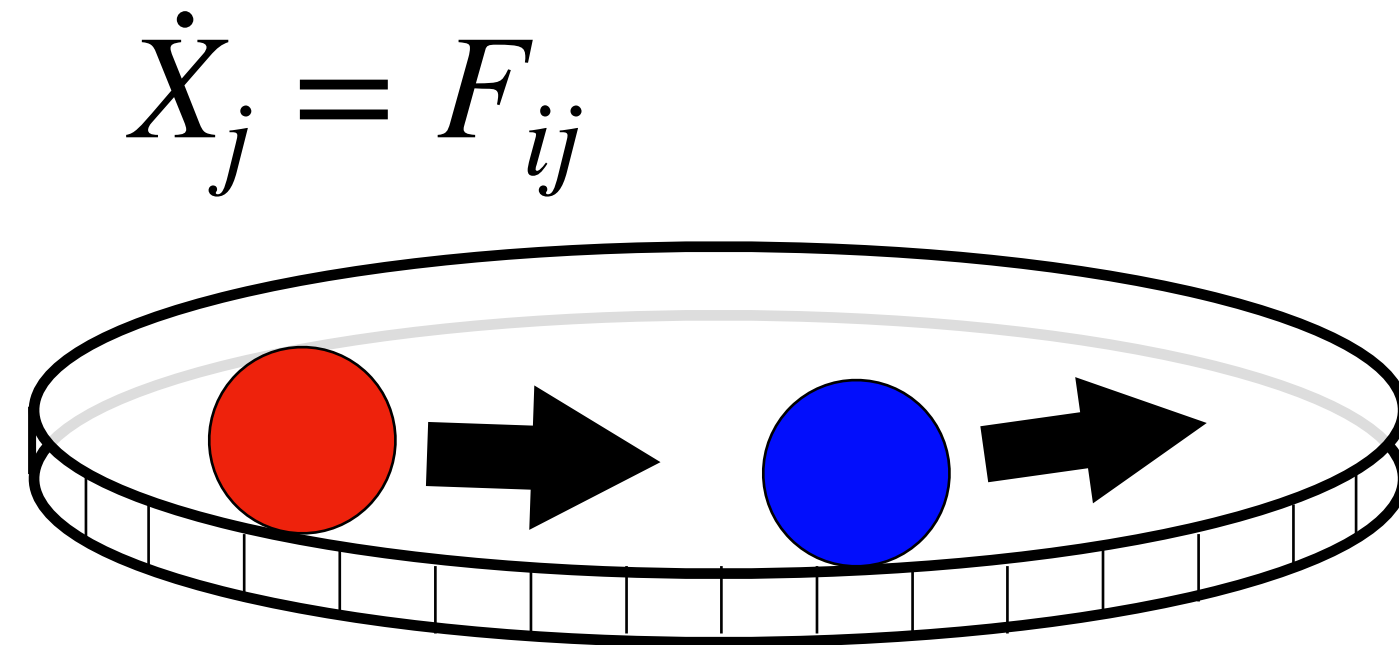


[Hanai, PRX (2024)]

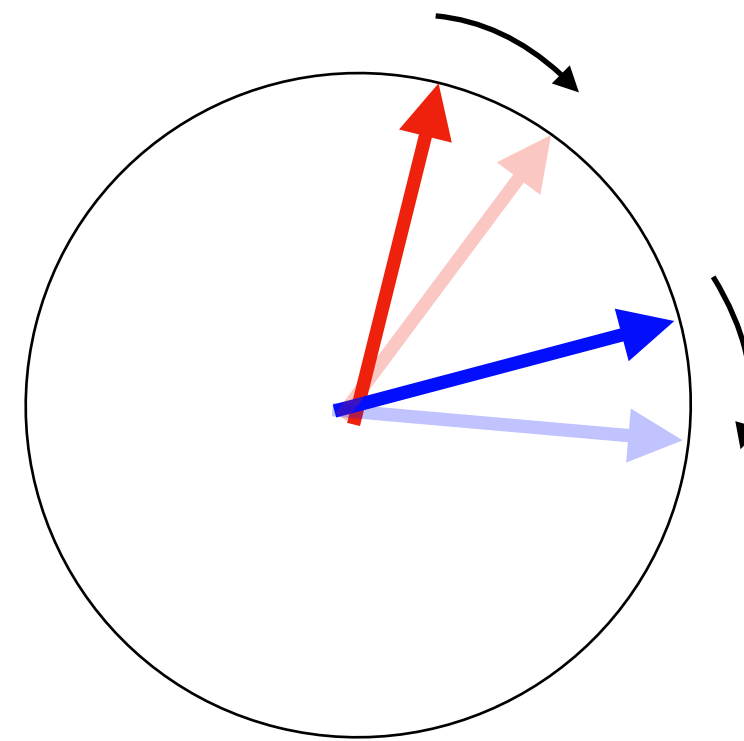
[Garcia Lorenzana, ArXiv (2024)]



Nonreciprocal $\rightarrow$ Symmetry breaking



$$F_{ij} = -J_{ij}(X_i - X_j)$$



$$F_{ij} = -a_{ij} \sin(X_i - X_j)$$

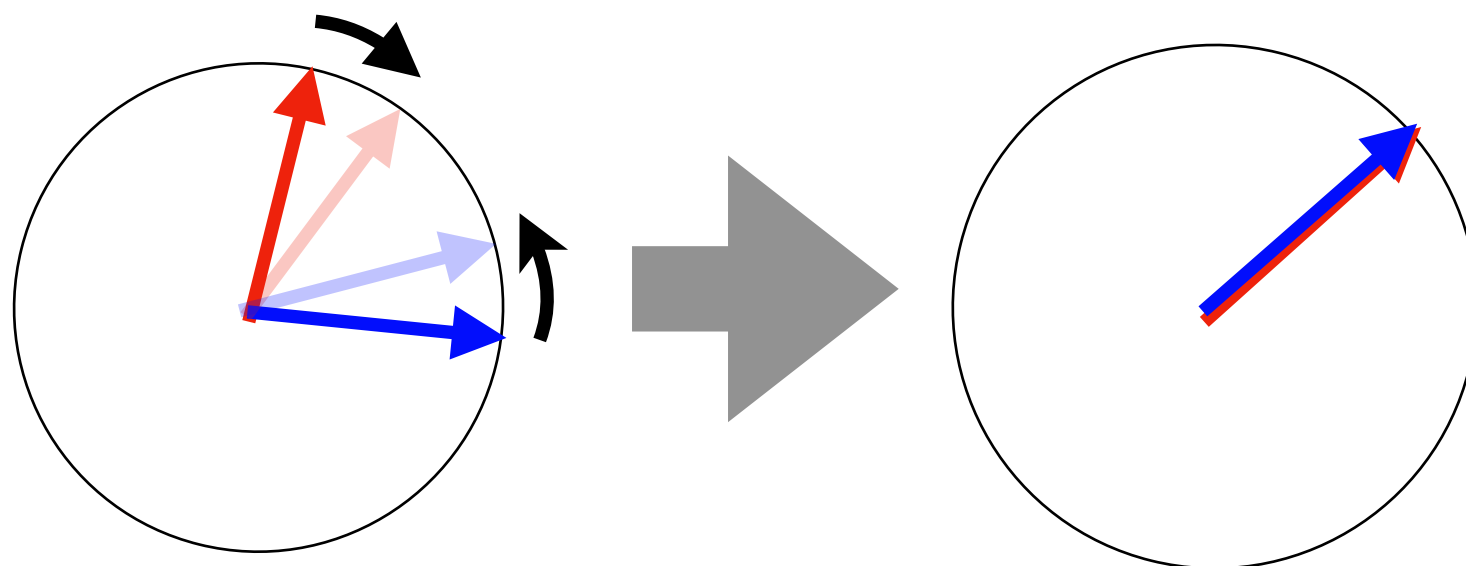
$$J_{ij} \neq J_{ji}$$

$$F_{ij} + F_{ji} \neq 0$$

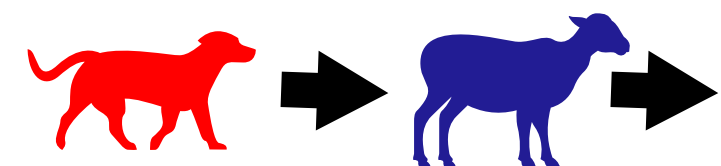
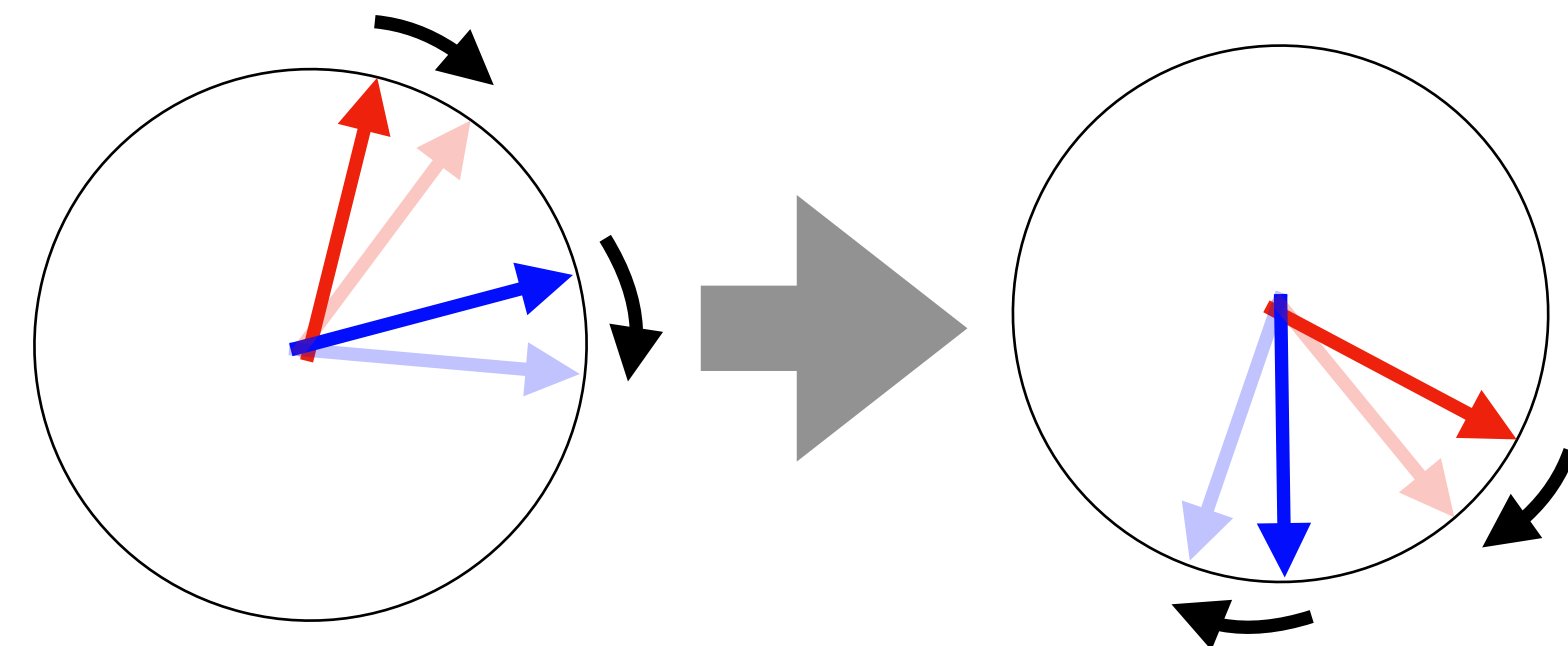
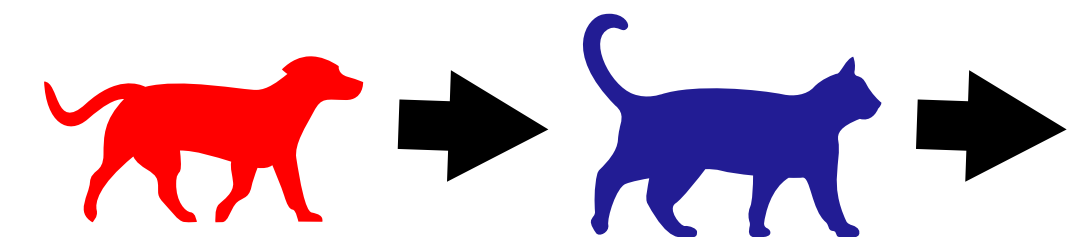
NR $\Rightarrow$ Linear momentum conservation broken

Particle Current?

$$J_{ij}J_{ji} > 0$$

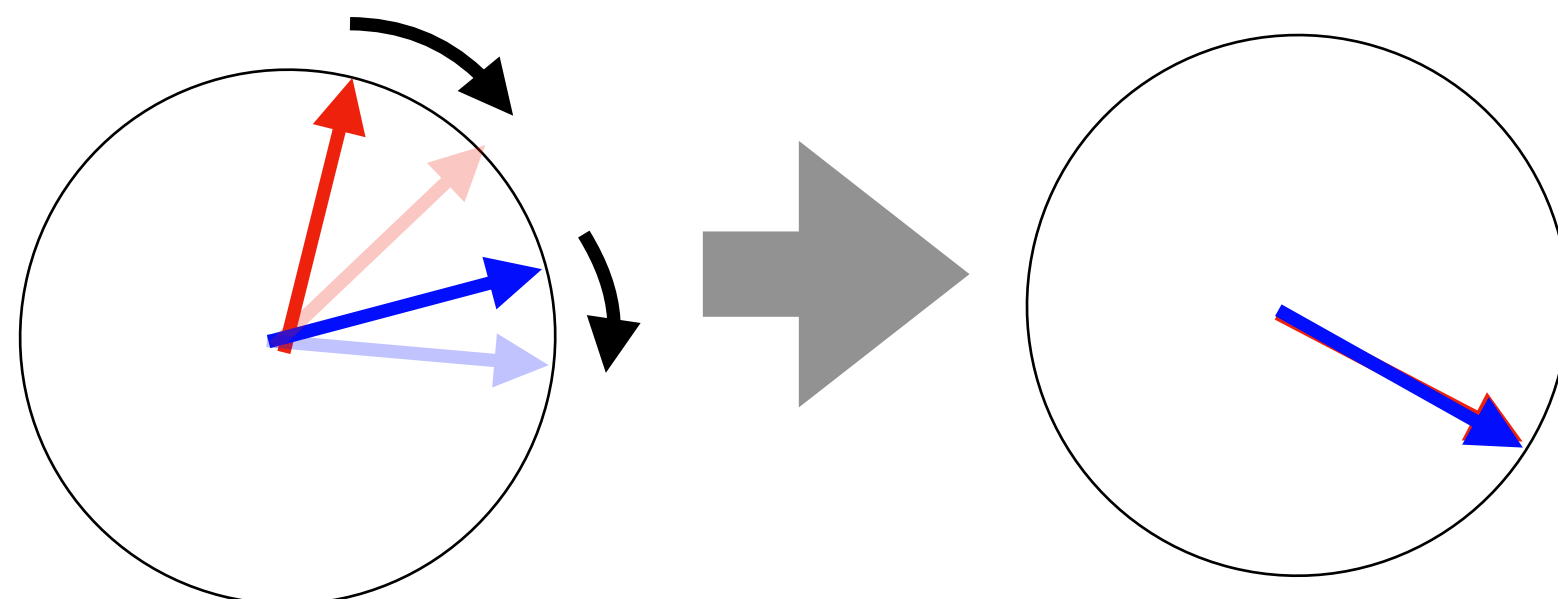


$$\text{Anti-reciprocal } J_{ij} = -J_{ji}$$



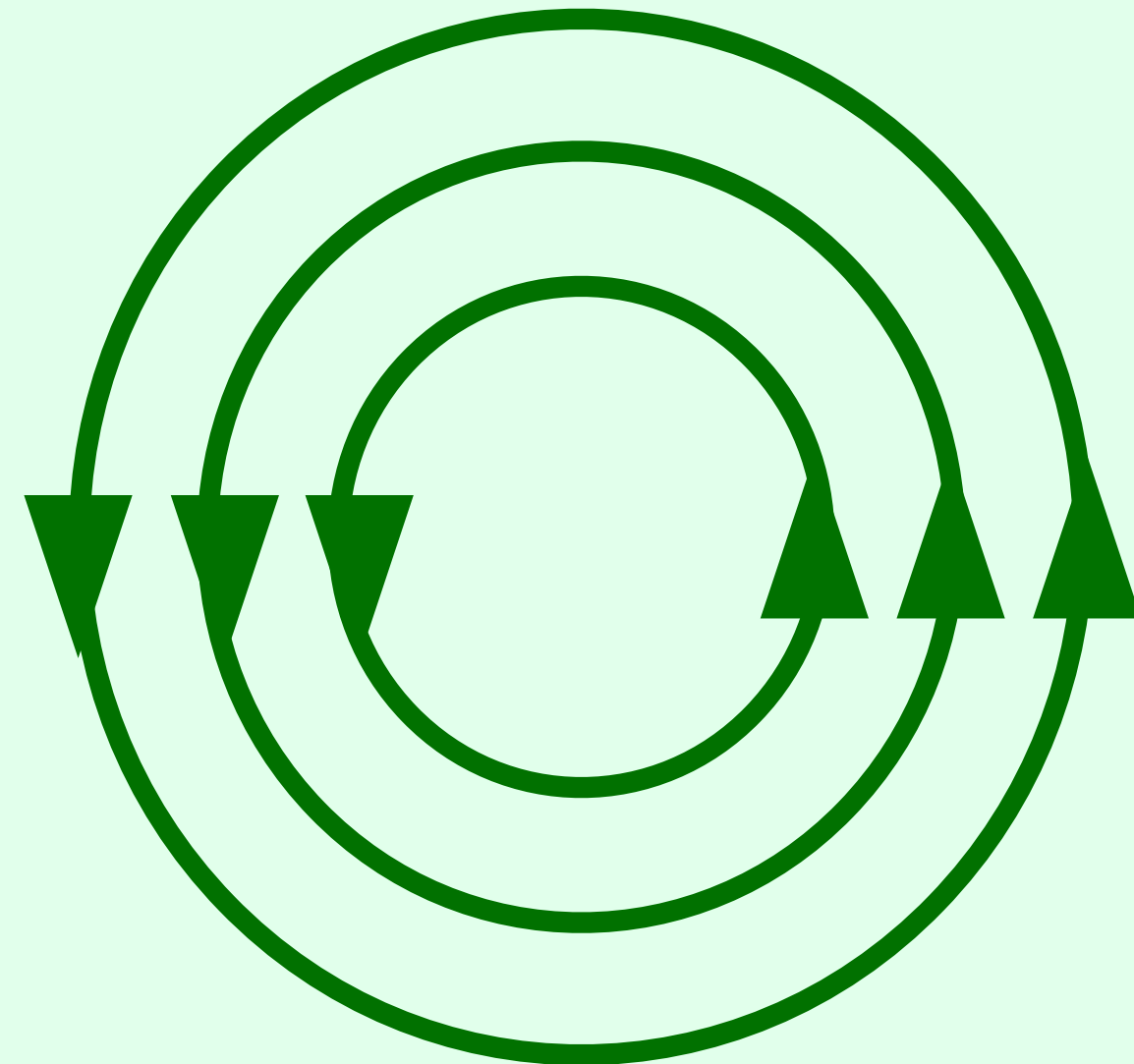
$$J_{ij}J_{ji} < 0$$

$$|J_{ij}| > |J_{ji}|$$



Nonreciprocal $\rightarrow$ Symmetry breaking

$$\dot{X}_i = F_{ii}$$



$$\dot{X} \propto \Delta\phi$$

Marginal orbits $\rightarrow$ Not stable
against noise $\rightarrow$ No persistent current

Rotation due to period domain +
“anti-reciprocity” $J_{ij} = -J_{ji}$

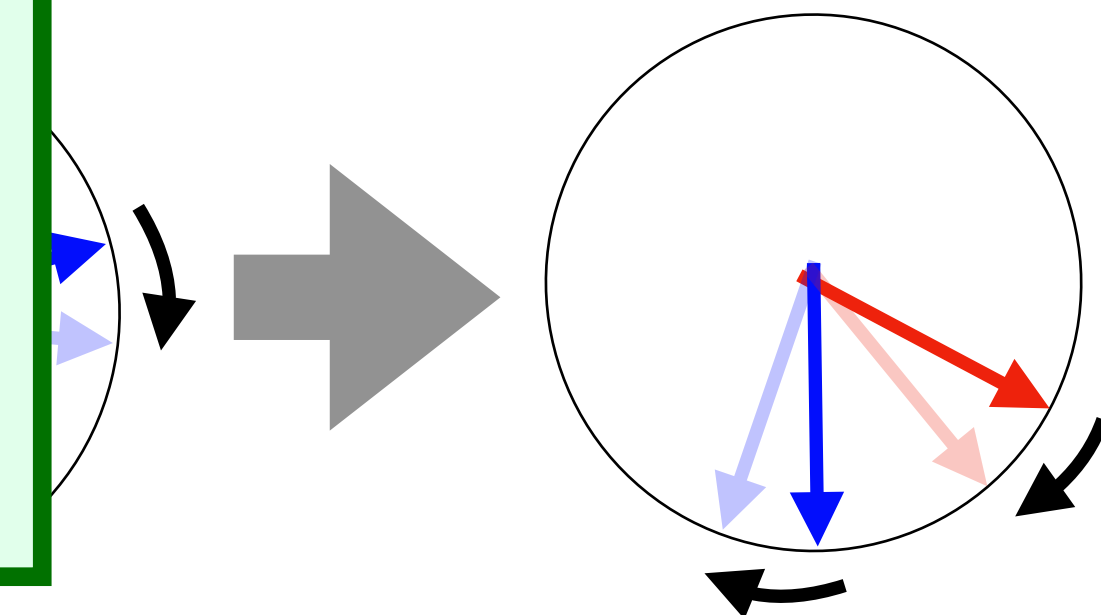
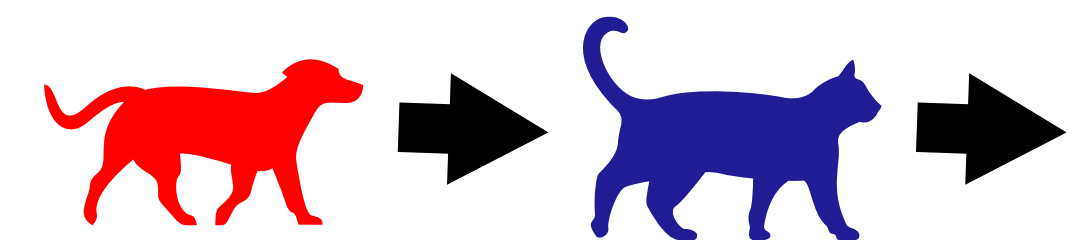
$$J_{ij} \neq J_{ji}$$

$$F_{ij} + F_{ji} \neq 0$$

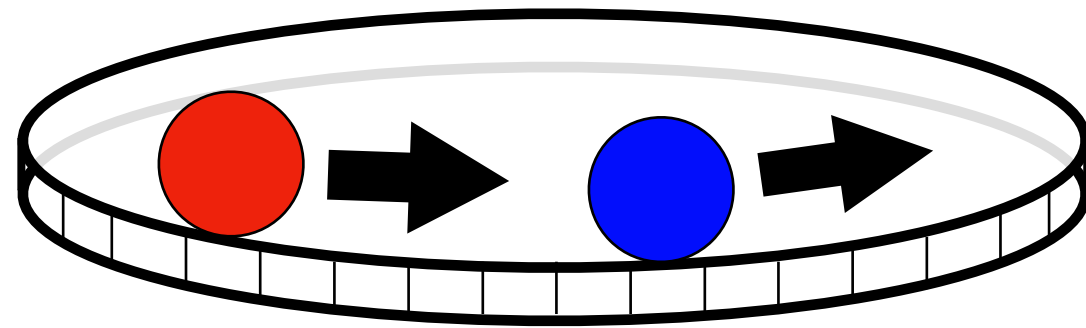
IR $\Rightarrow$ **Linear momentum
conservation broken**

Particle Current?

$$\text{cal } J_{ij} = -J_{ji}$$



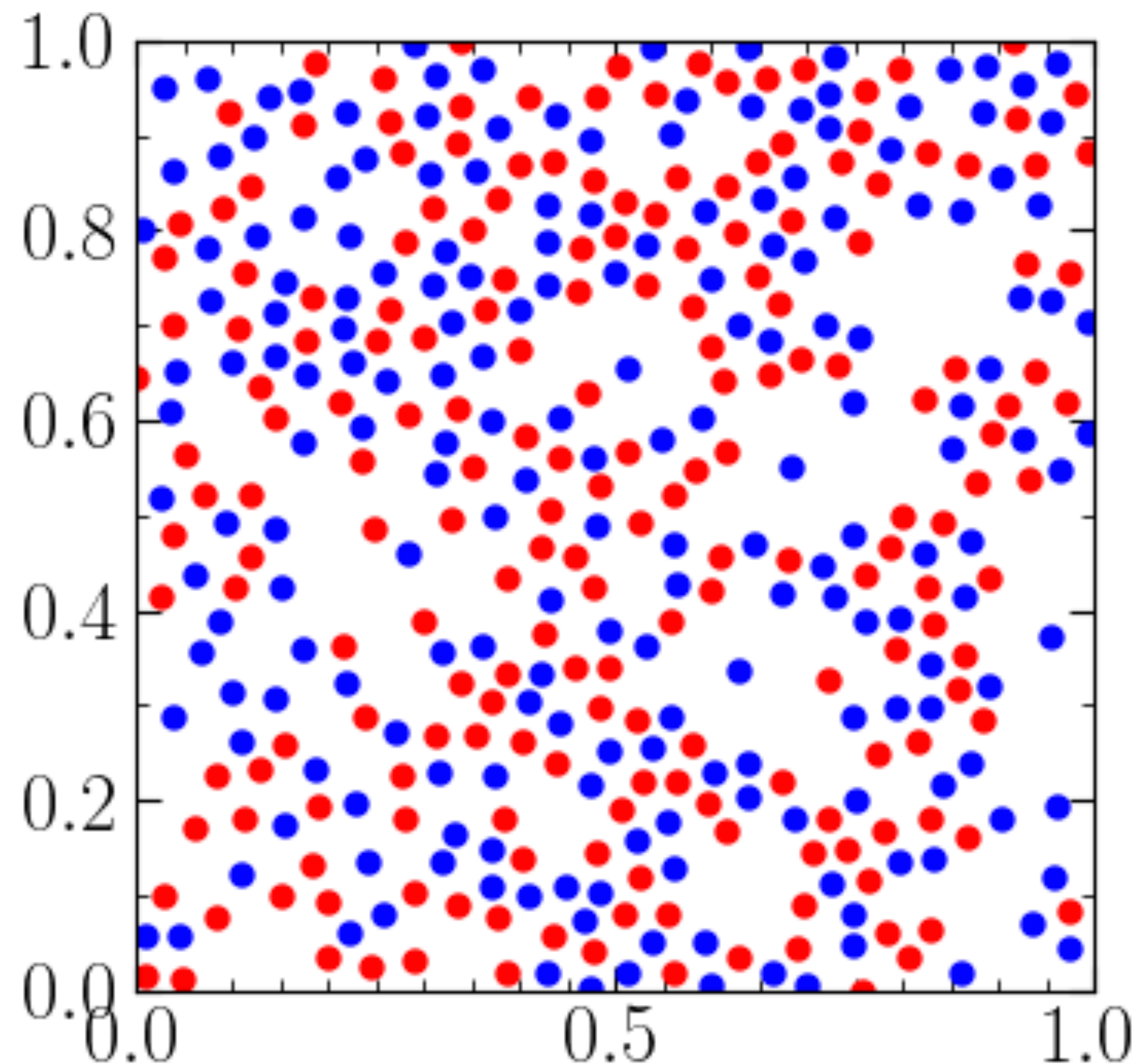
Nonreciprocal $\rightarrow$ Symmetry breaking



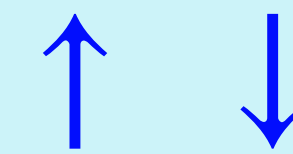
NR $\Rightarrow$ **Linear momentum conservation broken**

$\Rightarrow$ **Time-reversal symmetry breaking (TRSB)**

SL, Klapp, NJP (2020)



Non-reciprocity (NR)



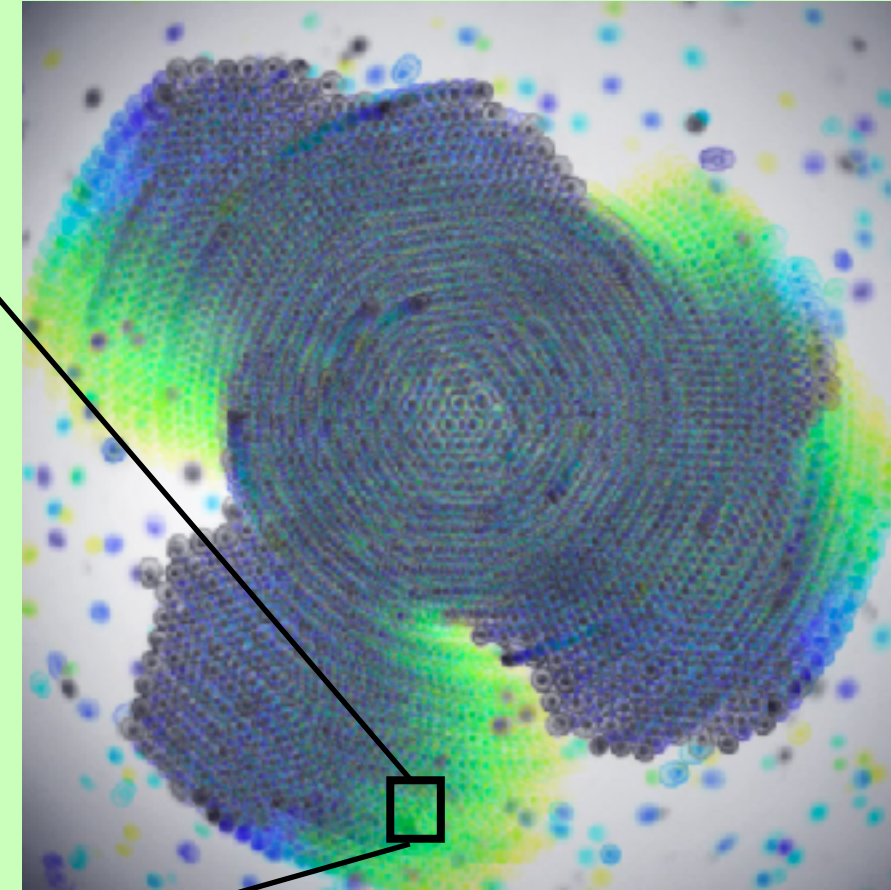
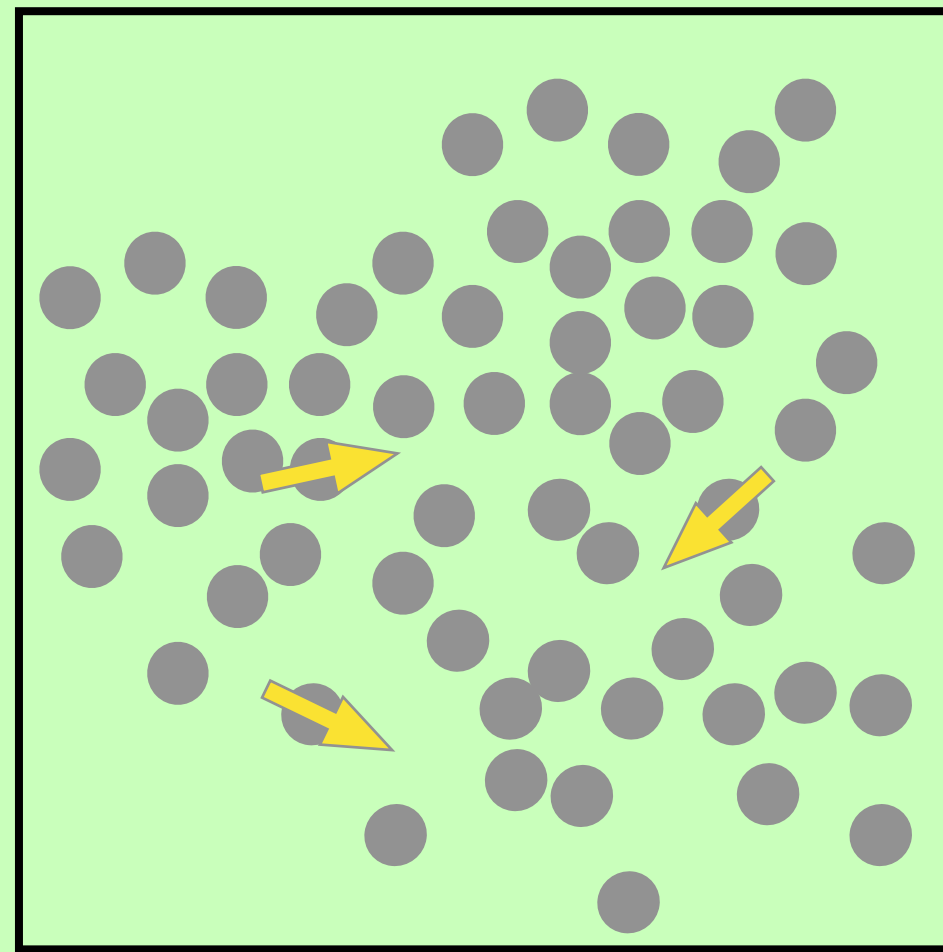
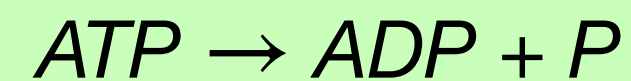
Time-reversal symmetry breaking (TRSB)

How much retrievable at different scales?

Nonreciprocal → Symmetry breaking

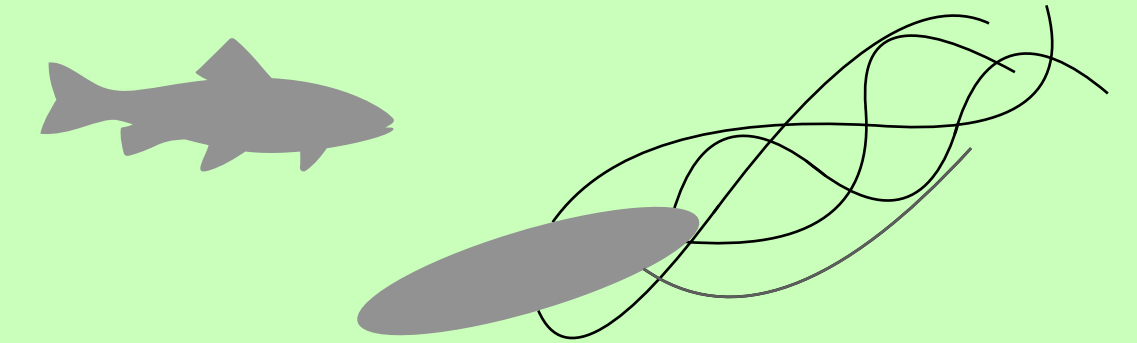
Emergence of macroscopic irreversible behaviour from microscale energy-input

Energy input at
particle scale
($\text{\AA} - \mu\text{m}$)



[Tan et al. Nature 2022]

Emergent
irreversible behaviour
at macro scale
($\mu\text{m} - \text{mm} - \text{m}$)

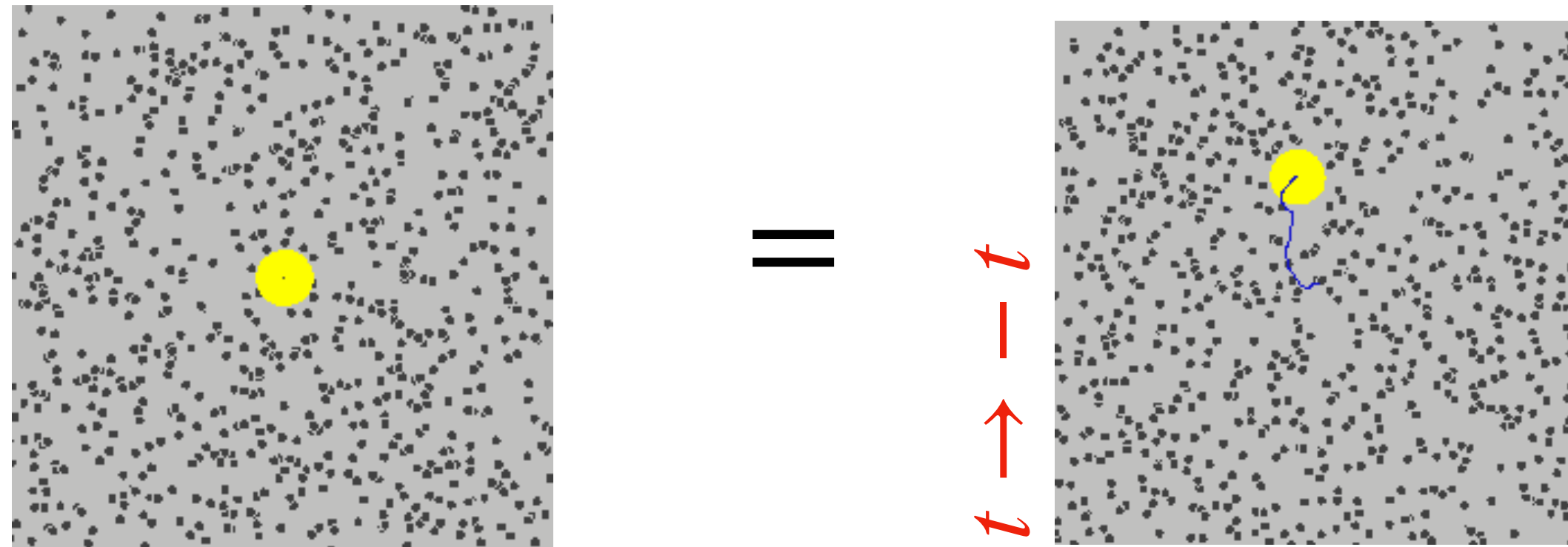


Time-reversal symmetry breaking (TRSB)

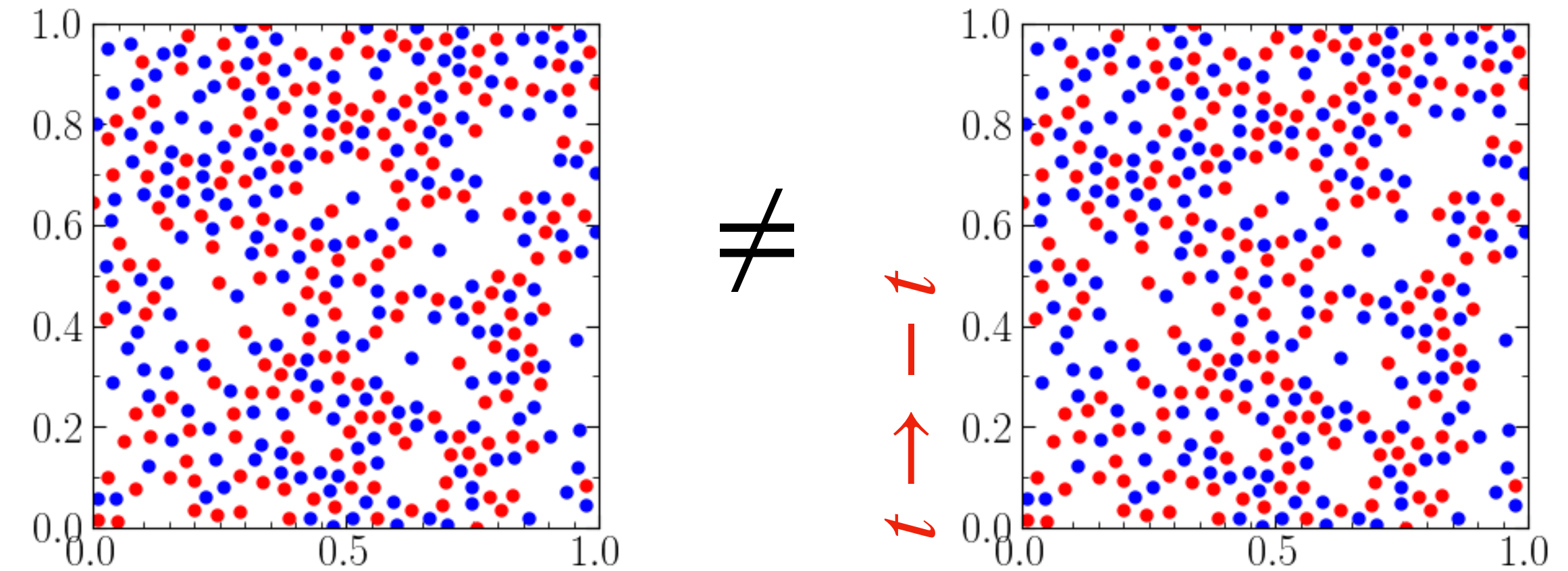
How much retrievable at different scales?

Quantifying non-equilibriumness via TRSB

EQUILIBRIUM: Time-reversal symmetry (TRS)



NON-EQUILIBRIUM: TRSB, “Arrow of Time”



• **(Informatic) entropy production rate¹:**

$$\dot{S} = \left\langle \ln \left(\frac{\mathbb{P}[X(t) |_0^T]}{\mathbb{P}_R[X_R(t) |_0^T]} \right) T^{-1} \right\rangle$$

$$\Rightarrow \langle e^{-S} \rangle = 1 \Rightarrow \langle S \rangle \geq 0$$

FTs, Generalisation of Second Law

$$\dot{S} = \int \frac{d\omega}{2\pi D} [\omega \tilde{C}(\omega) - 2D \tilde{R}'(\omega)]$$

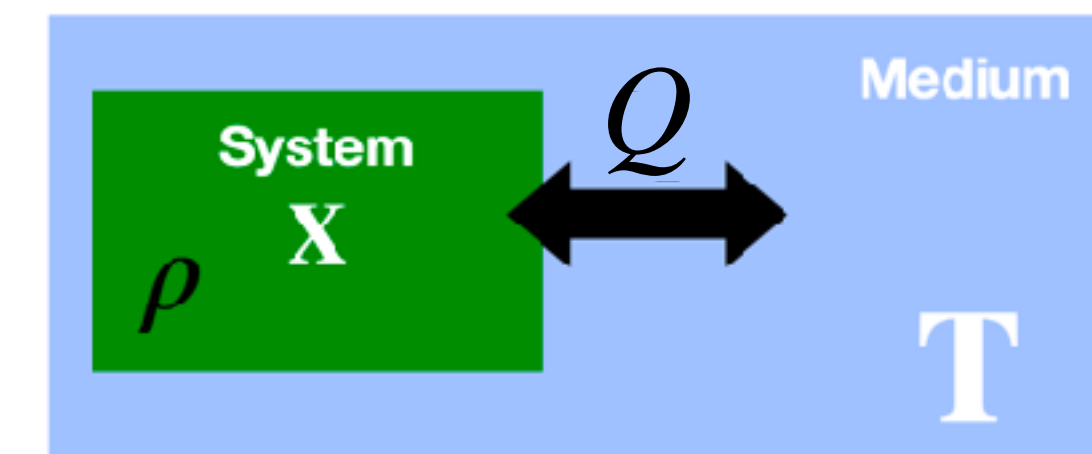
$X(t) |_0^T$: random trajectory $t \in [0, T]$, $X \in \mathbb{R}^d$

$\mathbb{P}$: path probability, Onsager – Machlup

$\mathbb{P}_R, X_R$: time reversal, $t \rightarrow -t$

Heat dissipation

$$\dot{S} = \dot{S}_{\text{shannon}} - \frac{\dot{Q}}{T} \geq 0$$



Quantifies TRSB = Distance from Equilibrium

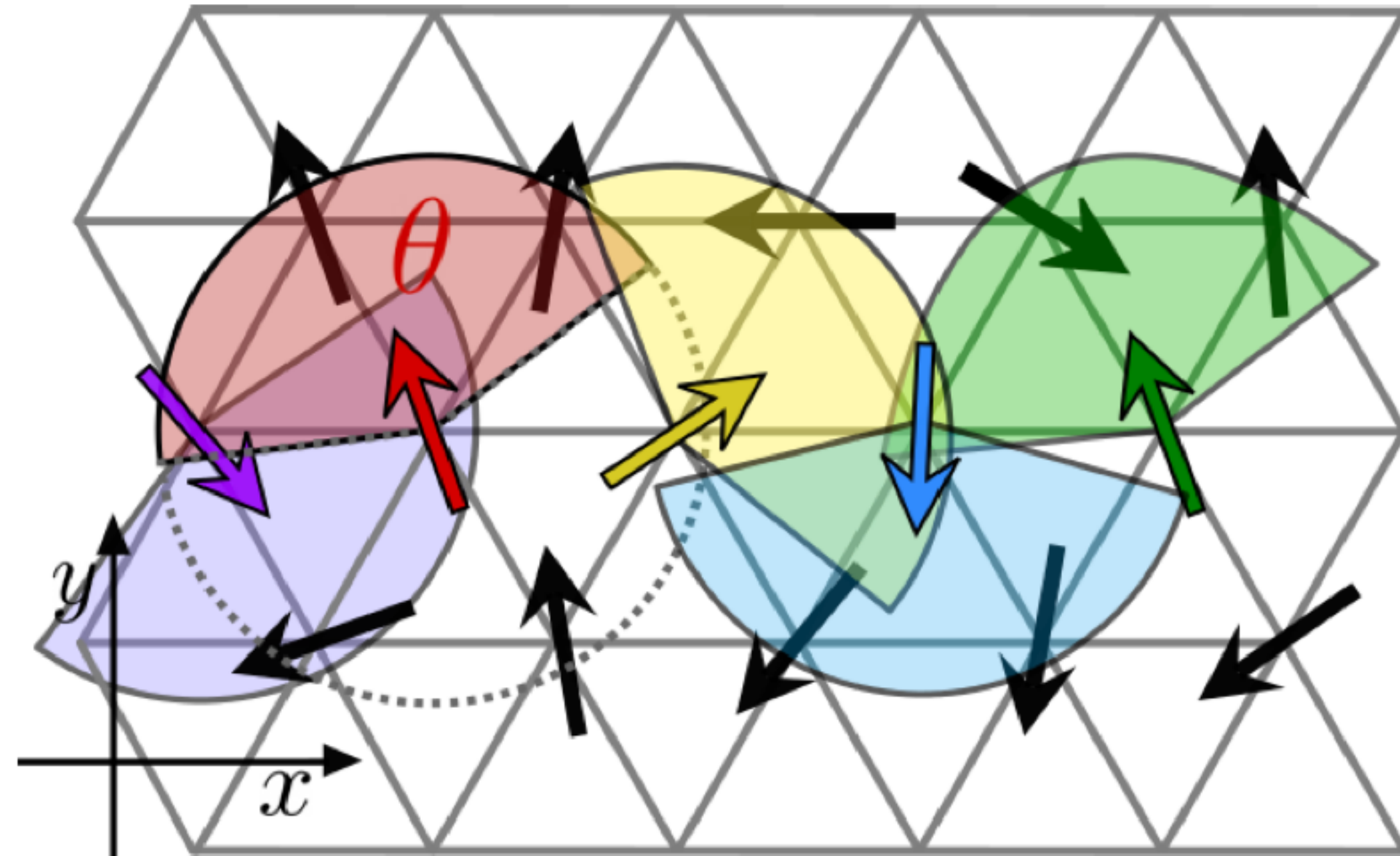
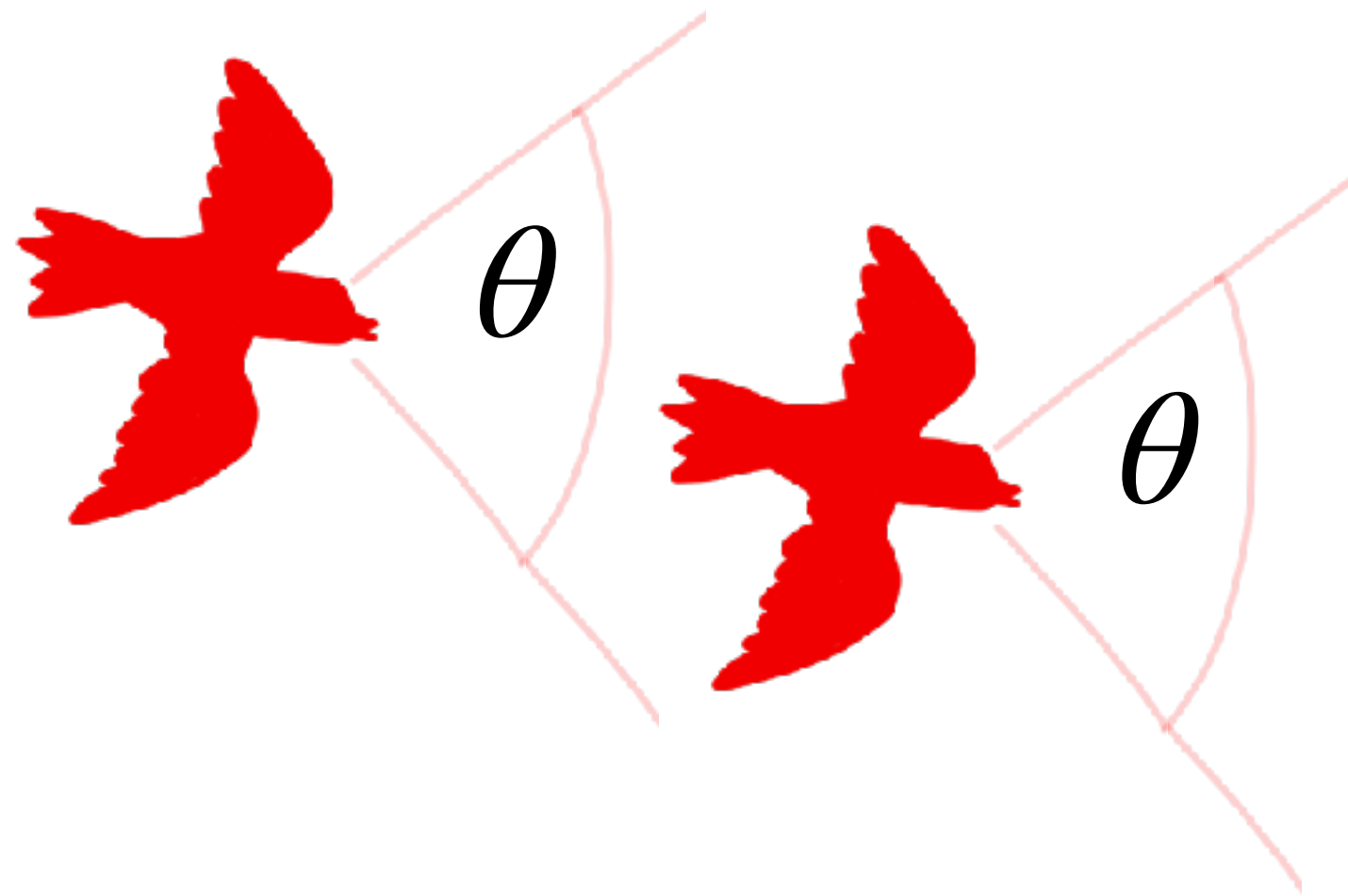
[Seifert, PRL (2005)], [Sekimoto, Springer (2010)]

[Harada, Sasa, PRL (2005)]

Outline

- (I) NR spin systems with vision-cone interactions
- (II) Two-Species system with inter-species nonreciprocity
- (III) Spin systems with quenched, NR disorder

One-species NR system with vision cones



SL, Klapp & Martynec, Phys. Rev. Lett. (2023)

$$E_i = - \sum_{\langle i,j \rangle} J_{ij}(\phi_i) \cos(\phi_i - \phi_j),$$

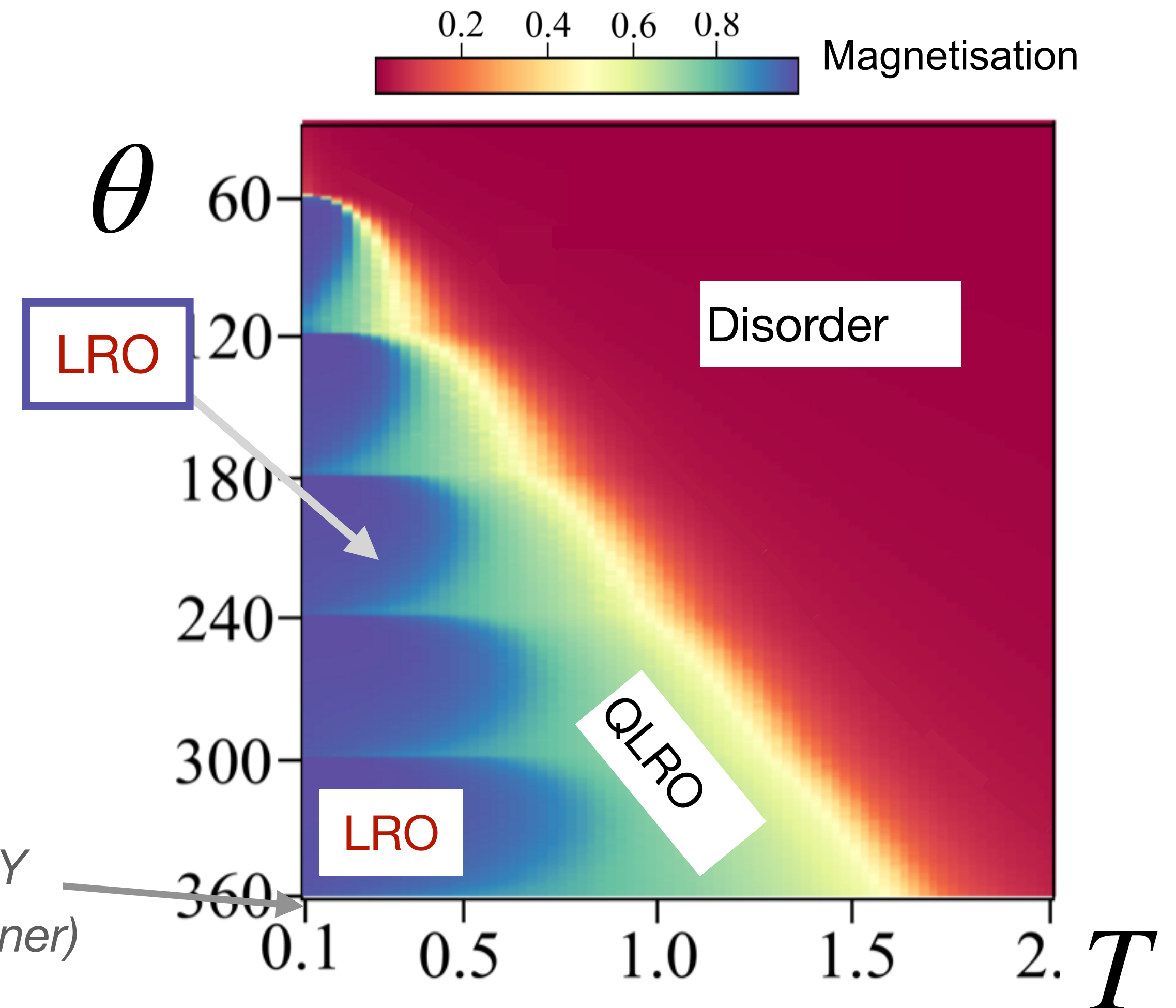
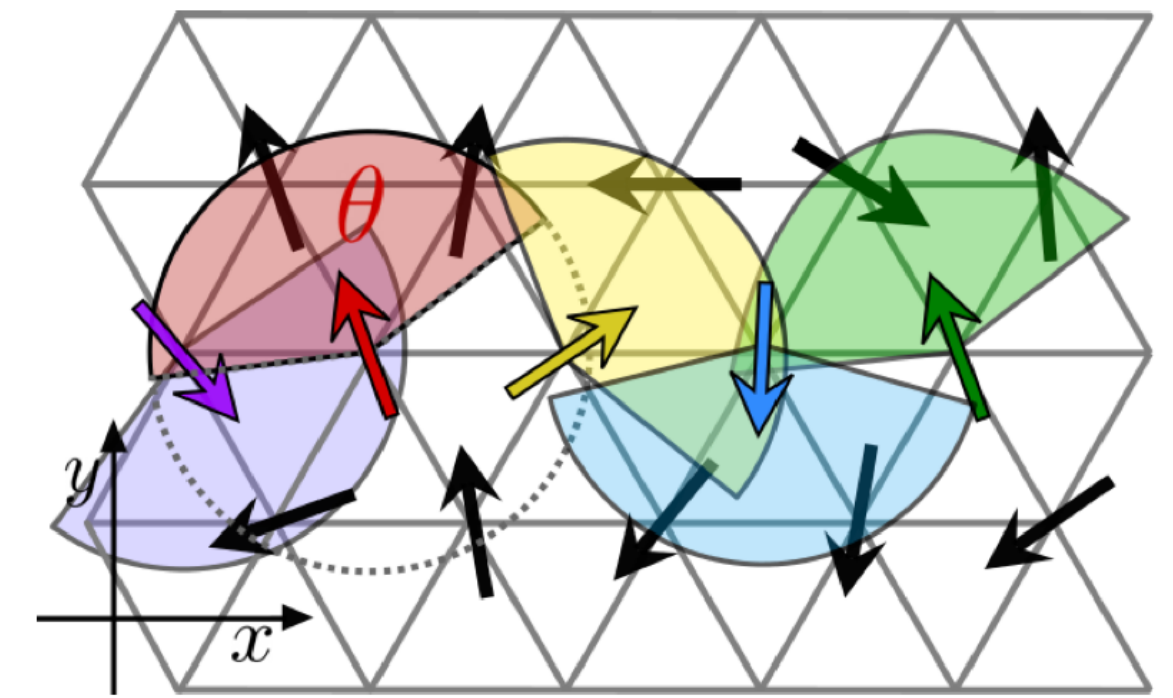
$$J_{ij}(\phi_i) = \begin{cases} J, & \text{if } j \text{ is "within vision cone"} \\ 0, & \text{else} \end{cases}$$

- **Local Energy function:** a move that lowers E_i can increase E_j
- 2D, nearest-neighbour coupling
- Monte Carlo (Glauber)

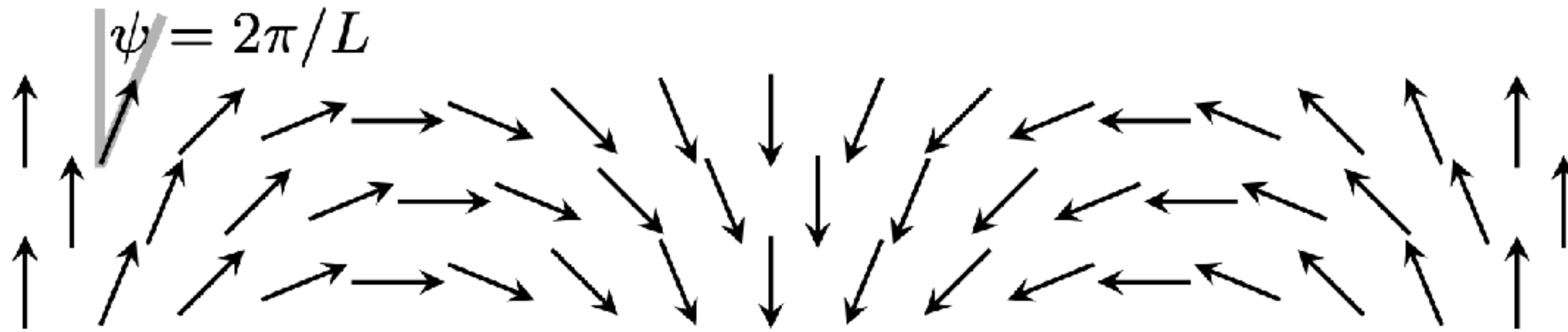
$$w(\phi_i \rightarrow \phi'_i) = [1 - \tanh([E_i(\phi_i) - E_i(\phi'_i)]/2T)] / 2$$

Emergence of true **Long-range order (LRO)** !

Vision-Cone XY model



Suppression of spin waves



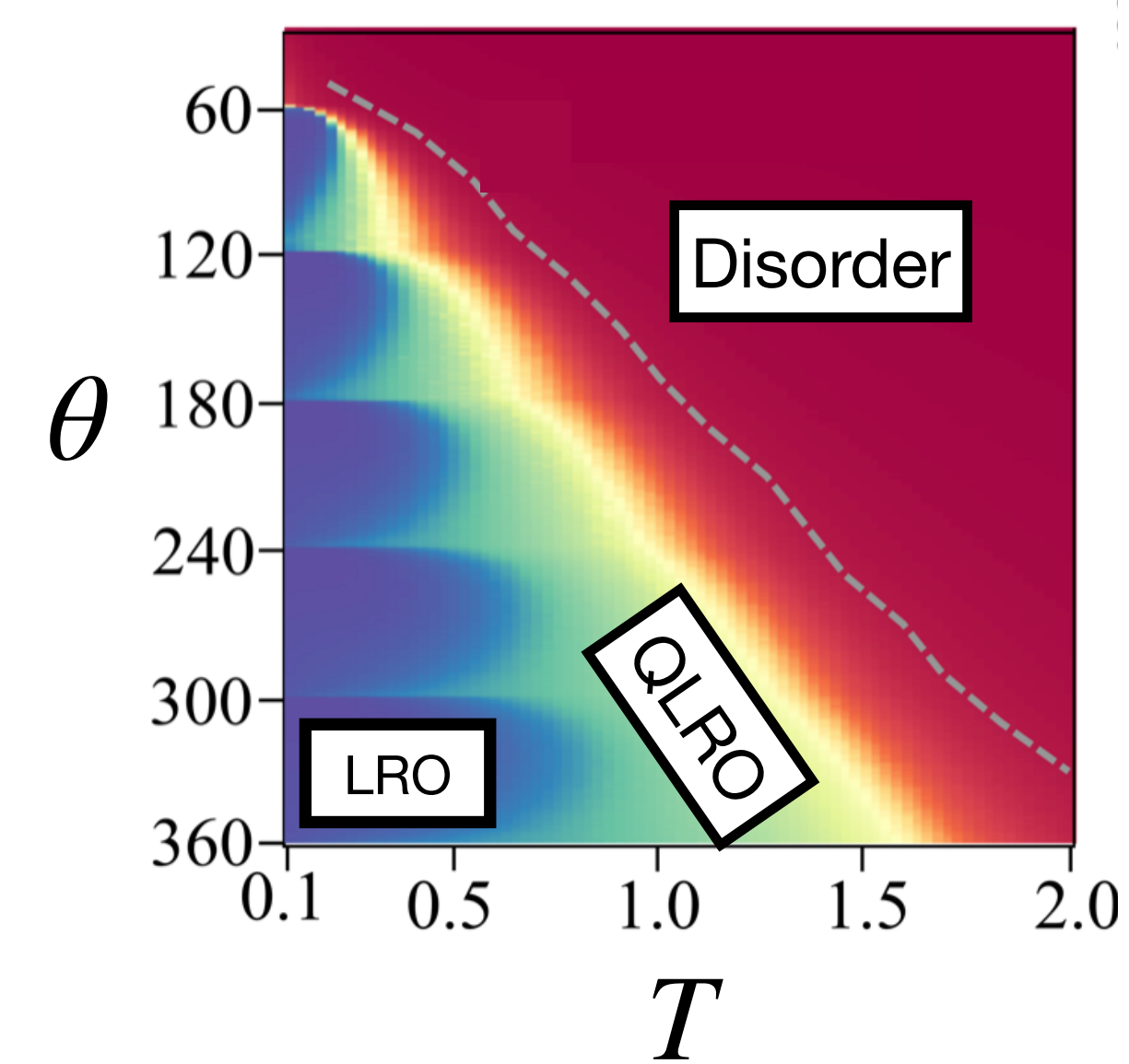
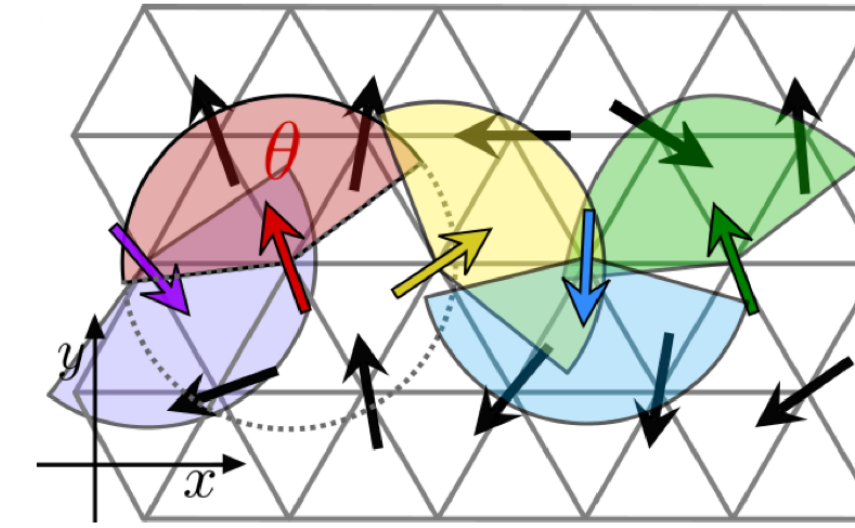
Cost to rotate spin i :

$$E_i \sim \psi^2 = 4\pi^2/L^2 + J$$

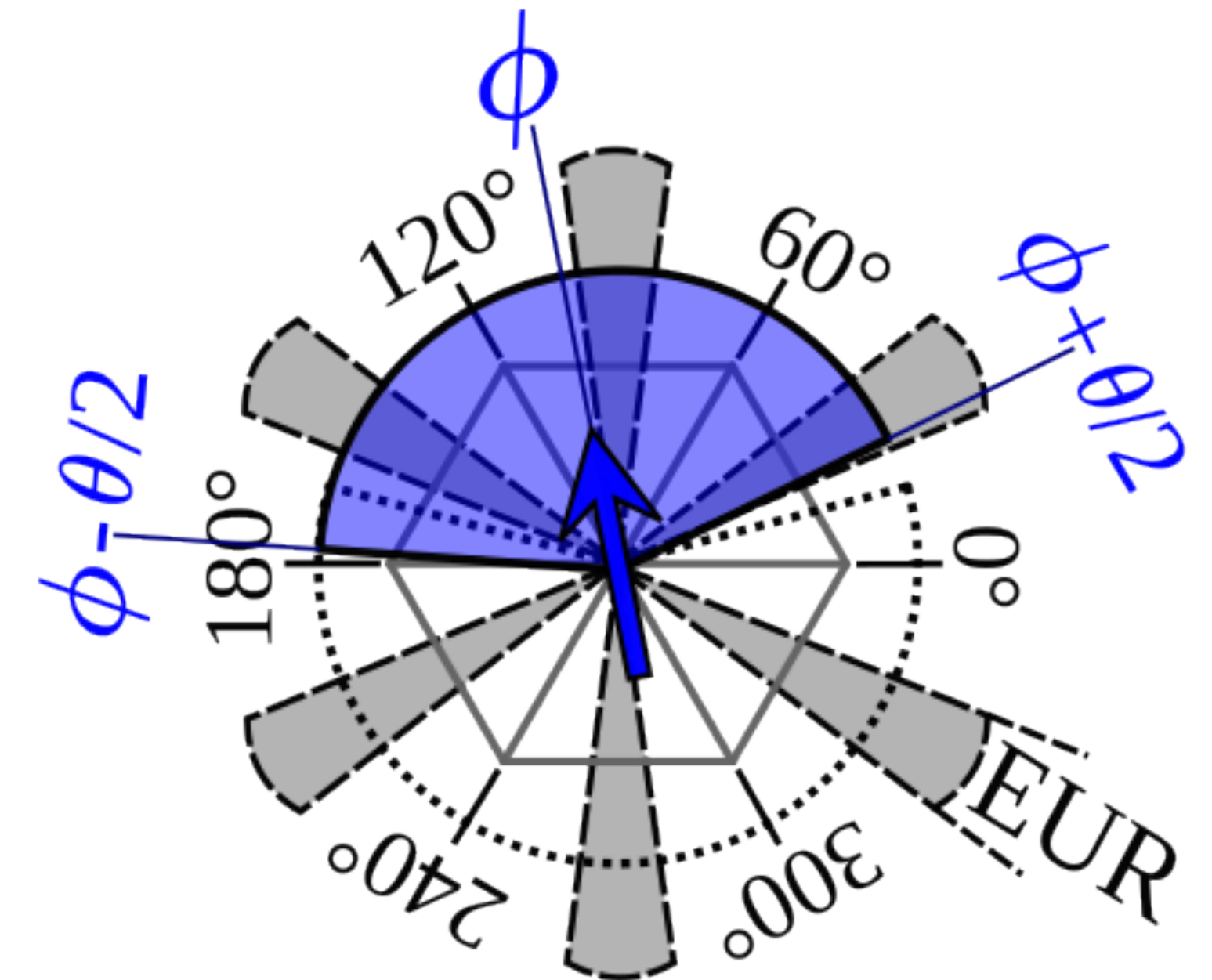
Total cost to excite spin wave:

$$E_{\text{total}} = 4\pi^2 + L^2 \eta J,$$

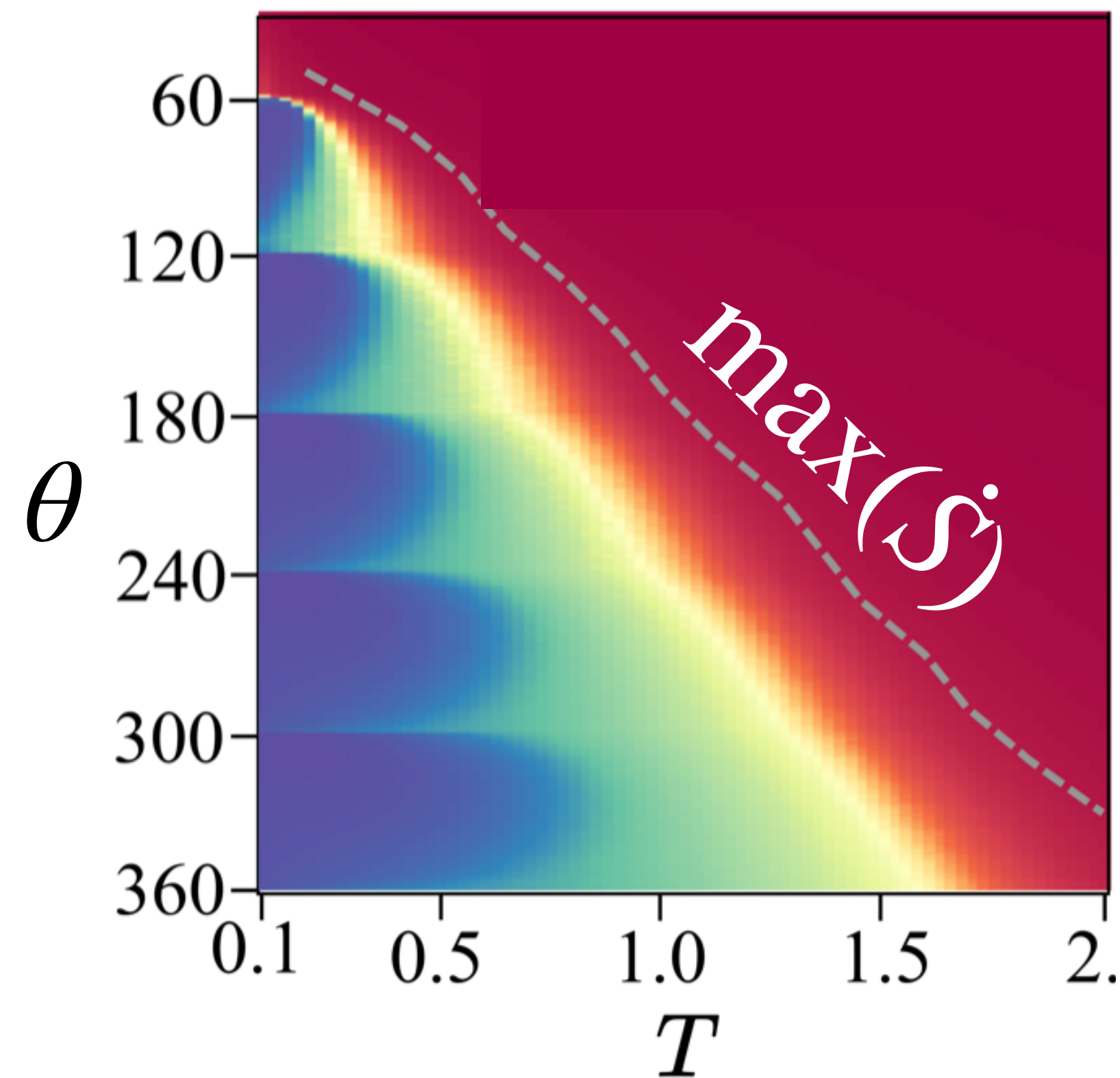
$$\eta \rightarrow 1 - \frac{[\theta \bmod (360^\circ/n)]}{(360^\circ/n)}$$



- “Energetically unfavourable regions” (EUR)
- Suppression of spin waves due to EUR $\rightarrow$ LRO
- No LRO at $\theta = n \cdot 60^\circ$



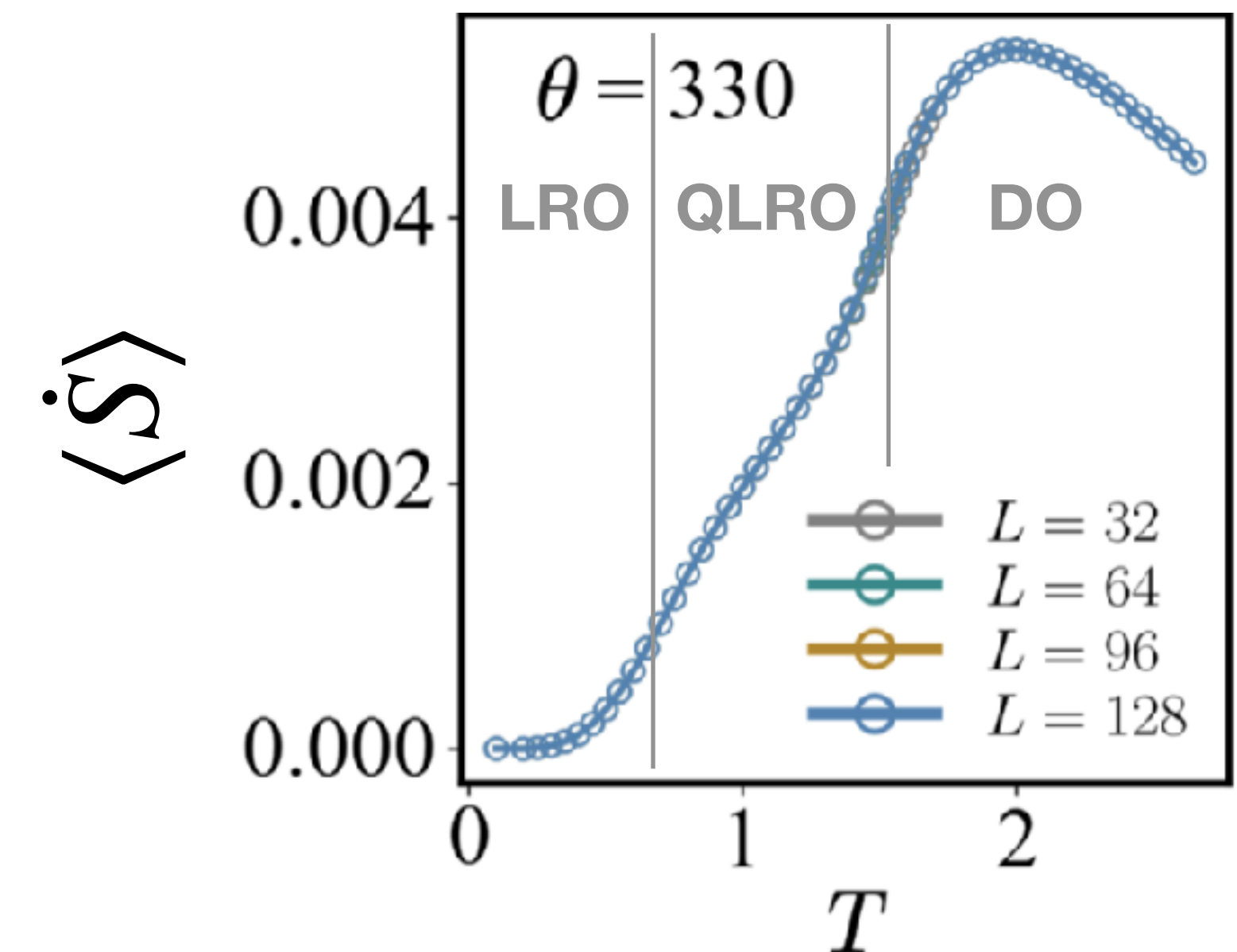
A nonreciprocal XY model — Arrow of time?



Net entropy production:

$$\langle \dot{S} \rangle = \left\langle \ln \frac{w[\phi_i(t) \rightarrow \phi_i(t + dt)]}{w[\phi_i(t + dt) \rightarrow \phi_i(t)]} \right\rangle / dt$$

↓
Unidirectional
propagation
of defect lines
↓



Unidirectional propagation of defects

→ Manifestation of nonequilibrium-ness in **collective fluctuations**

A nonreciprocal XY model — devil in the details

PHYSICAL REVIEW LETTERS **135**, 088303 (2025)

Ordering and Defect Cloaking in Nonreciprocal Lattice XY Models

Pankaj Popli^{1,*}, Ananyo Maitra^{2,3,†} and Sriram Ramaswamy^{4,5,‡}

¹*Centre for Condensed Matter Theory, Department of Physics, Indian Institute of Science, Bengaluru 560 012, India*

²*Laboratoire de Physique Théorique et Modélisation, CNRS UMR 8089, CY Cergy Paris Université, F-95032 Cergy-Pontoise Cedex, France*

³*Laboratoire Jean Perrin, Sorbonne Université and CNRS, F-75005 Paris, France*

⁴*Centre for Condensed Matter Theory, Department of Physics, Indian Institute of Science, Bengaluru 560 012, India*

⁵*International Centre for Theoretical Sciences, Tata Institute of Fundamental Research, Bangalore 560 089, India*

PHYSICAL REVIEW LETTERS **135**, 088302 (2025)

Inescapable Anisotropy of Nonreciprocal XY Models

Dawid Dopierala^{1,2,*}, Hugues Chaté^{3,4,2,†}, Xia-qing Shi^{5,‡} and Alexandre Solon^{2,§}

¹*TCM Group, Cavendish Laboratory, JJ Thomson Avenue, Cambridge CB3 0HE, United Kingdom*

²*Sorbonne Université, CNRS, Laboratoire de Physique Théorique de la Matière Condensée, 75005 Paris, France*

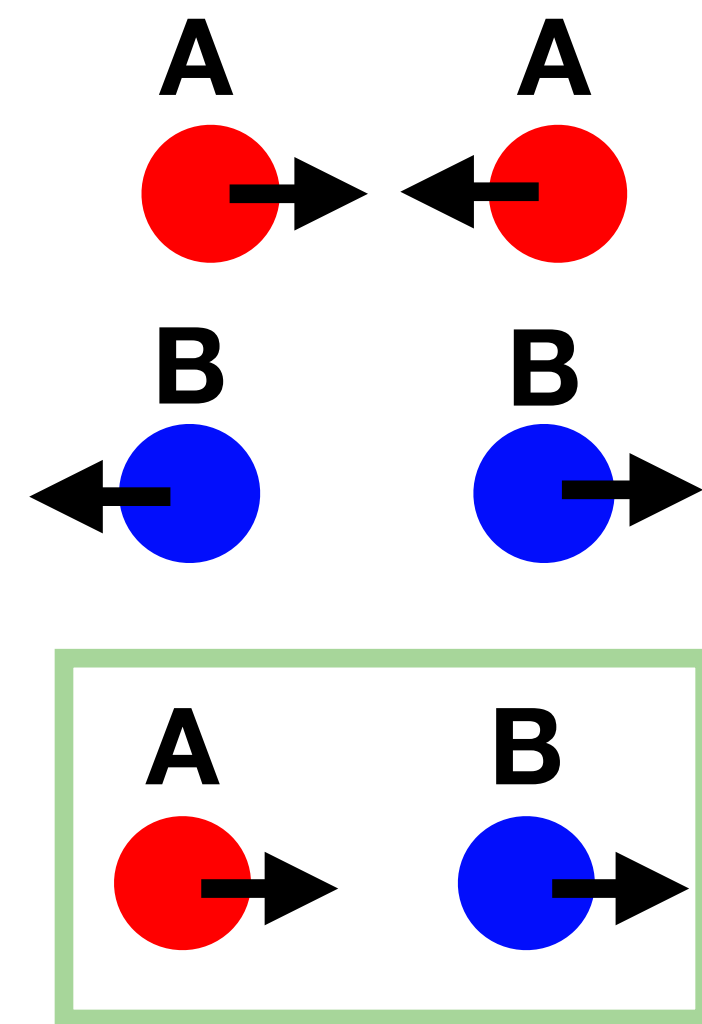
³*Service de Physique de l'Etat Condensé, CEA, CNRS Université Paris-Saclay, CEA-Saclay, 91191 Gif-sur-Yvette, France*

⁴*Computational Science Research Center, Beijing 100094, China*

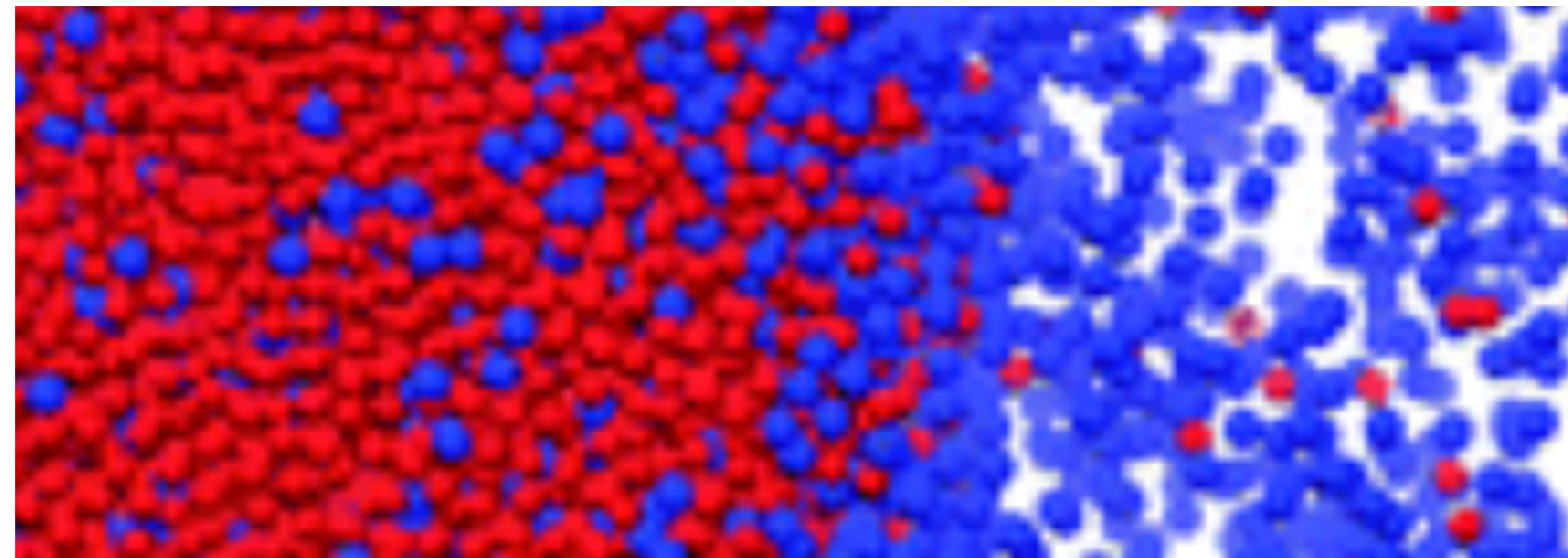
⁵*Center for Soft Condensed Matter Physics and Interdisciplinary Research & School of Physical Science and Technology, Soochow University, Suzhou 215006, China*

- Z. Liu, et al. "Dynamic approach to the two-dimensional nonreciprocal XY model with vision cone interactions." PRE (2025).
- Y. Du, et al. "Collective behavior of the non-reciprocal vision cone XY lattice-gas model." Newton (2025).
- L. Dadhichi, et al. "Nonmutual torques and the unimportance of motility for long-range order in two-dimensional flocks." PRE (2020).
- Y-B. Sh, et al. "General Hamiltonian description of nonreciprocal interactions." arXiv preprint arXiv:2505.05246 (2025).
- ...

Two-Species system with inter-species nonreciprocity



Binary fluids with NR interactions



Example: Nonreciprocal Cahn Hilliard model

$t \in \mathbb{R}, r \in \mathbb{R}^d, \phi_{i=1,2,..}(r, t) \in \mathbb{R}$ scalar fields (densities)

$$\dot{\phi}_i = - \nabla \cdot \left(- \nabla \frac{\delta F}{\delta \phi_i} + \kappa_{ij} \nabla \phi_j + \sqrt{2\epsilon} \Lambda_i \right)$$

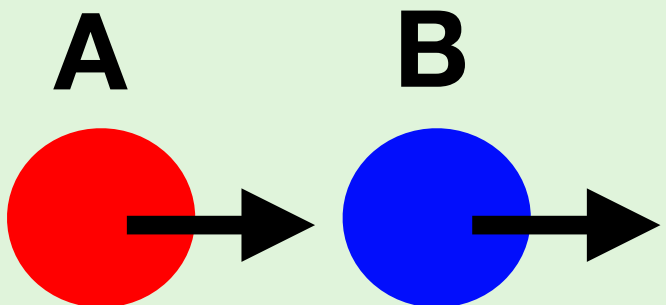
Gaussian space-time white noise

$$\langle \Lambda_i(x, t) \Lambda_j(x', t') \rangle = \delta_{ij} \delta(x - x') \delta(t' - t)$$

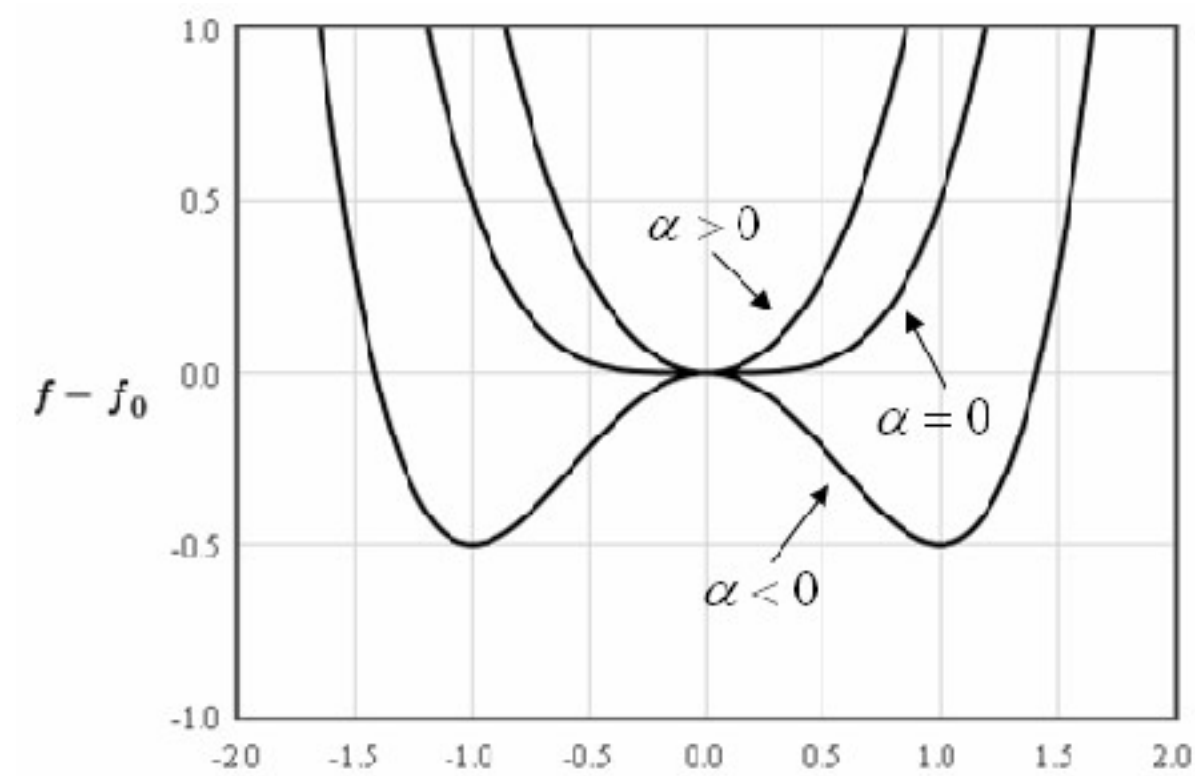
NR:

$$\kappa_{ij} \neq \kappa_{ji}$$

“Non-equilibrium term” (NR) : $\kappa_{AB} = \delta, \kappa_{BA} = -\delta$



“Equilibrium term”

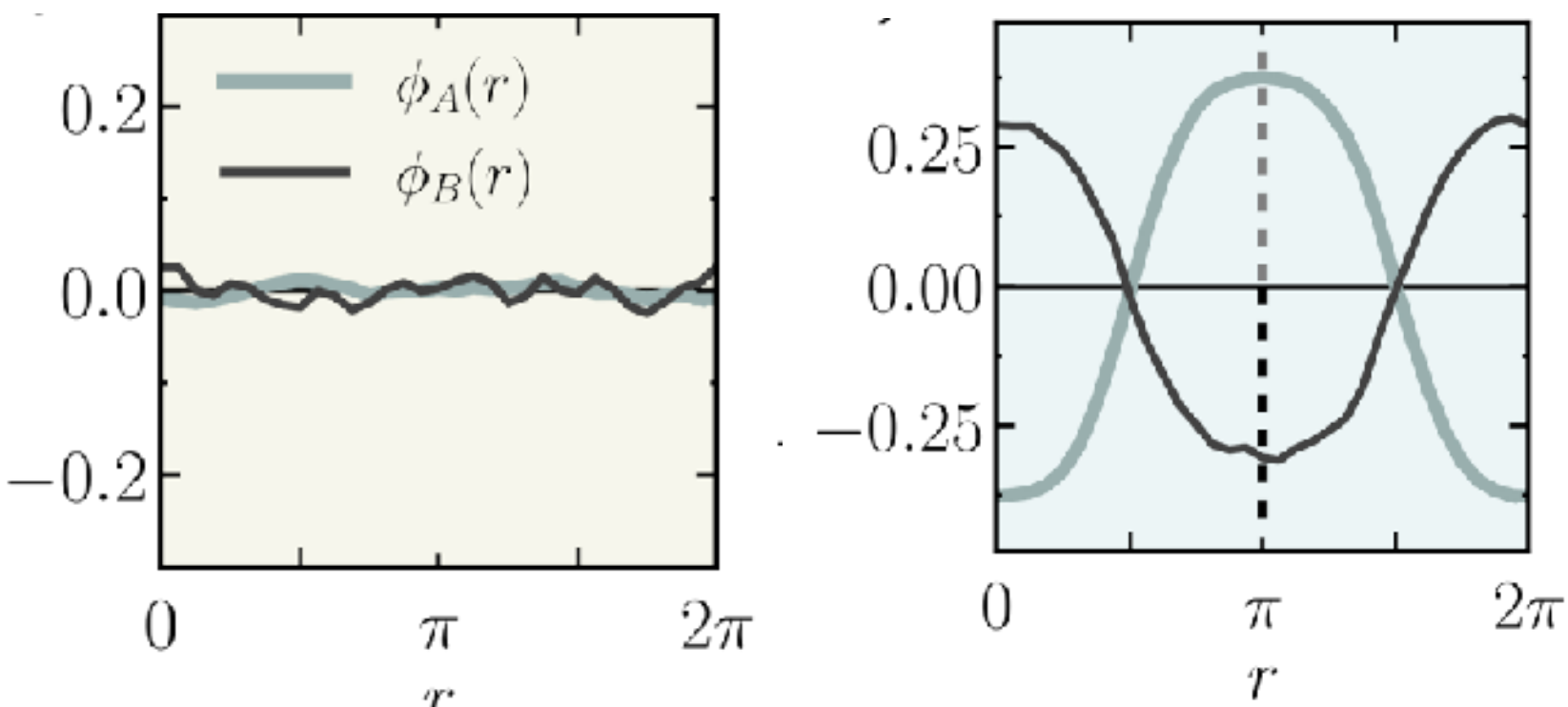


Example:

- Cahn-Hilliard (ϕ^4 theory)
- $\phi_{A,B}$ = concentrations

$$F = \frac{1}{2} \int_V dx \left[\sum_{i=A,B} \frac{1}{6} \phi_i^4 + \alpha_i \phi_i^2 + \gamma_i (\nabla \phi_i)^2 - \kappa \phi_A \phi_B \right]$$

Demixing transition

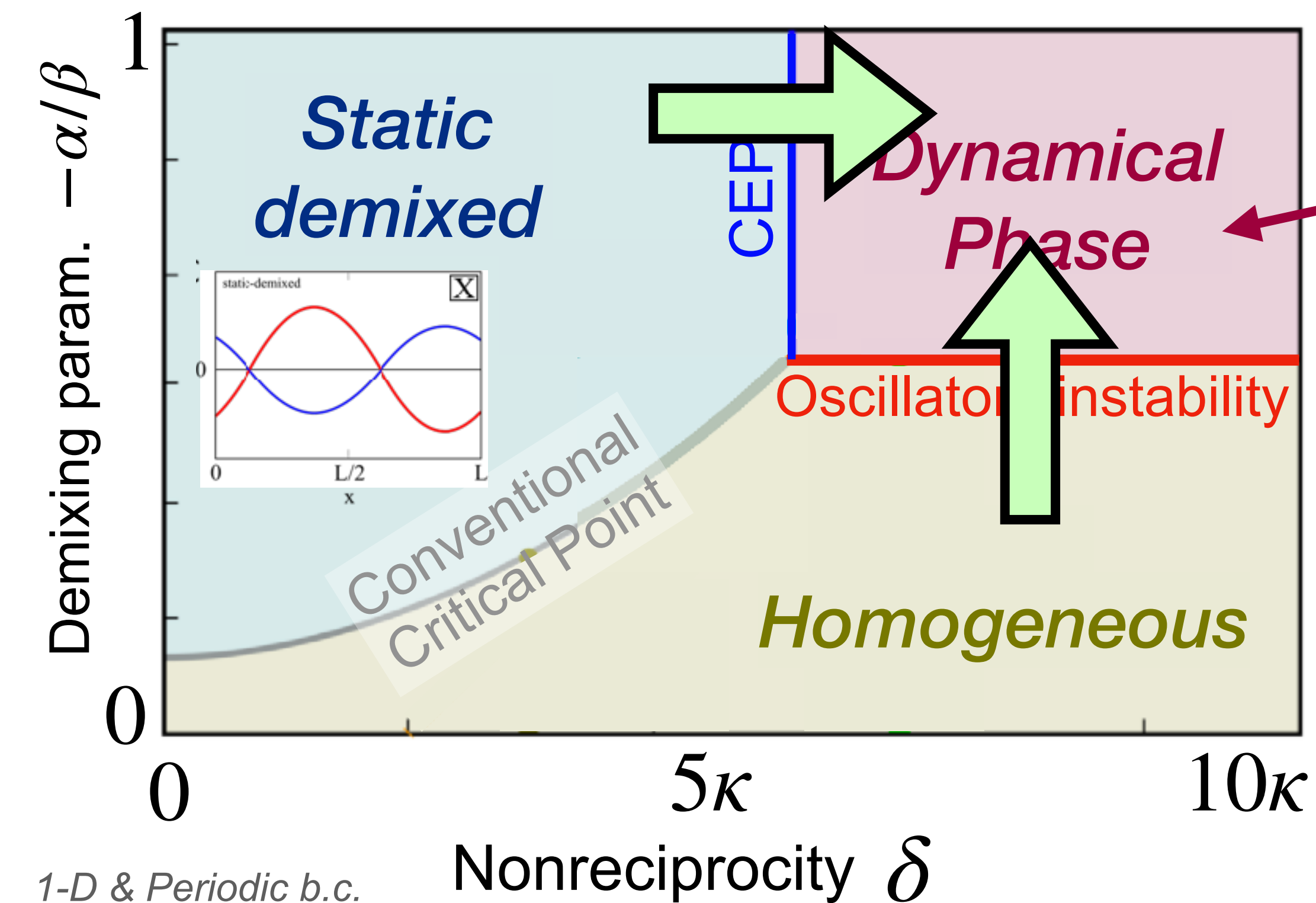


[You et al., PNAS (2020)]
 [Saha et al, PRX (2020)]
 [Mason et al., ArXiv (2024)]
 [Mandal et al., PRE (2024)]
 [Brauns et al., PRX (2025)]

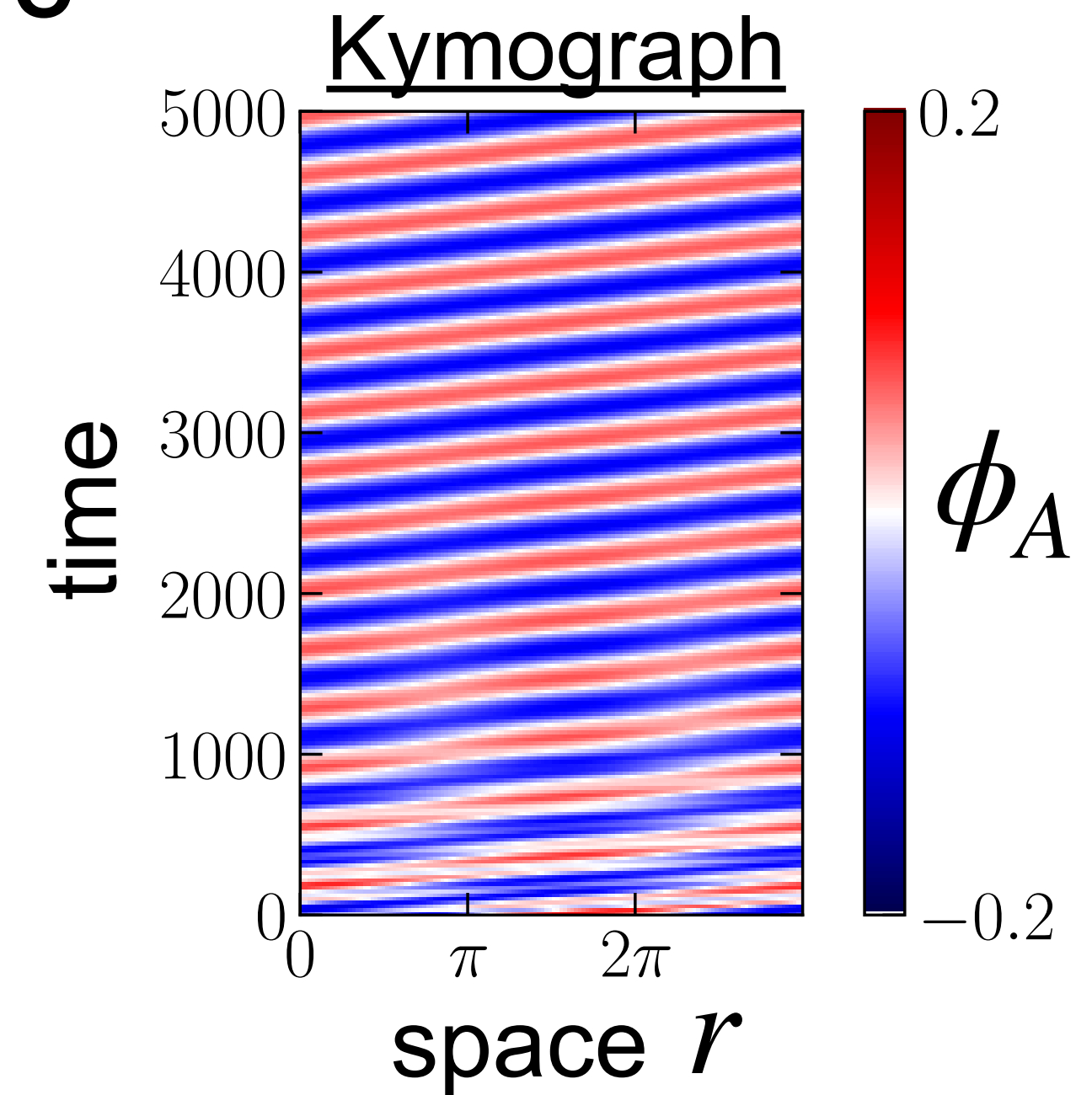
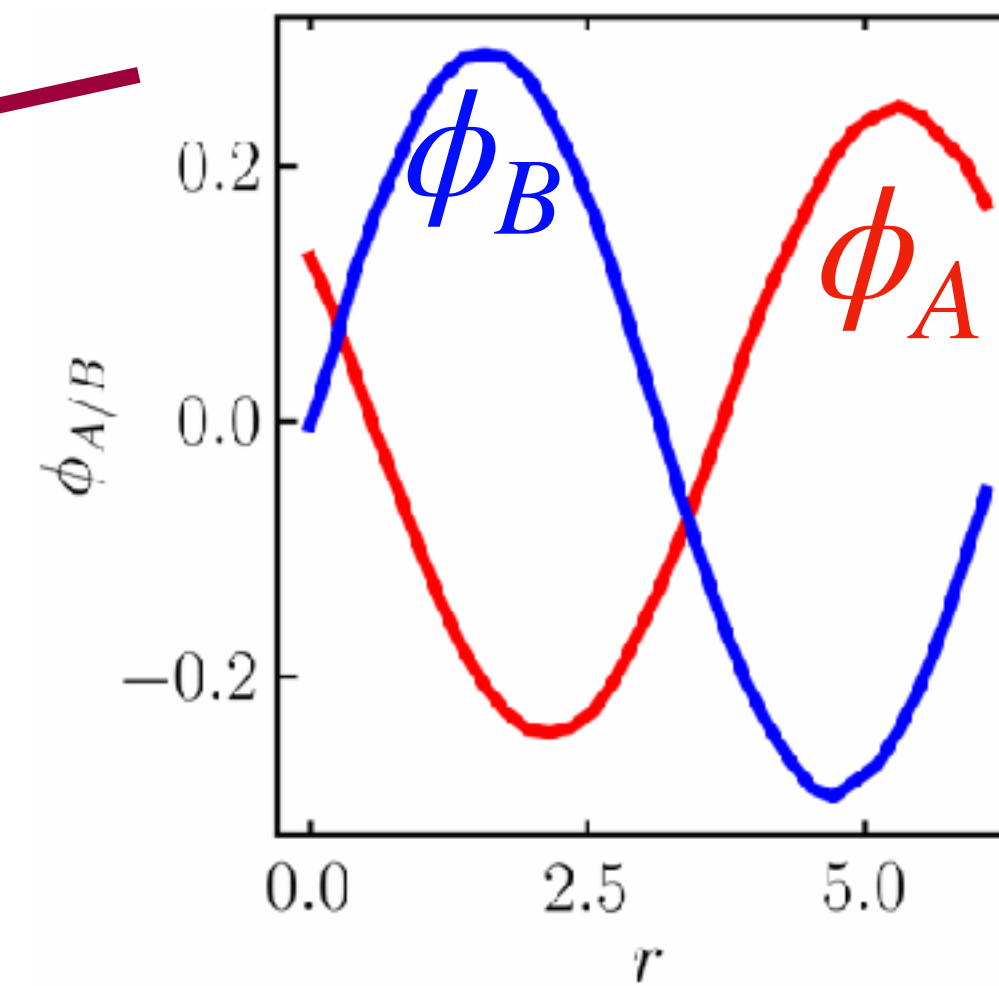
$$\dot{\phi}_A = \nabla[(\alpha + \phi_A^2 - \gamma \nabla^2) \nabla \phi_A + (1 - \delta) \nabla \phi_B] + \sqrt{2\epsilon} \Lambda_A$$

$$\dot{\phi}_B = \nabla[(\beta + \cancel{\phi_B^2} - \cancel{\gamma \nabla^2}) \nabla \phi_B + (1 + \delta) \nabla \phi_A] + \sqrt{2\epsilon} \Lambda_B$$

NR Cahn-Hilliard model



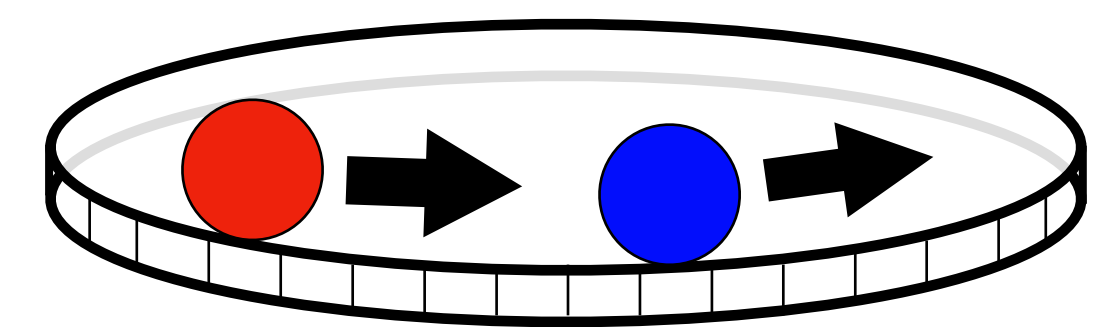
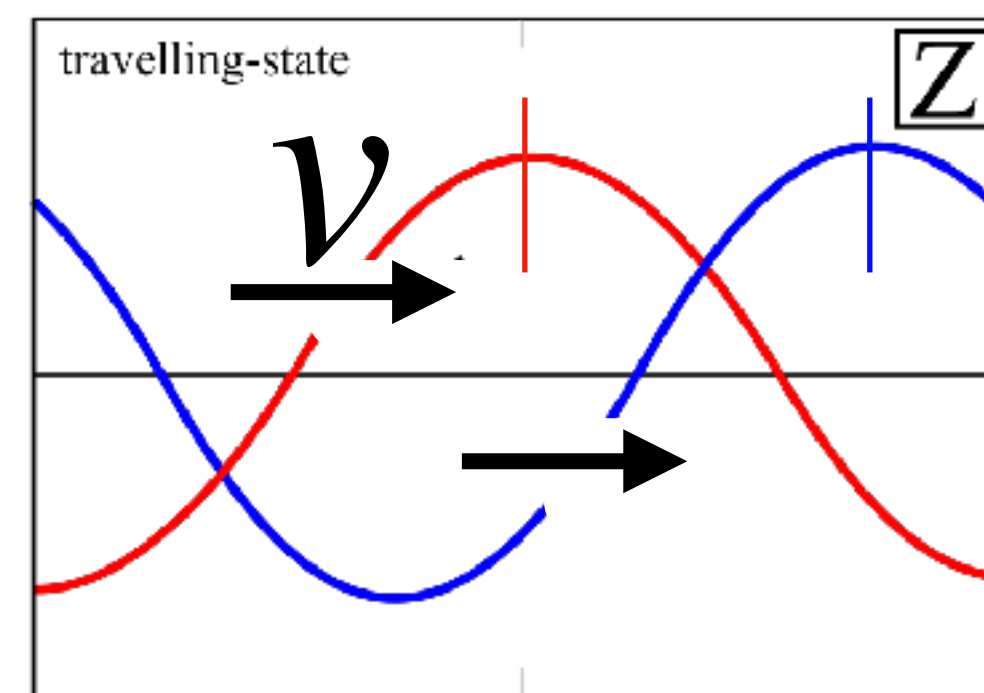
Travelling wave



Dynamical phase:
sustained mass flows

$$\Delta\theta^\pi = [\theta_A^1 - (\theta_B^1 + \pi)]$$

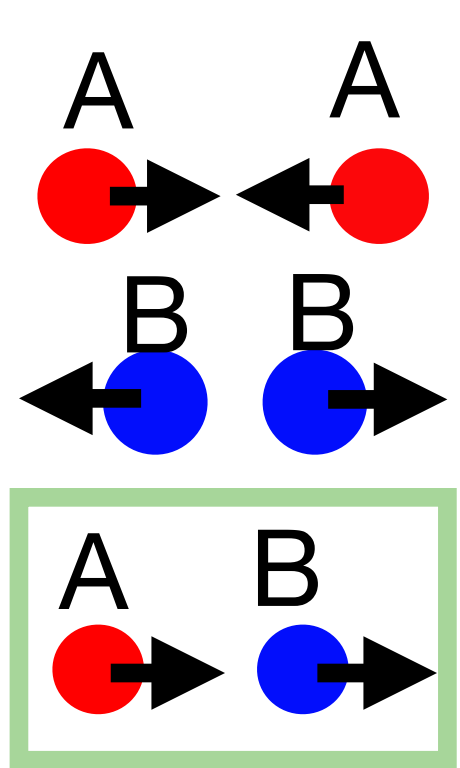
$$|v| \approx \sqrt{|\delta - \delta_c|}$$



Combination of NR +
many-body effects $\rightarrow$ net flow

[You et al, PNAS (2020)]
[Saha et al, PRX (2020)]
[Brauns et al, PRX (2024)]

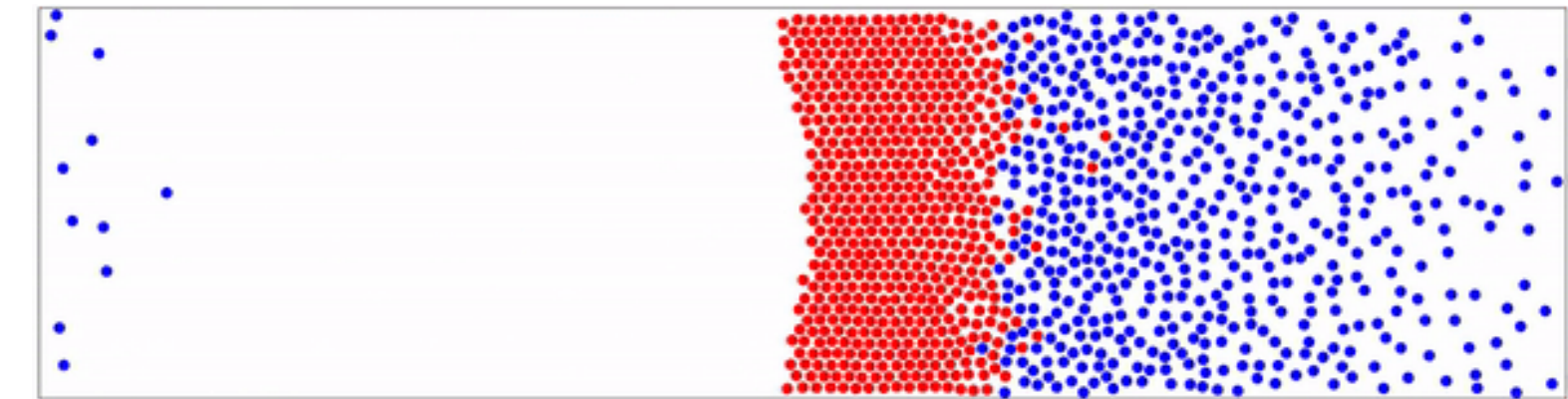
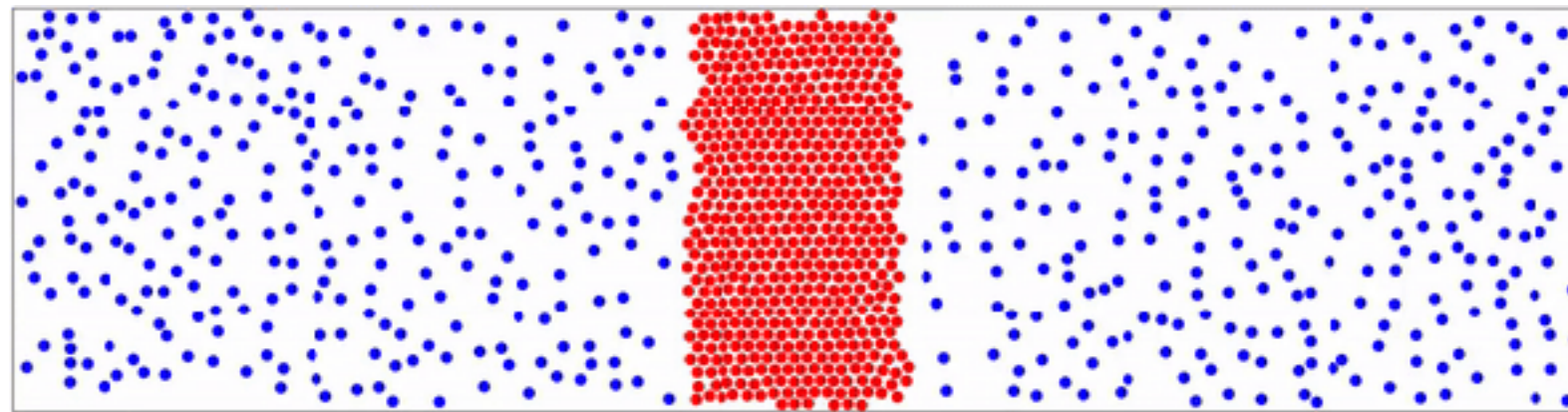
Emergence of TRSB at macroscopic scale



Static phase

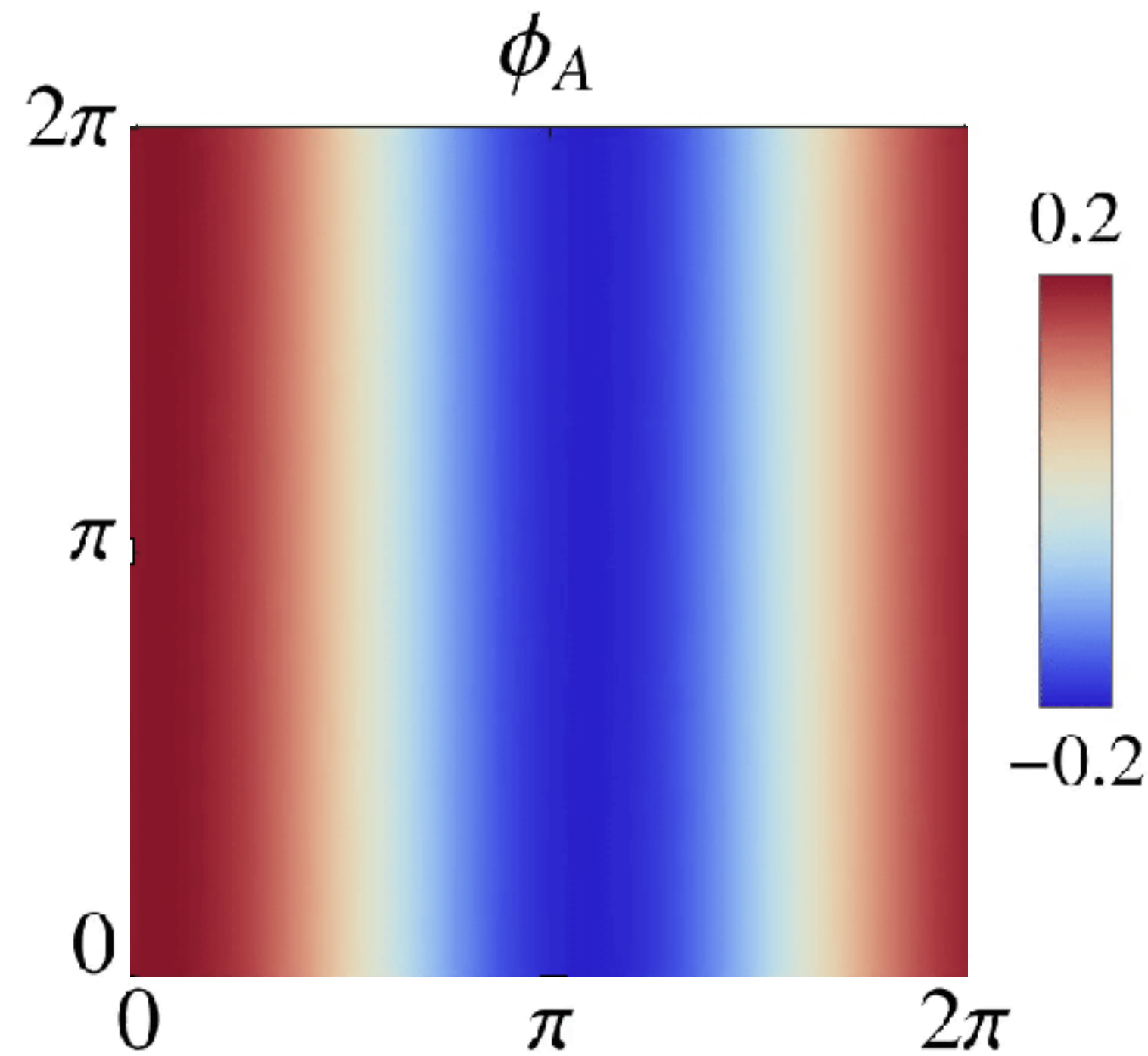
Increase Nonreciprocity

Dynamic phase

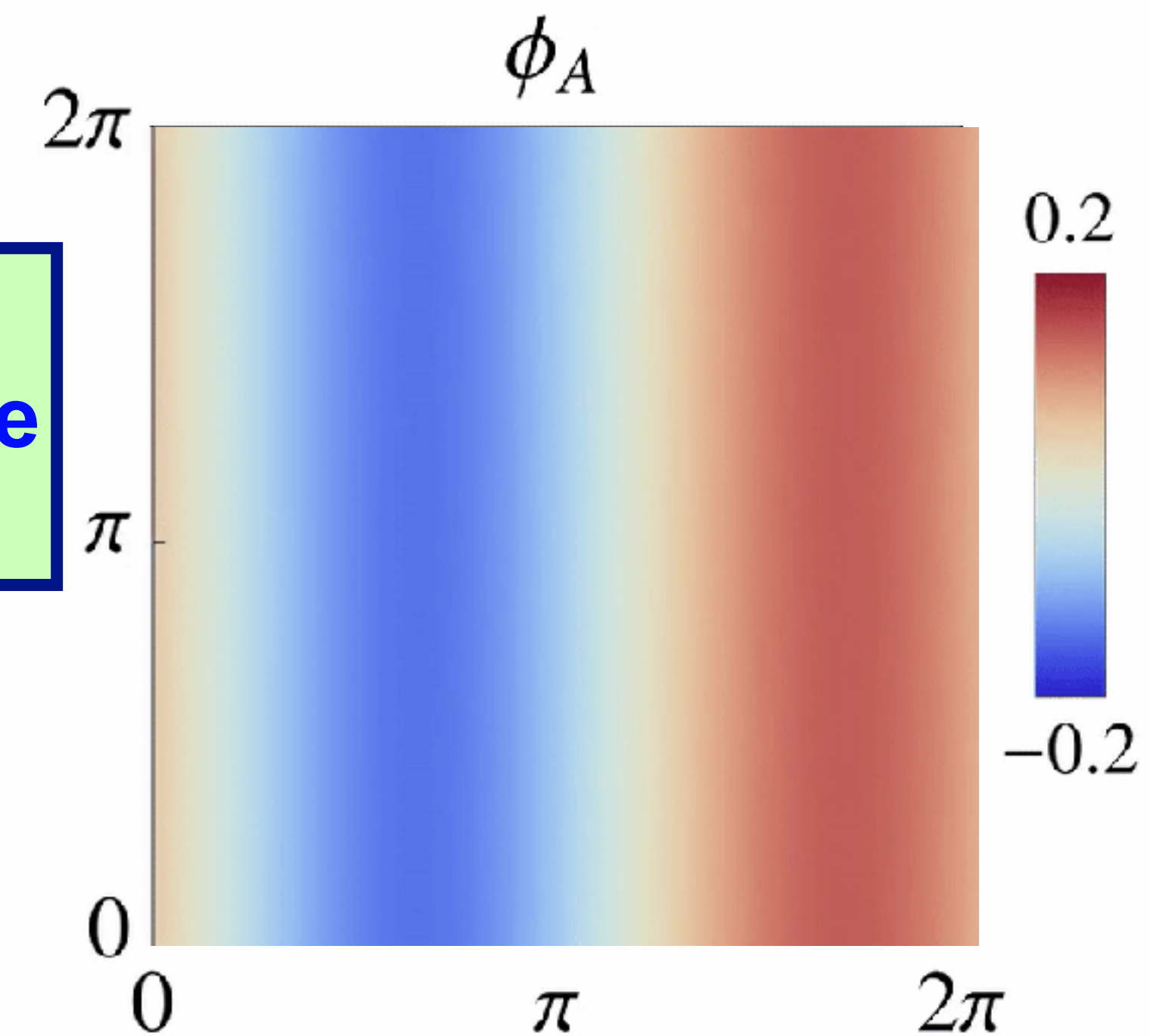


Critical
threshold

Emergence of
TRSB at collective
level?



“Equilibrium-like” phase



Clearly nonequilibrium

Entropy production in travelling wave phase

NRCH

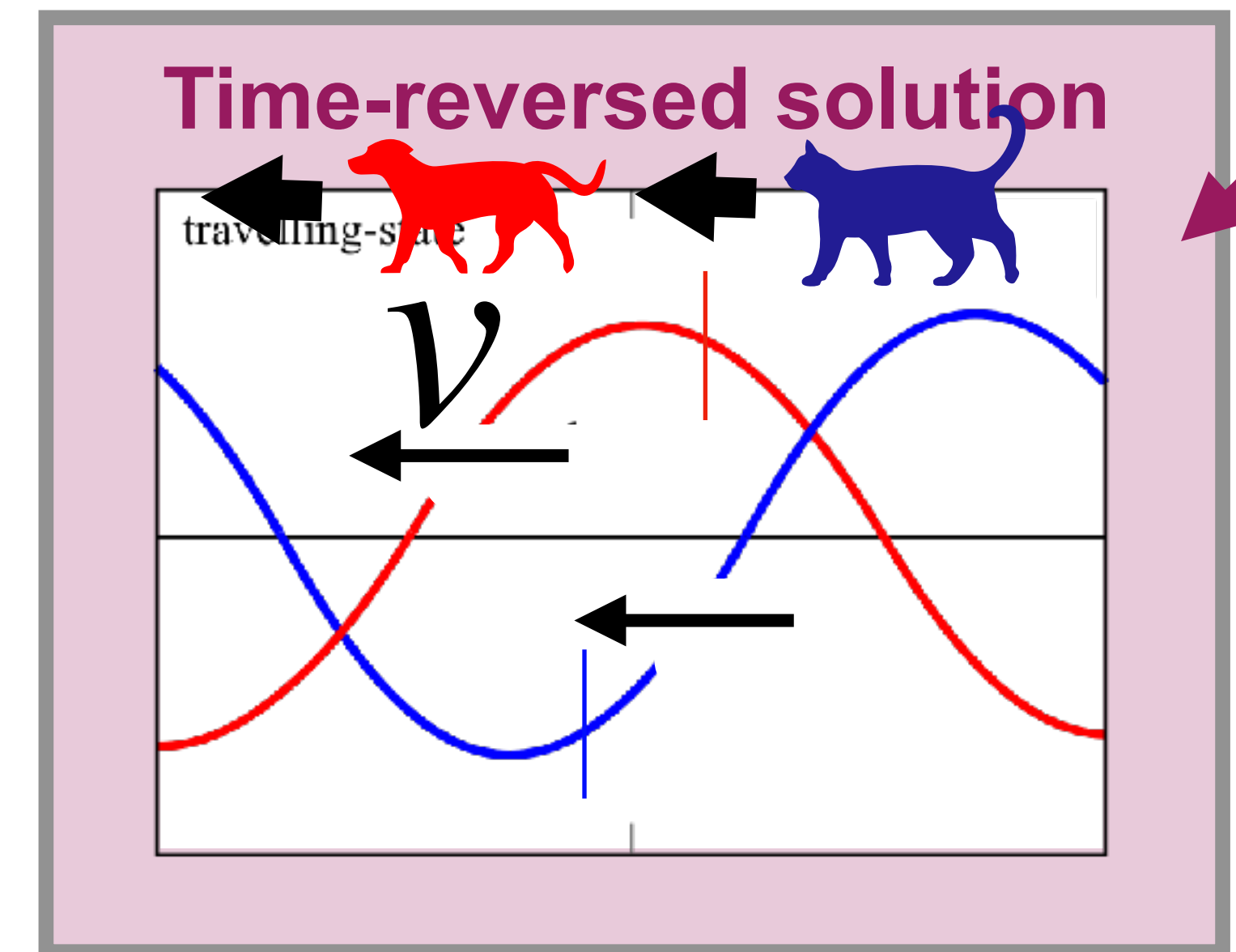
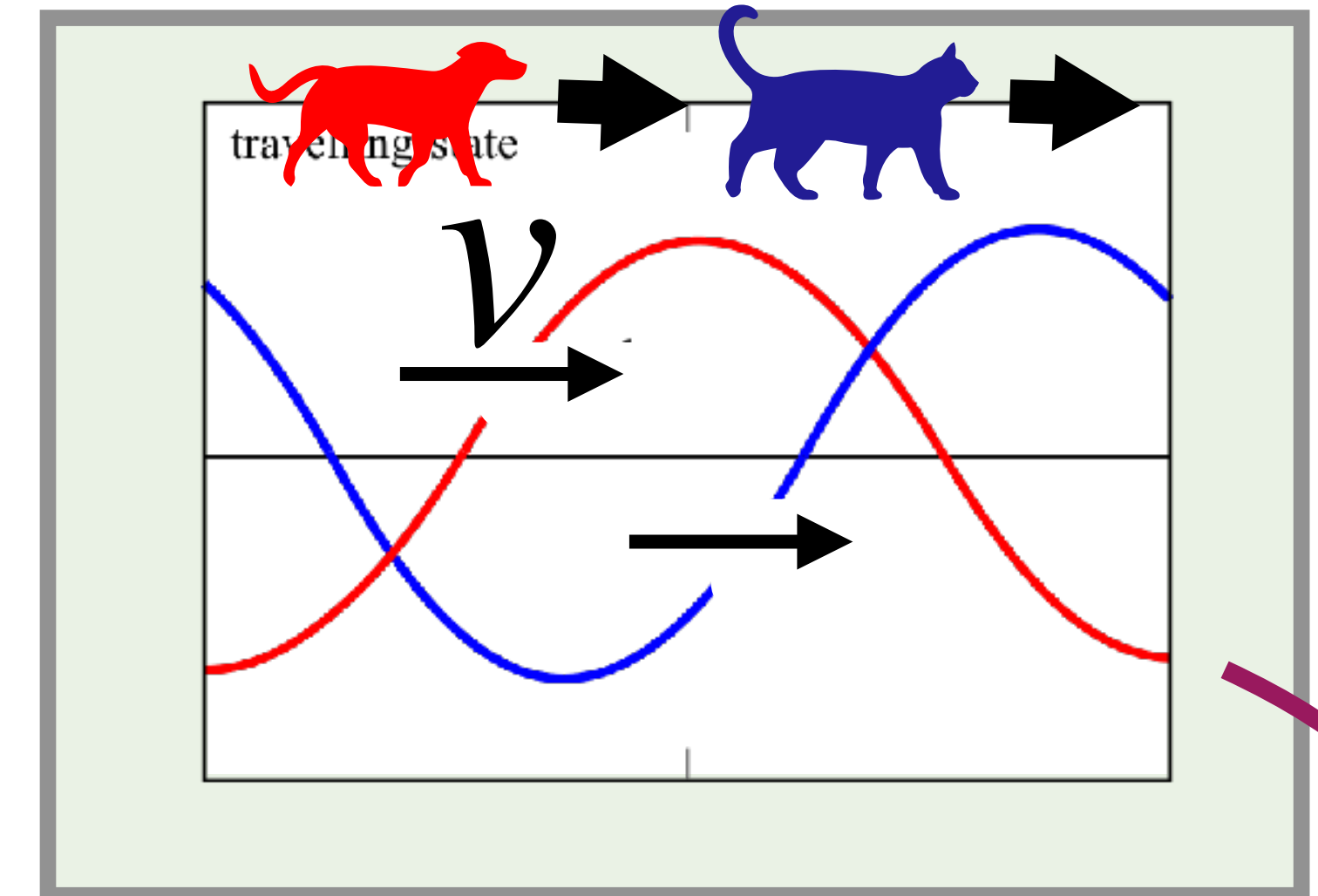
In Dynamical Phase

$$\dot{S} = \left\langle \ln \frac{P[\{\phi_{A,B}(t')\}_0^t]}{\hat{P}[\{\hat{\phi}_{A,B}(t')\}_0^t]} / t \right\rangle \sim \frac{1}{\epsilon}$$

Diverges in
thermodynamic limit!
“Macroscopic
arrow of time”

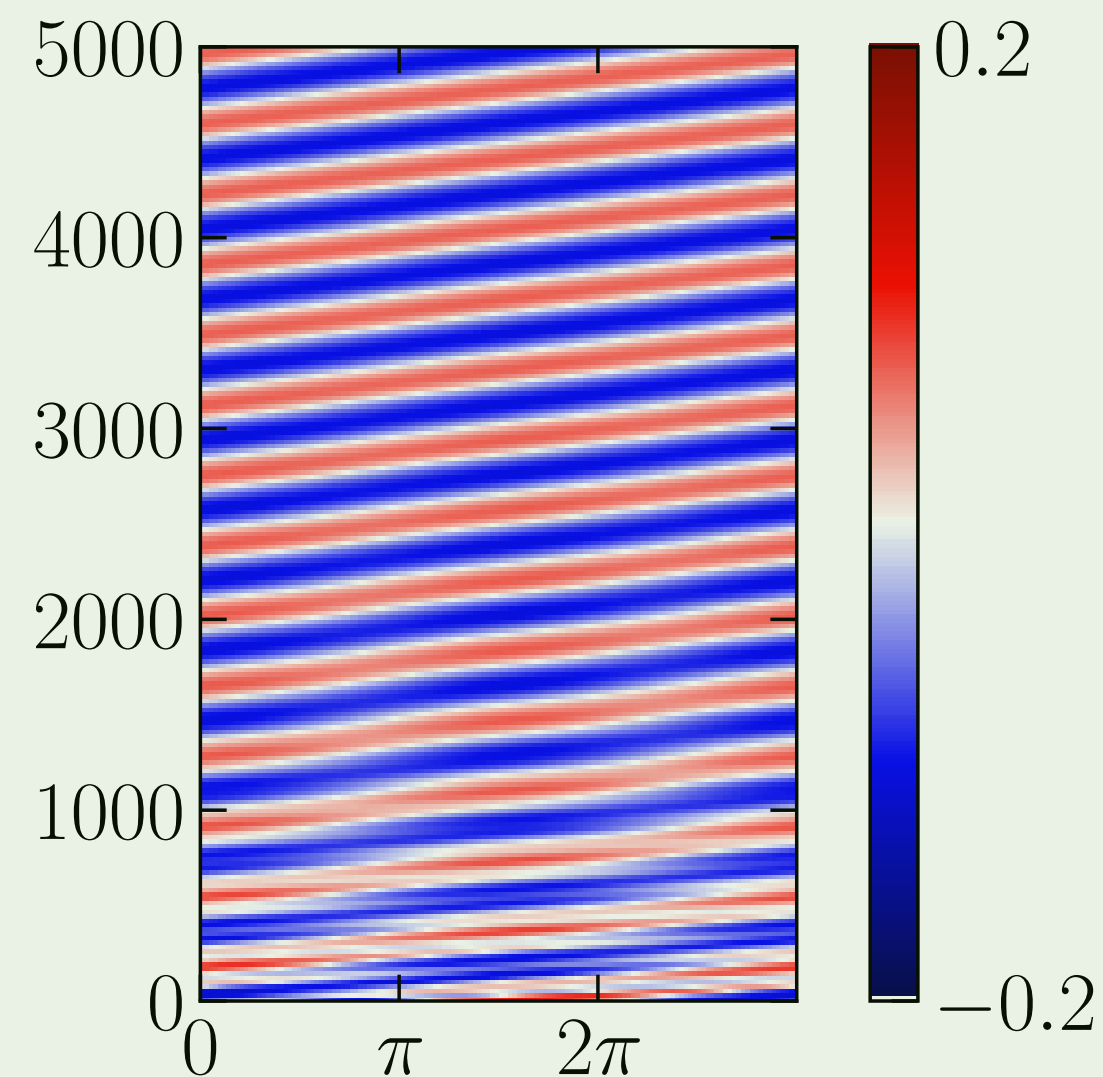
$$t \rightarrow -t$$

But sign(v) spontaneously
broken symmetry!

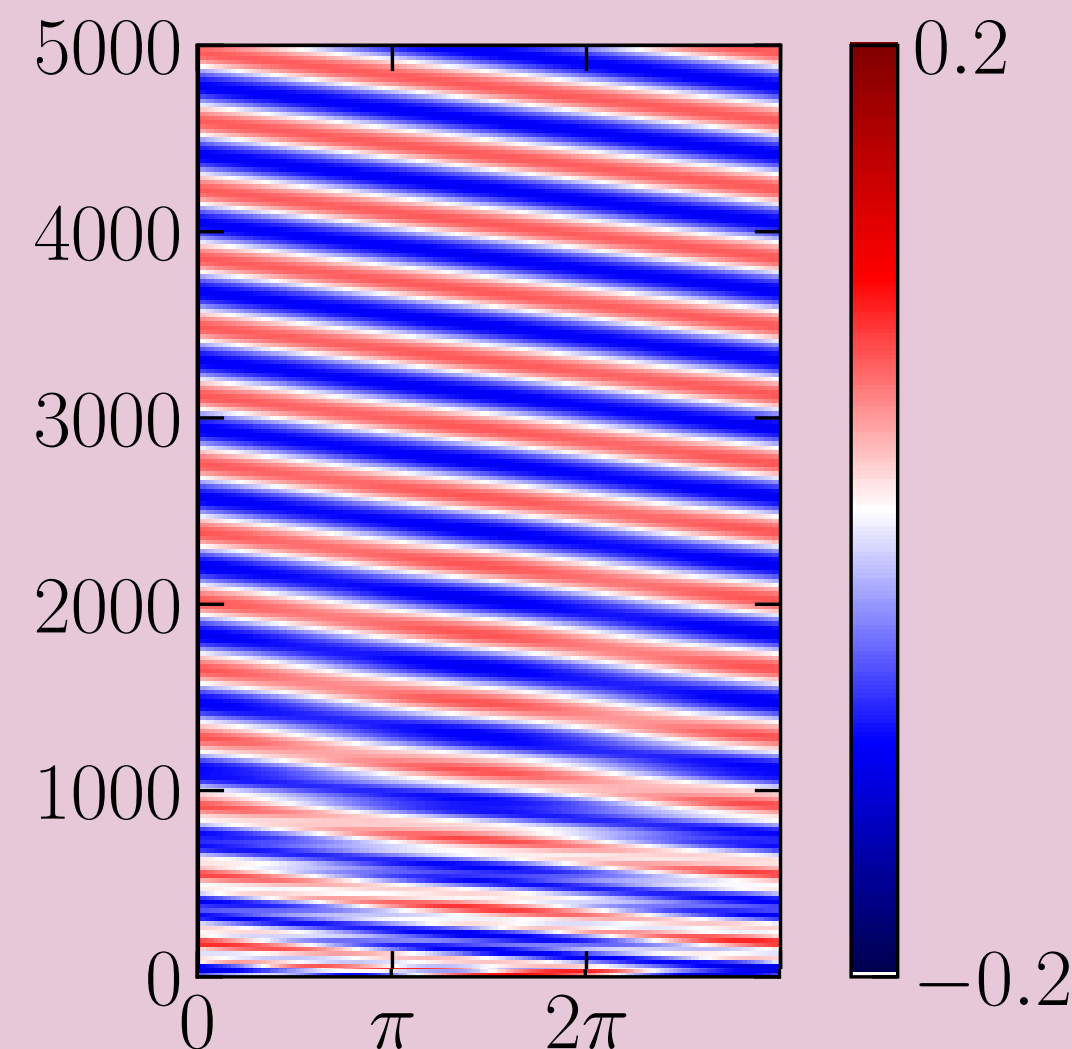


Characteristic Phase Shift

TRSB $\leftrightarrow$ PT-symmetry breaking



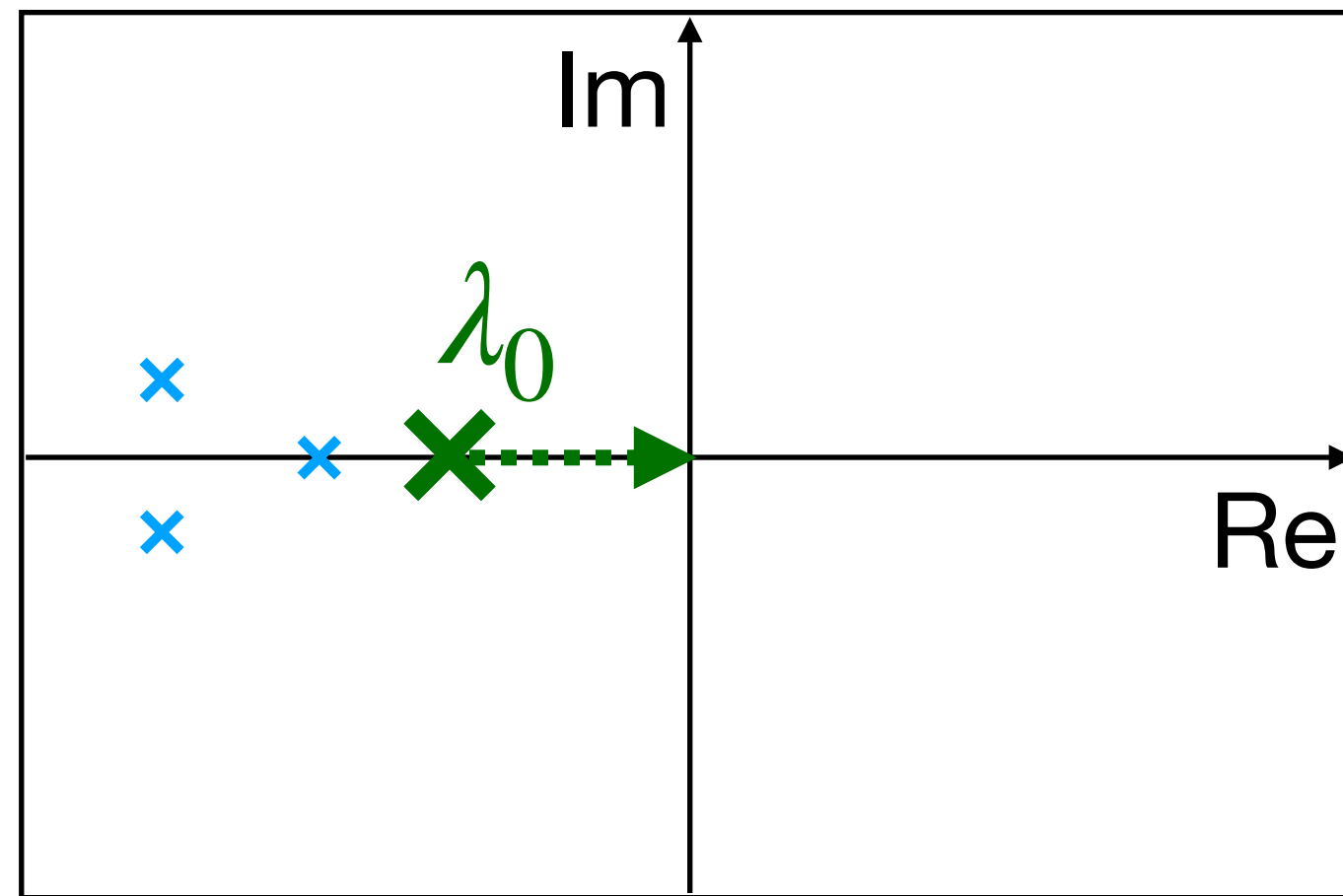
Time-reversed solution



3 different transition scenarios

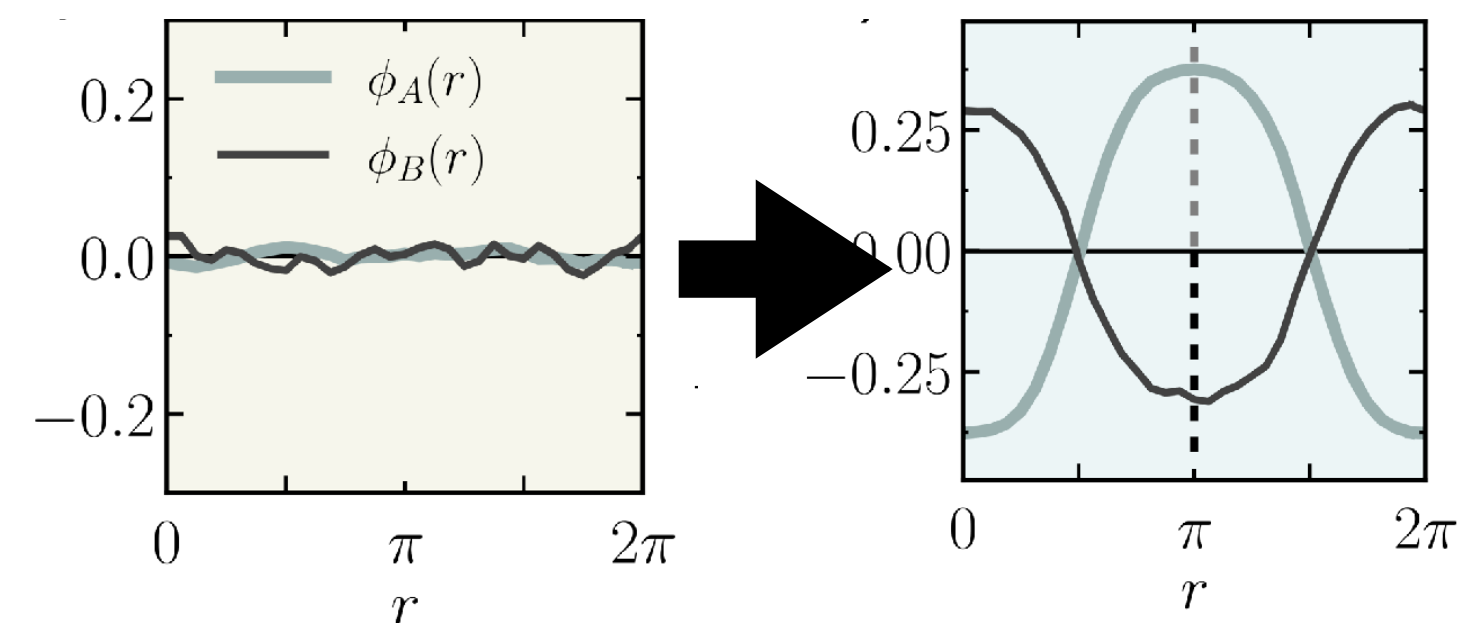
Static—Static transition

Conventional Critical Point



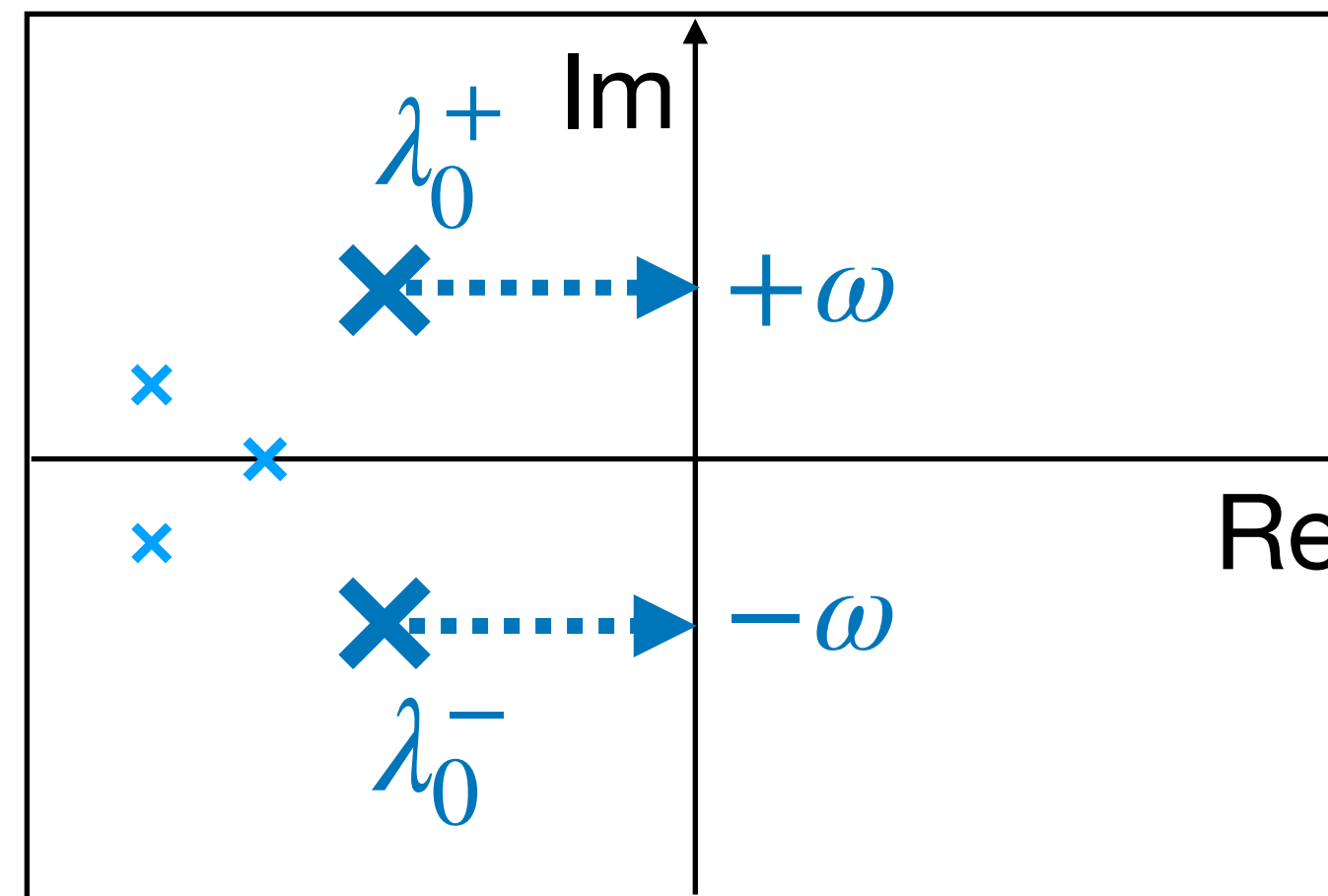
Primary Bifurcation

Hom. → Static demixed



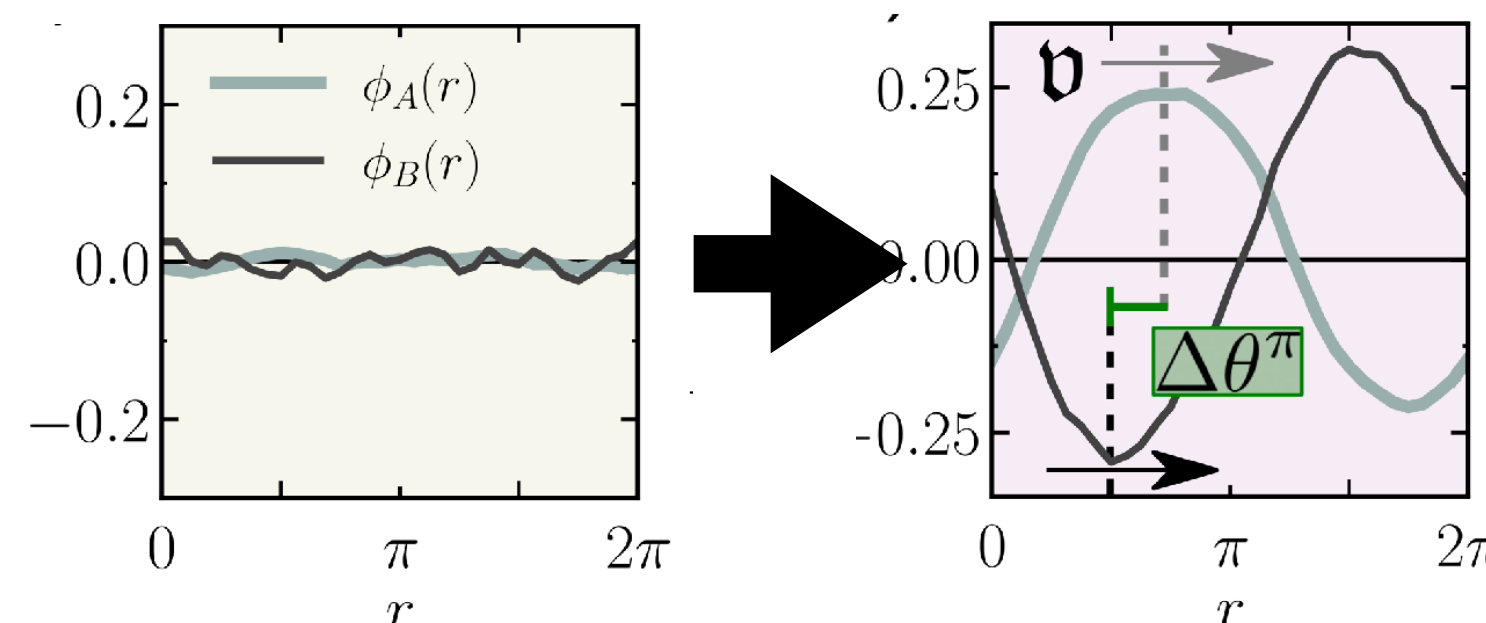
Static—Dynamic transitions

Oscillatory Instability

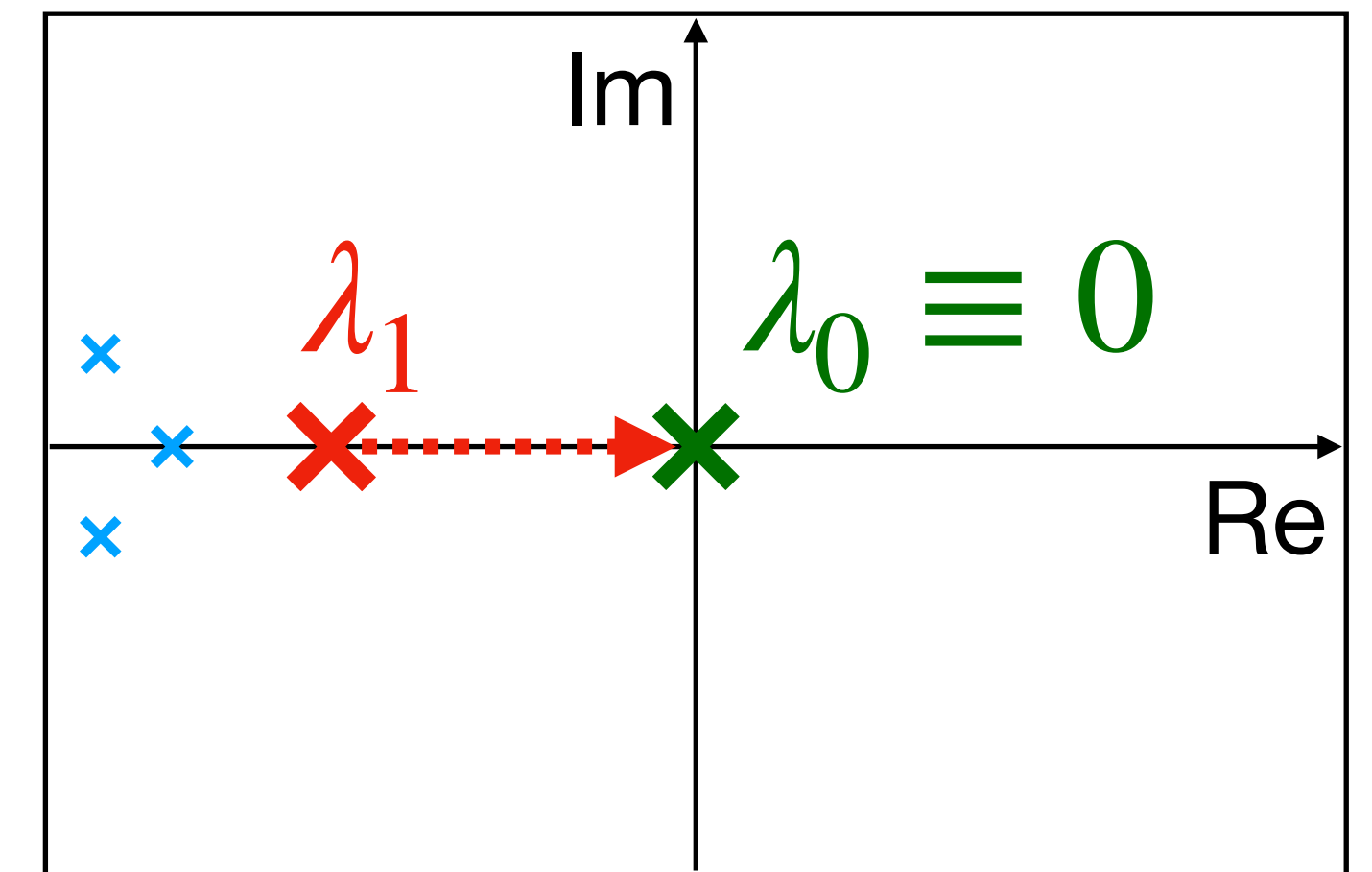


Primary Bifurcation

Hom. → Travelling wave

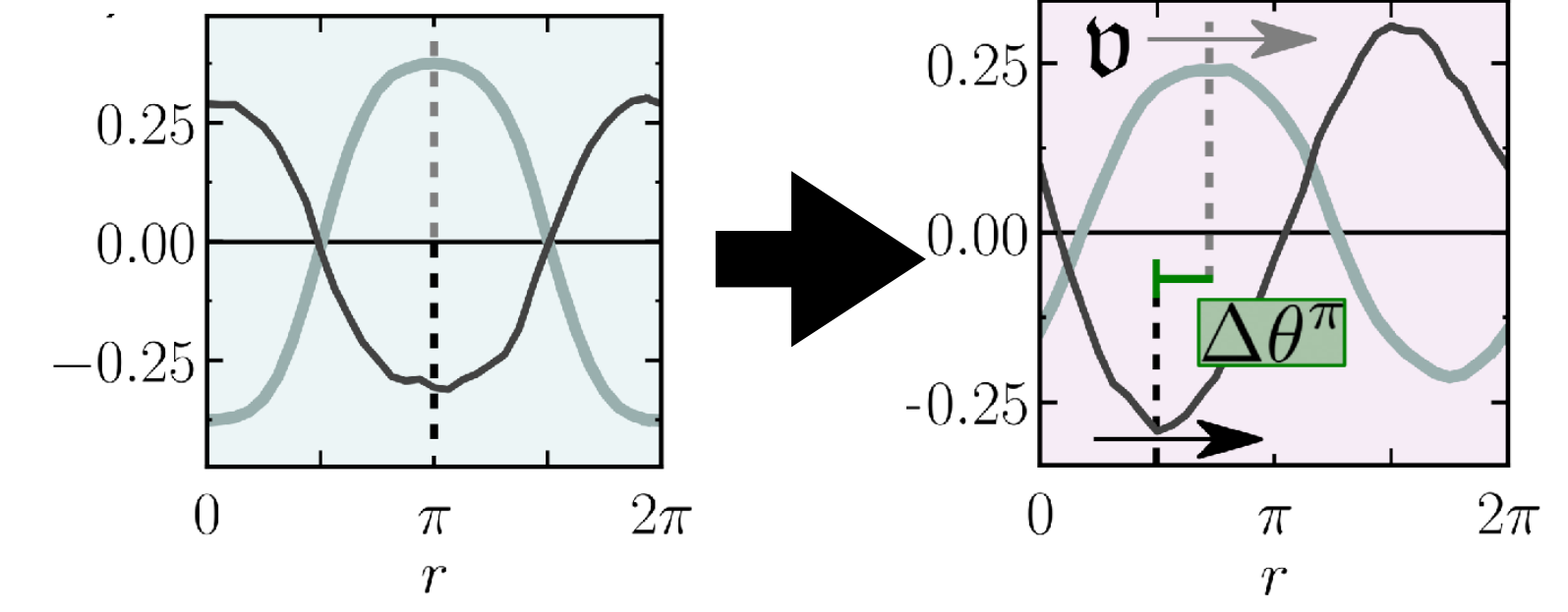


Critical Exceptional Point (CEP)



Secondary Bifurcation

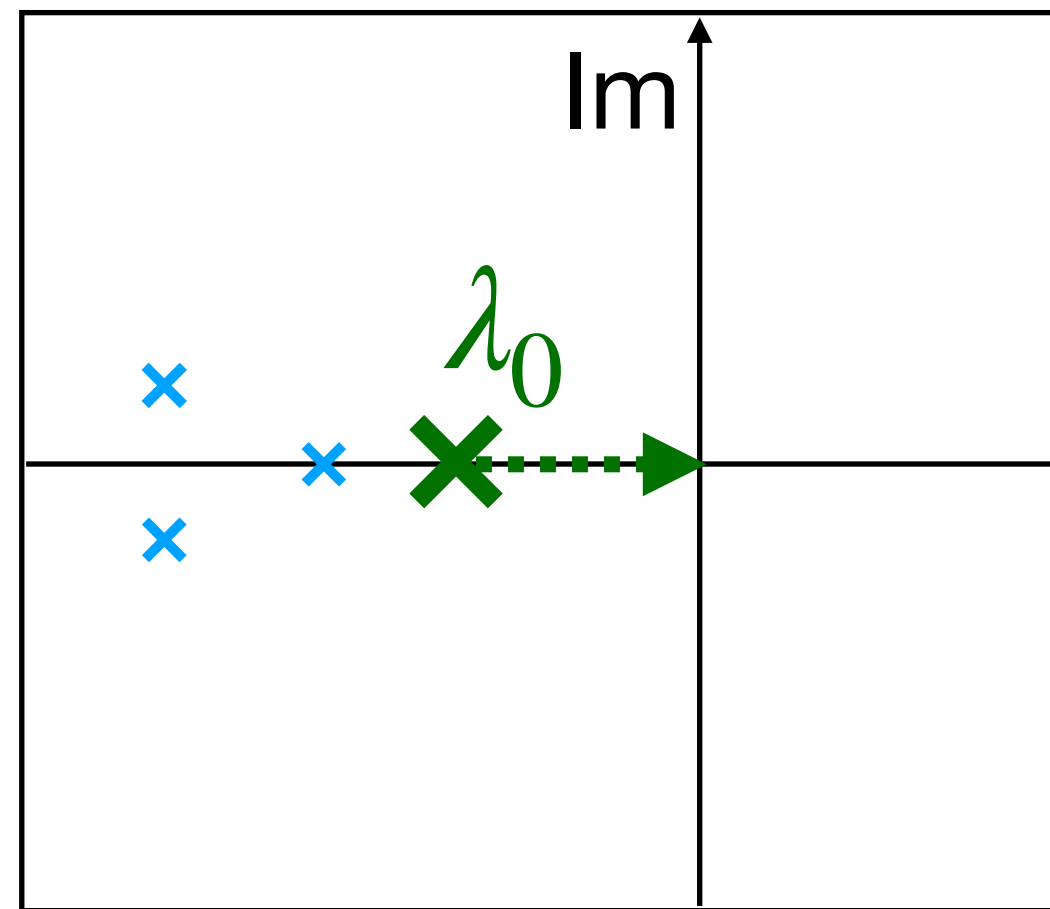
Static dem. → Travelling wave



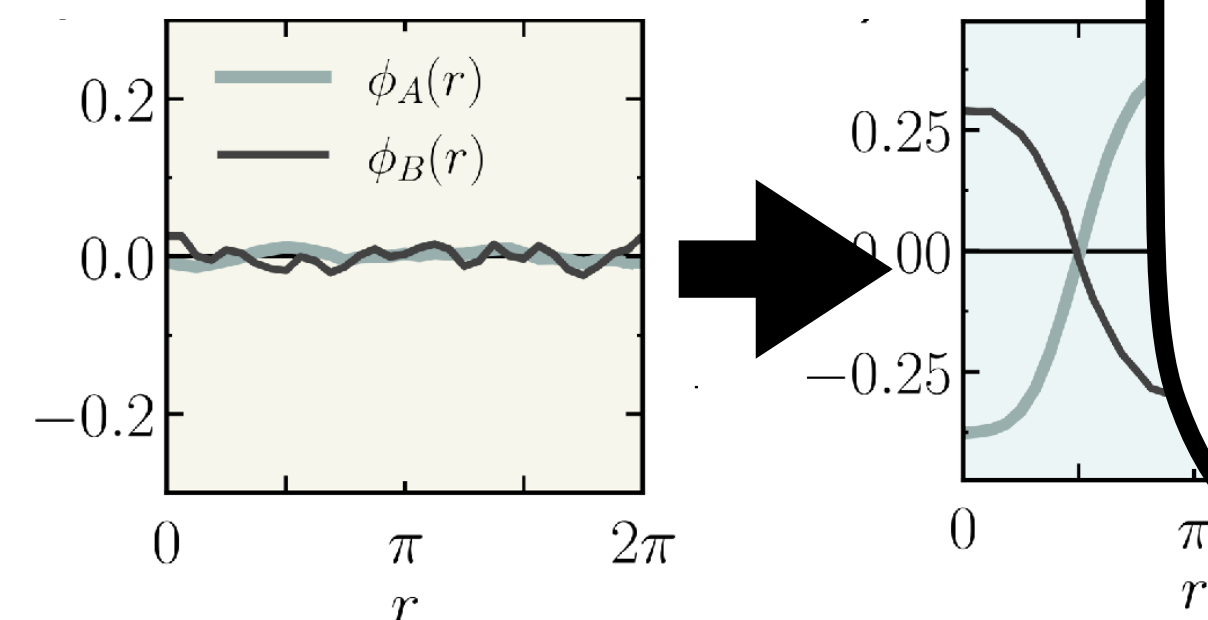
3 different transition scenarios

Static—Static transition

Conventional Critical

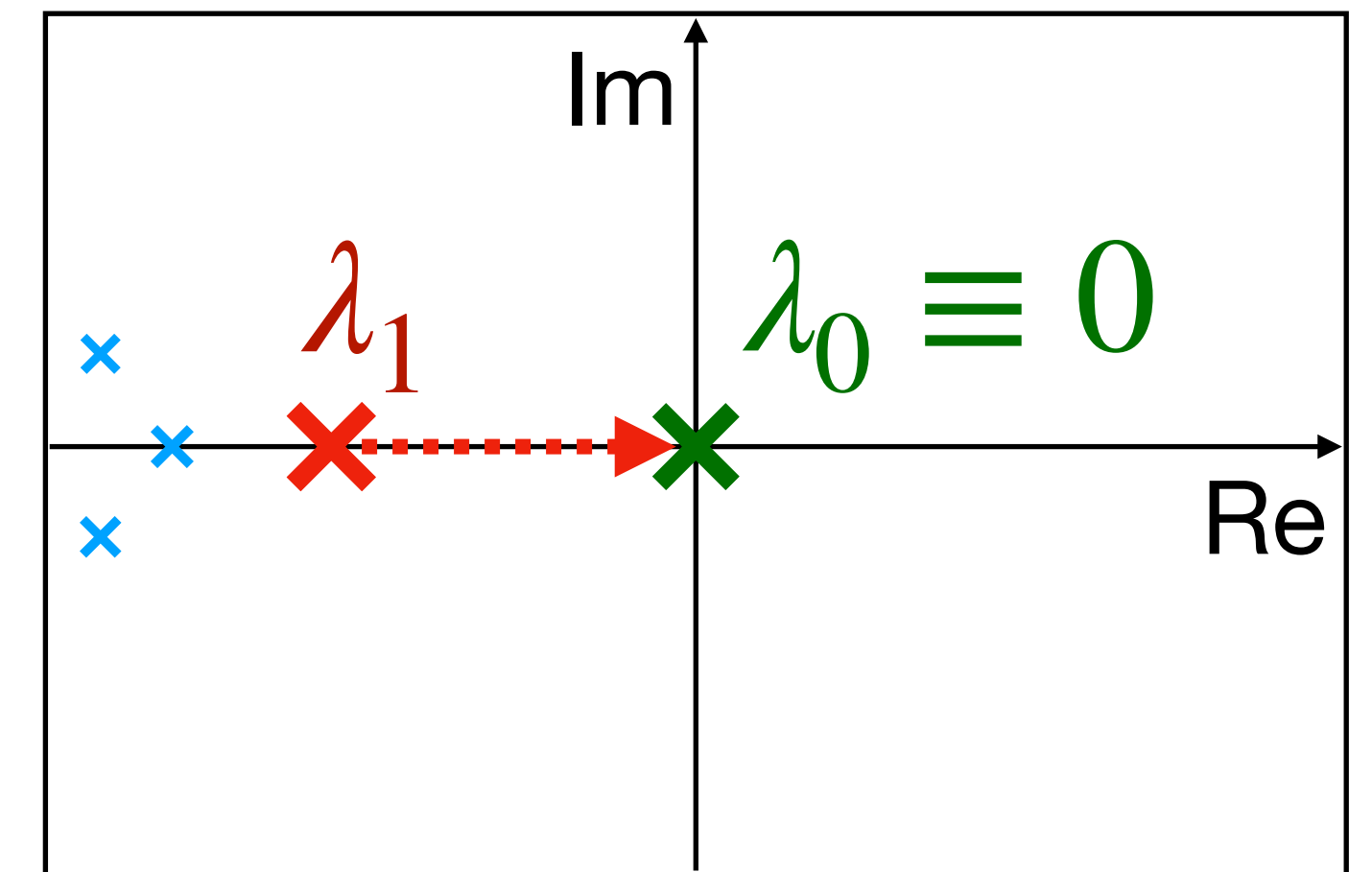


Primary Bifurcation
Hom. → Static dem.

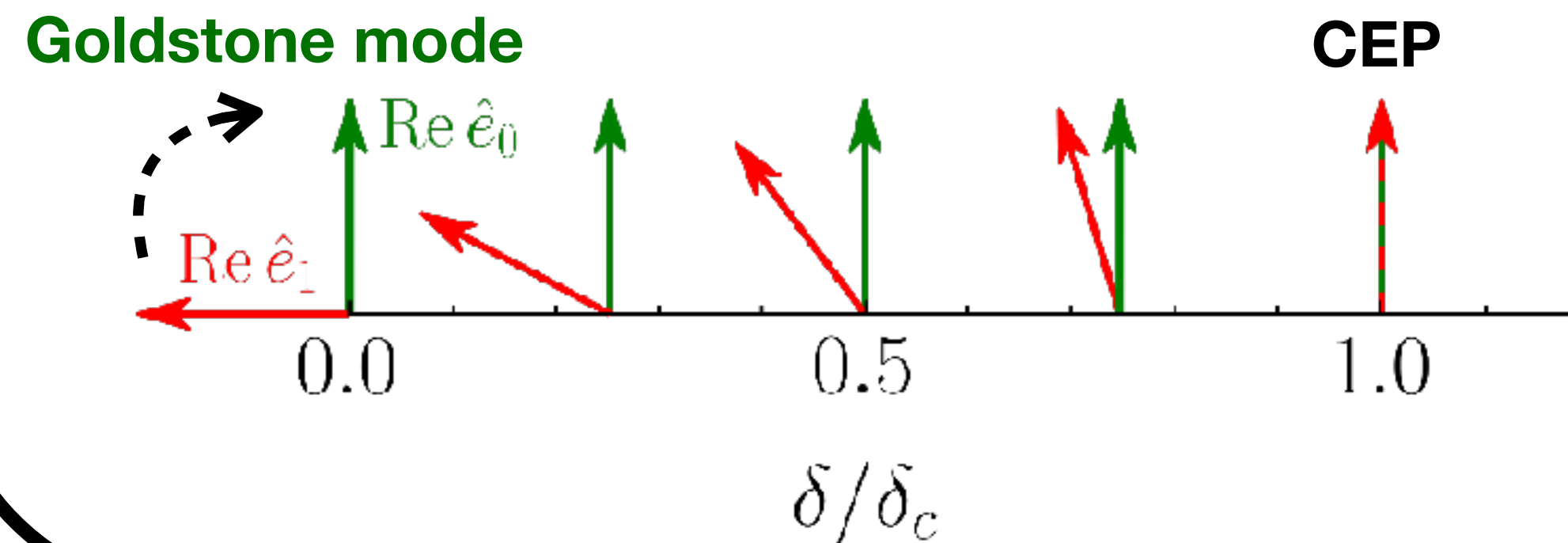
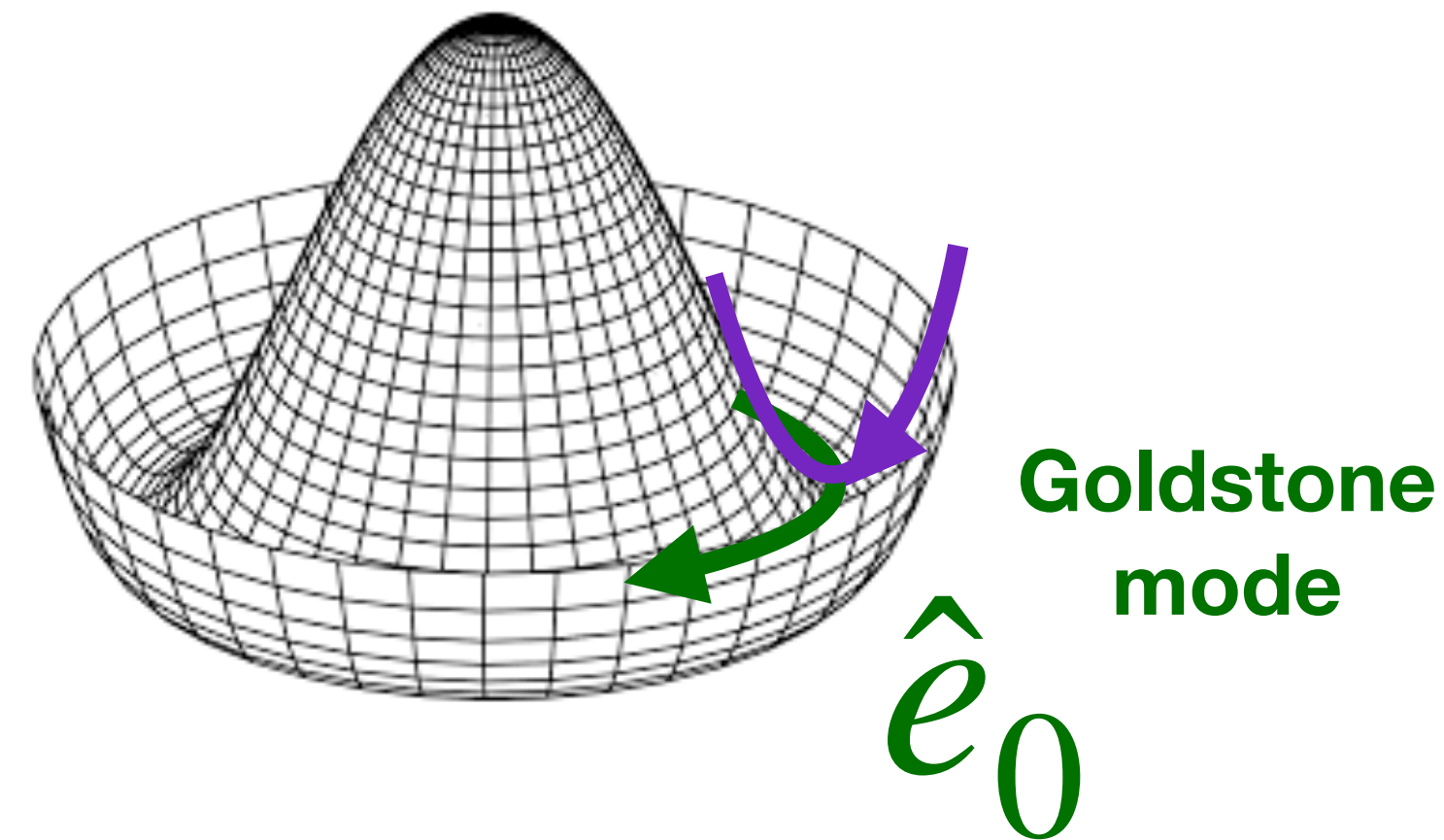
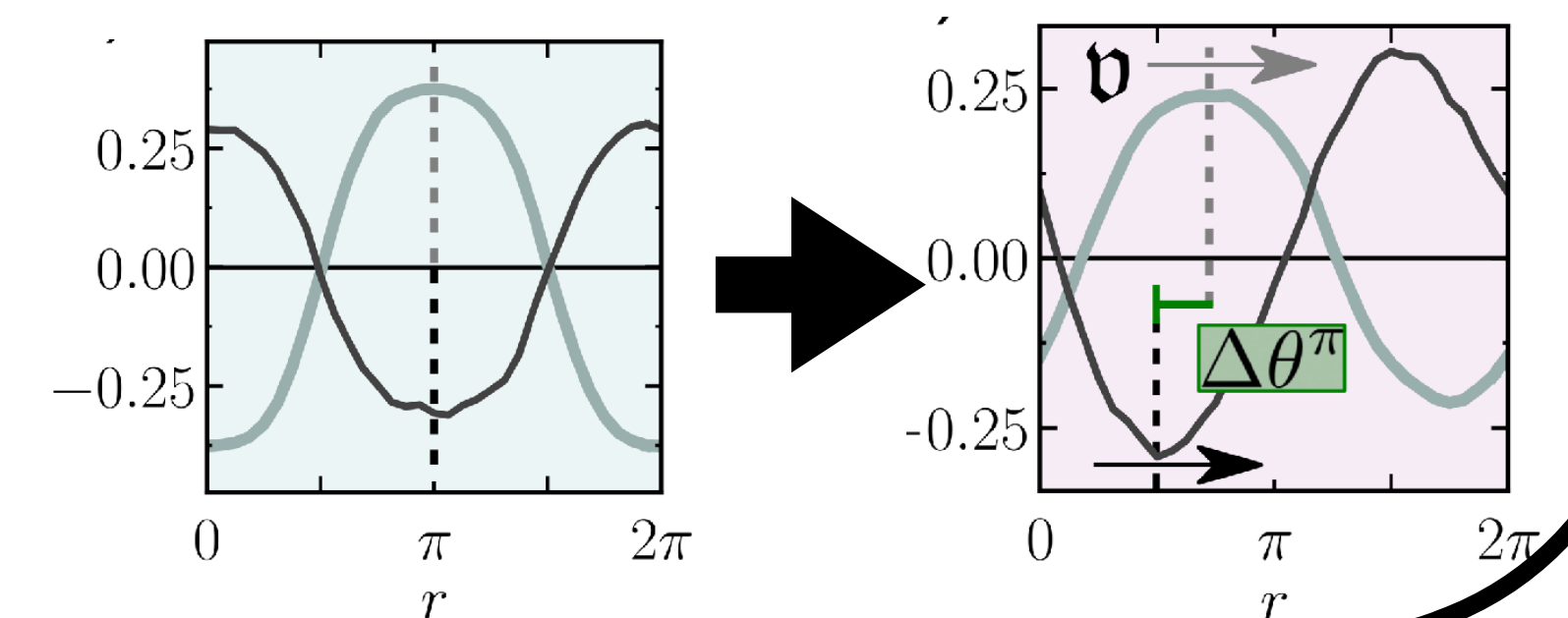


Static—Dynamic transitions

Critical Exceptional Point (CEP)



Secondary Bifurcation
Static dem. → Travelling wave



Characteristics of fluctuations close to transitions

- Typical size $\chi = \lim_{\epsilon \rightarrow 0} \sum_i \int_V dr \left\langle (\phi_i - \phi_i^*)^2 \right\rangle / \epsilon$

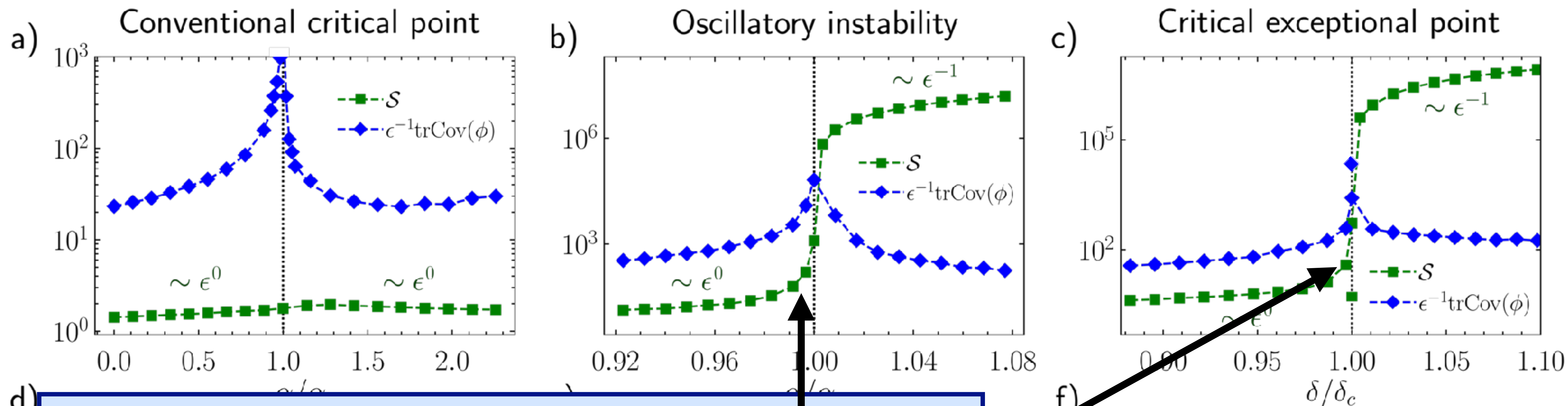
$$\dot{\phi}_i = - \nabla \cdot \left(J_i^d + \sqrt{2\epsilon} \Lambda_i \right)$$

- TRSB (“Nonequilibrium-ness”) $\dot{S} = \left\langle \ln \frac{\mathbb{P}[\phi(t) |_0^T]}{\mathbb{P}_R[\phi_R(t) |_0^T]} / T \right\rangle$
[Nardini et al., PRX (2017)]

Static-to-Static	• Conventional Critical Points:	$\chi \sim \text{Re}[\lambda_0^k] \rightarrow \infty$	$\dot{S}$ finite
Static-to-Dynamic	• Oscillatory Instabilities:		$\dot{S} \sim \frac{\omega^2}{ q^l ^2} \chi \rightarrow \infty$
	• Critical Exceptional Point:		$\dot{S} \sim \frac{1}{ q^l ^2} \chi \rightarrow \infty$

Prediction: $\dot{S}$ heralds approach to dynamical phase transition

- In dynamical phase: $\dot{S} \sim \epsilon^{-1} (1/\tau)^2$

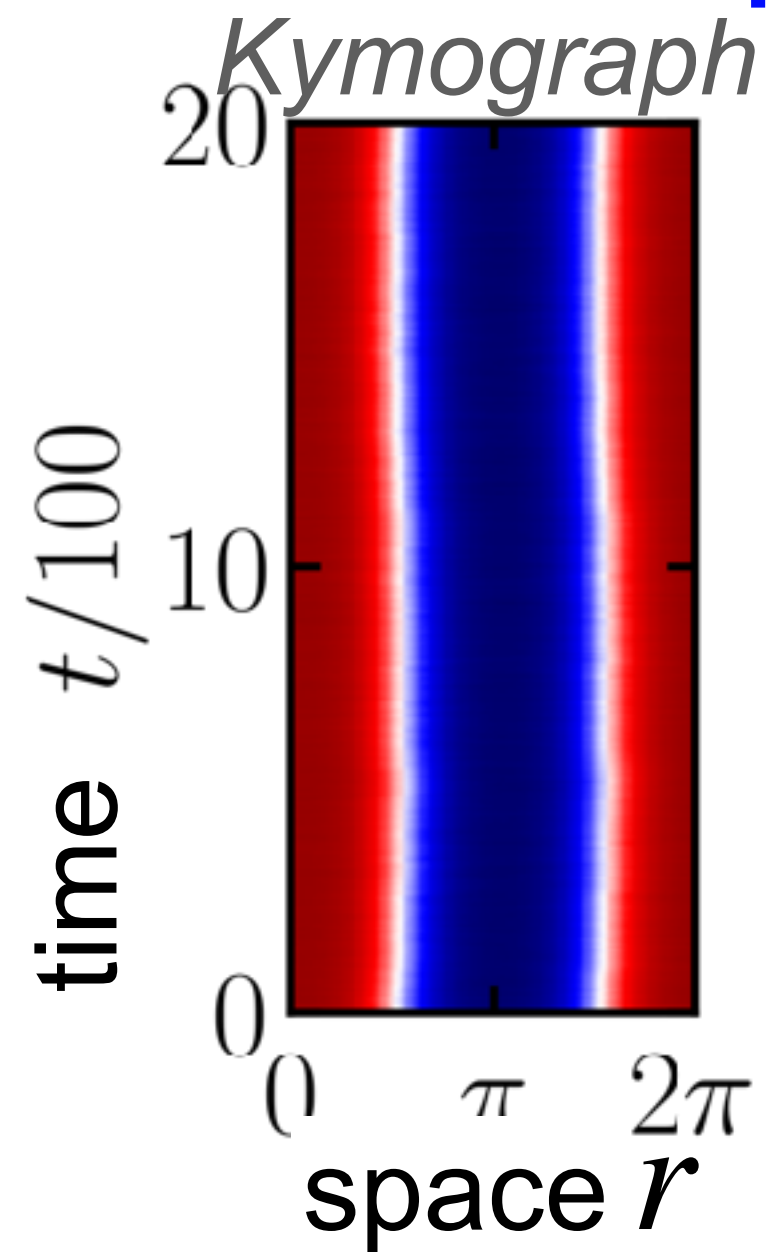


both diverge

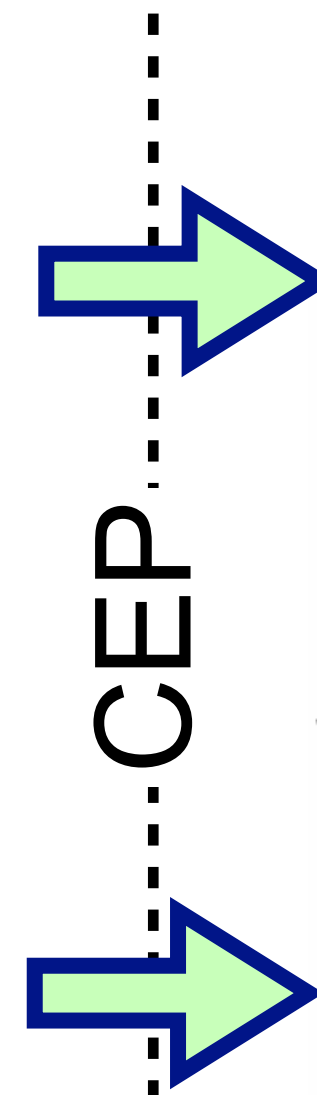
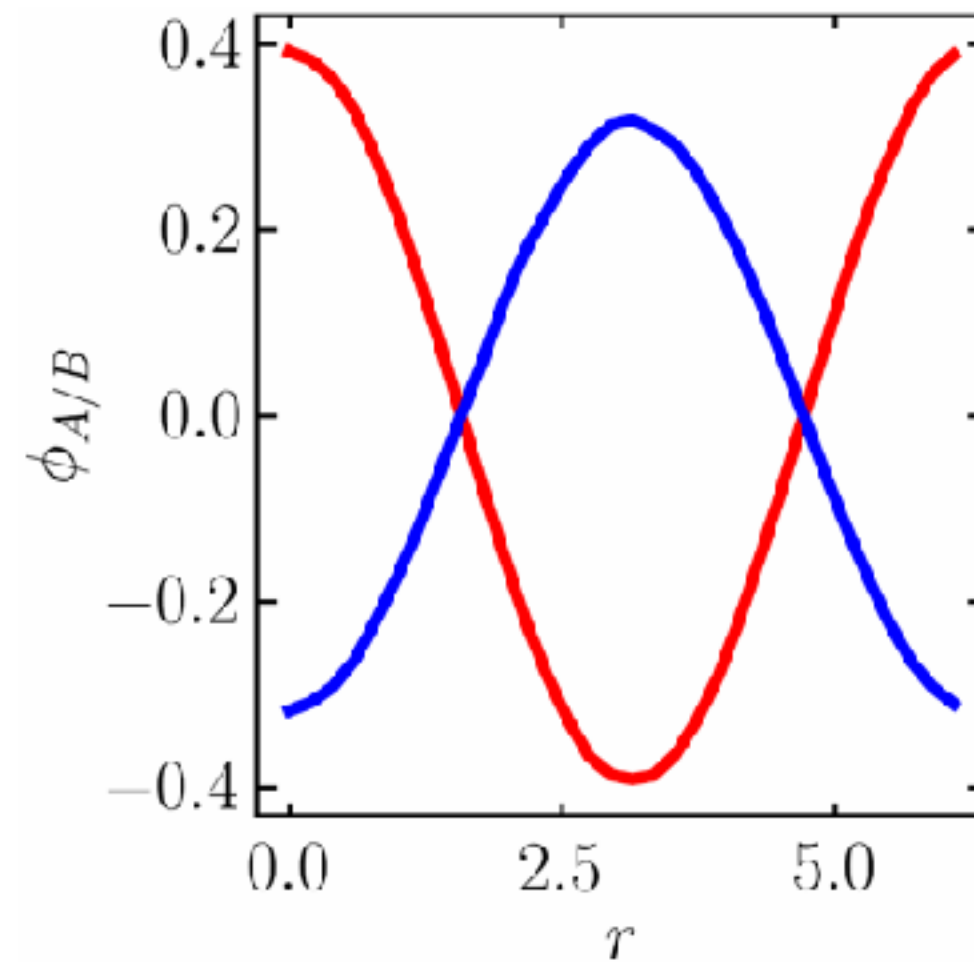
$$\dot{S} \sim \chi \sim \text{Re}[\lambda_0^k]^{-1}$$

production ($S \propto \epsilon^0$)

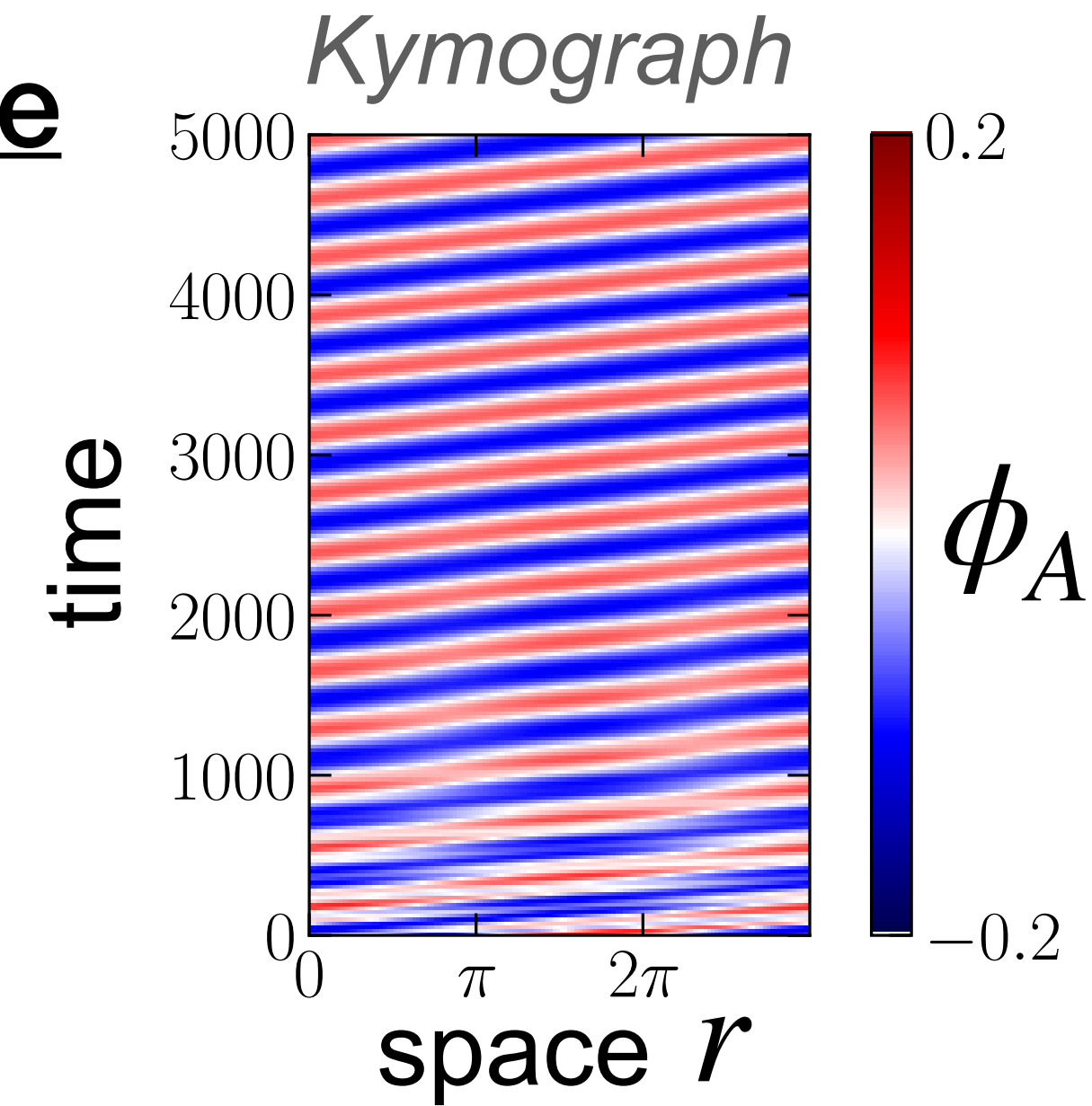
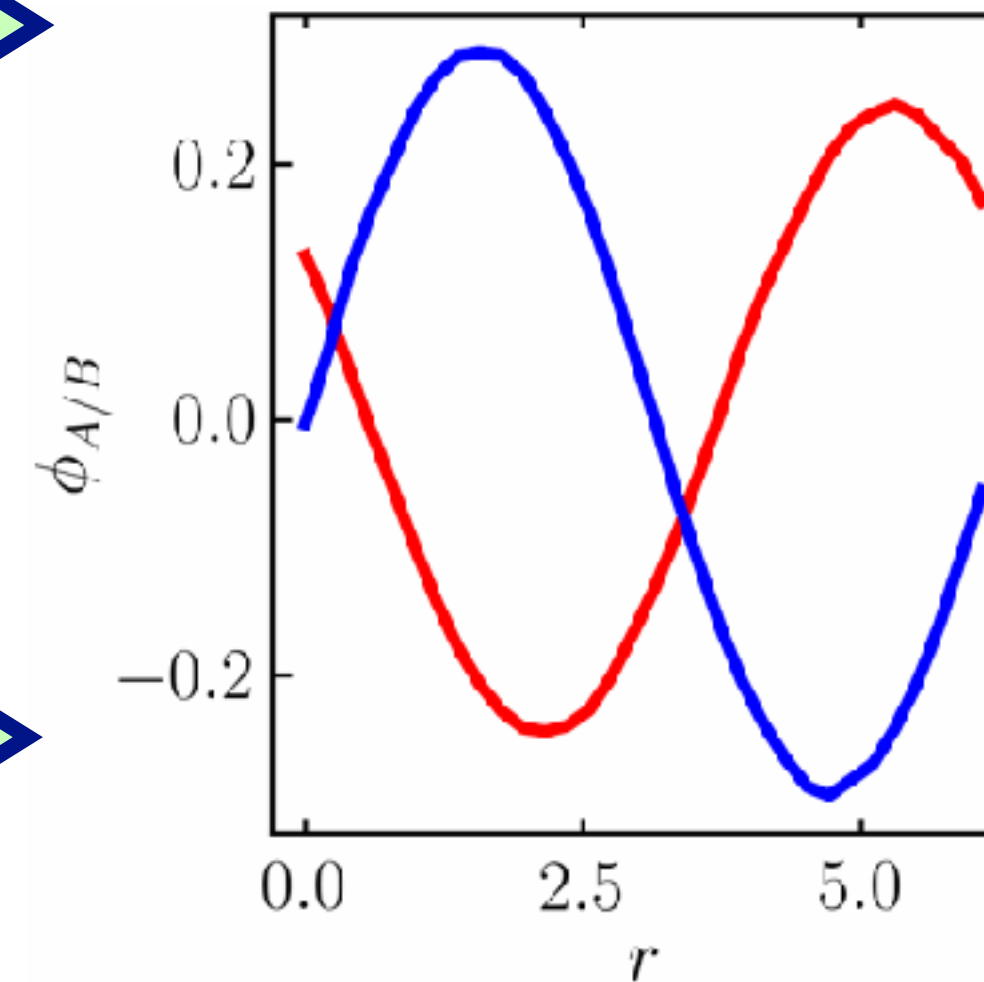
Spectral singularity: Critical Exceptional Point (CEP)



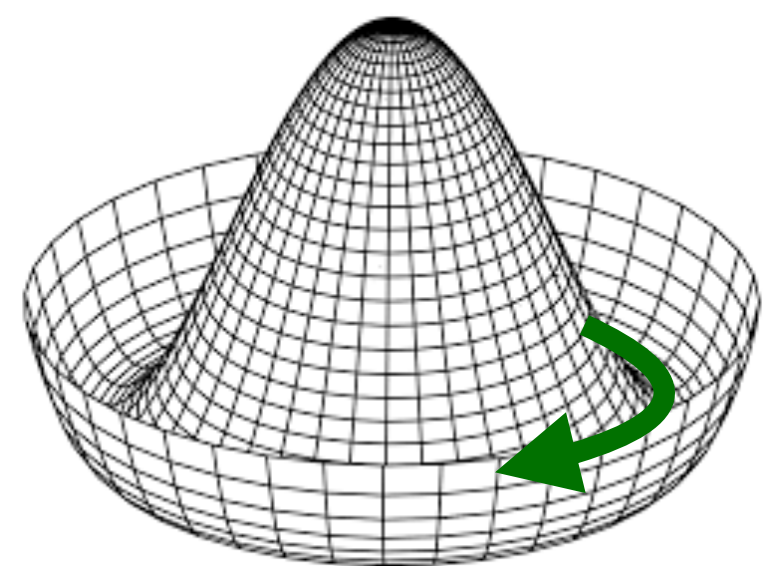
Static demixed



Travelling wave

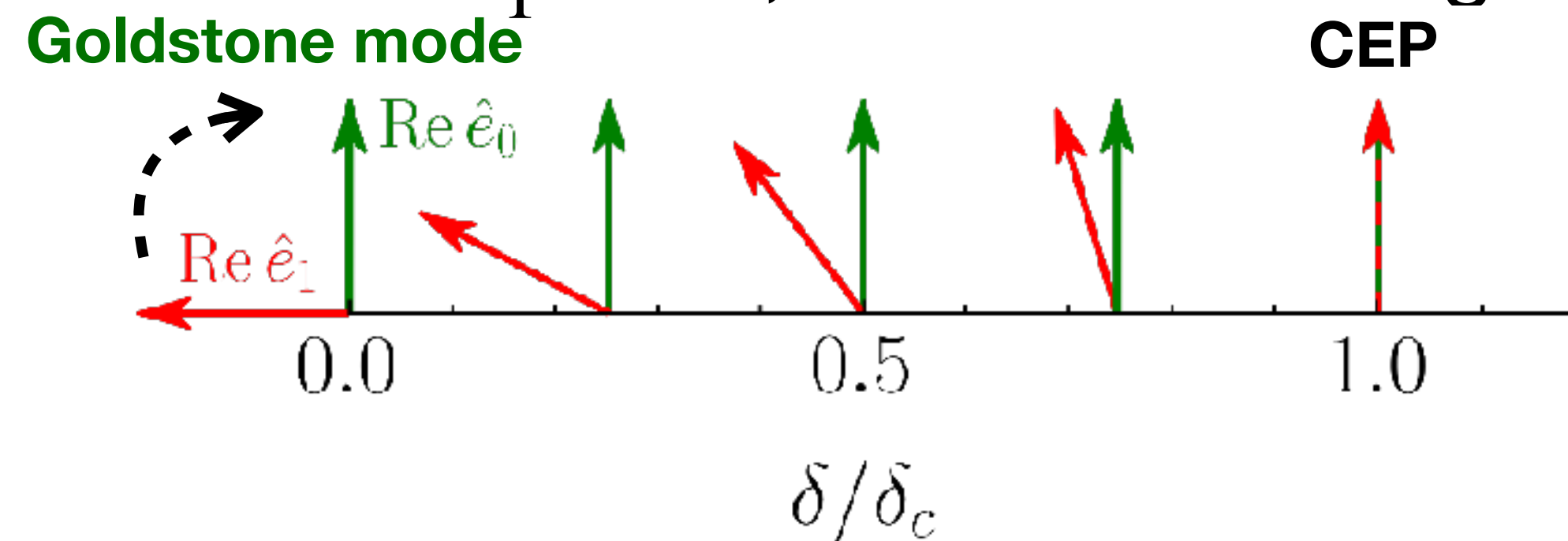


- Before CEP: Broken translational symmetry
→ Goldstone mode ($\hat{e}_0$, $\lambda_0 = 0$)



Goldstone mode
 $\hat{e}_0$

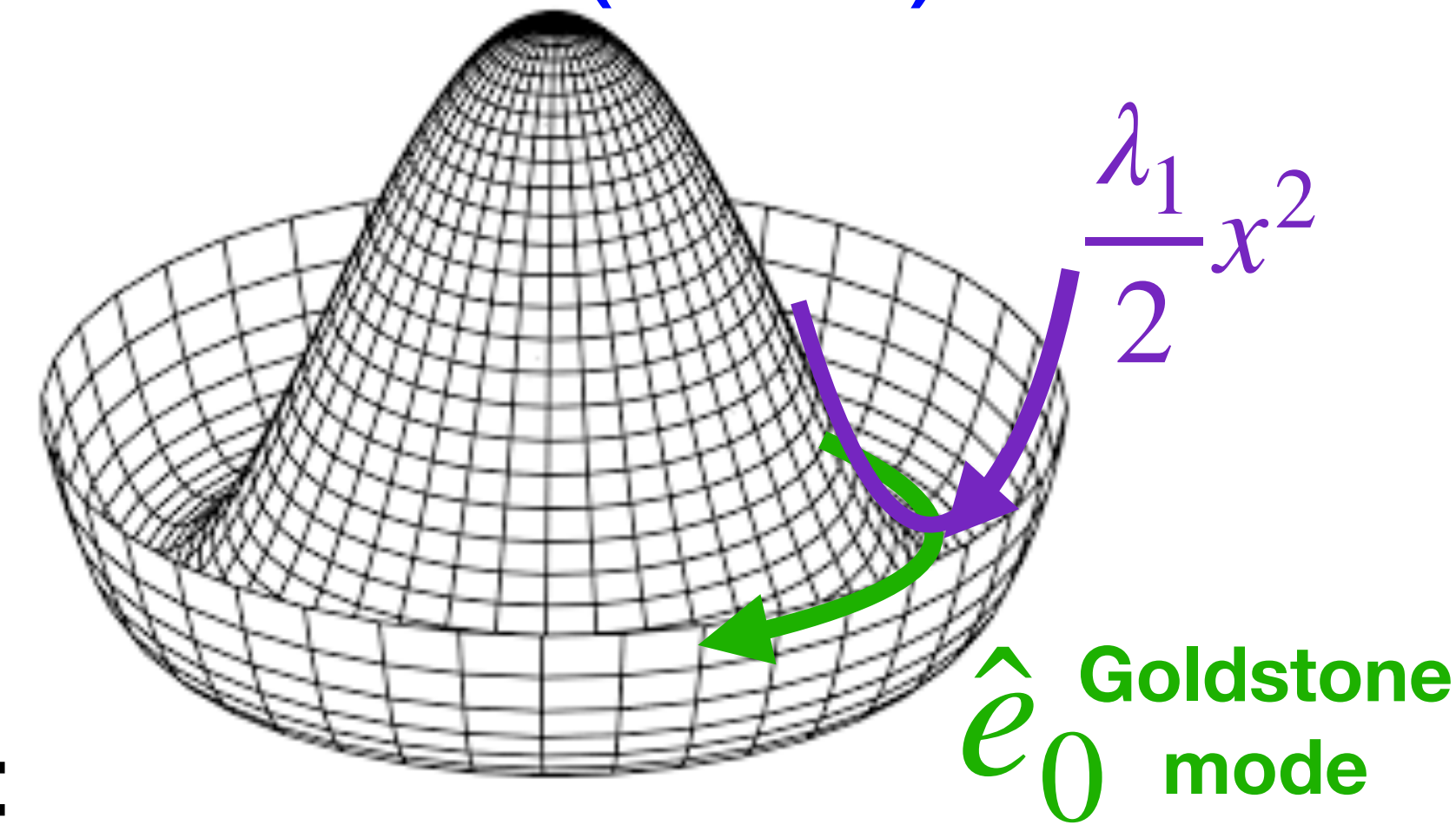
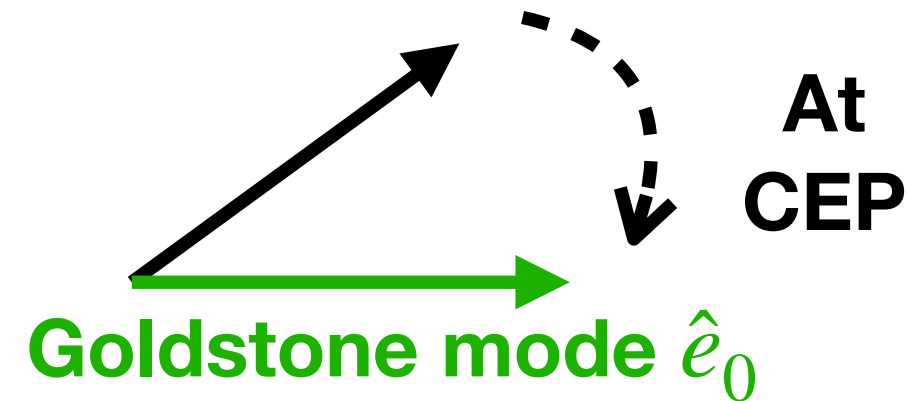
- At CEP: another eigenmode of $\mathcal{L}$ becomes unstable: $\lambda_1 \rightarrow 0$, while coalescing with $\hat{e}_0$



Entropy production at “Critical Exceptional Point” (CEP)

- Dynamics of relevant modes in eigenbasis of $\mathcal{L}$ $\{\hat{e}_0, \hat{e}_1\}$:

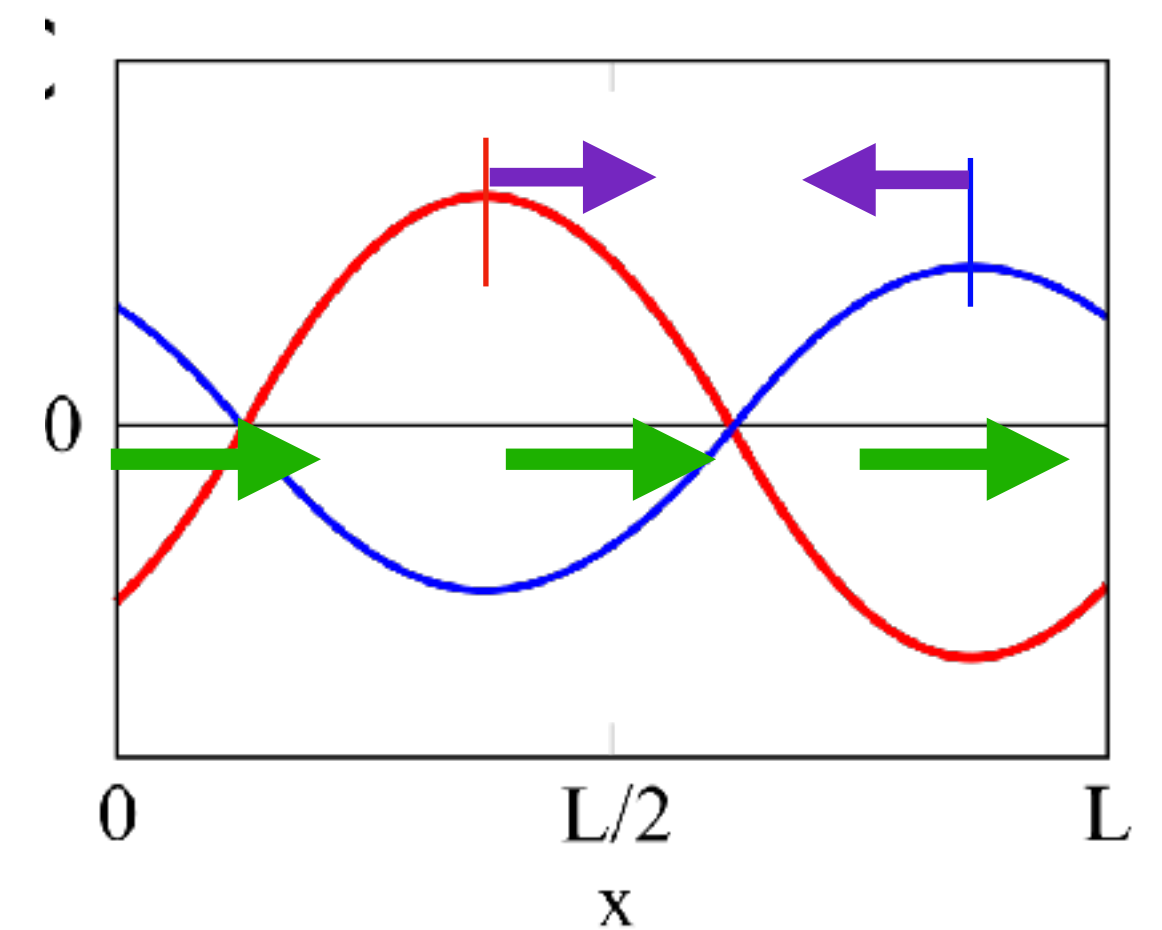
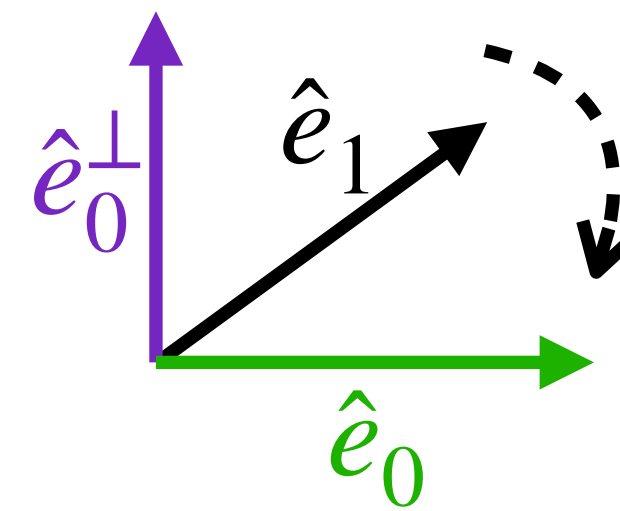
$$\frac{d}{dt} \begin{pmatrix} \psi_0 \\ \psi_1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & \lambda_1 \end{pmatrix} \begin{pmatrix} \psi_0 \\ \psi_1 \end{pmatrix} + \underline{\xi}$$



- Dimension reduction of $\mathcal{L}$ at CEP: Change basis to $\{\hat{e}_0, \hat{e}_0^\perp\}$:

$$\frac{d}{dt} \begin{pmatrix} \theta_c \\ \Delta\theta^\pi \end{pmatrix} = \begin{pmatrix} 0 & 2\delta \\ 0 & \lambda_1 \end{pmatrix} \begin{pmatrix} \theta_c \\ \Delta\theta^\pi \end{pmatrix} + \underline{\xi}$$

$$\Delta\theta^\pi \propto [\theta_A^1 - (\theta_B^1 + \pi)]$$



- Small perturbation ϵ in direction $\hat{e}_0^\perp$ yields:

$$\theta_c(t) = \frac{\epsilon 2\delta}{\lambda_1} (e^{\lambda_1 t} - 1)$$

$$\Delta\theta^\pi(t) = \epsilon e^{\lambda_1 t}$$

- At CEP: $\lambda_1 \rightarrow 0$: Giant noise amplification

$\Rightarrow$ *Squeezing patterns* $\rightarrow$ *Shifting patterns along Goldstone*

$\Rightarrow$ *TRSB fluctuations* $\rightarrow$ **Real space interpretation?**

Entropy production at “Critical Exceptional Point” (CEP)

- *Towards CEP: Patterns move ballistically*

$$\frac{d}{dt} \begin{pmatrix} \theta_c \\ \Delta\theta^\pi \end{pmatrix} = \begin{pmatrix} 0 & 2\delta \\ 0 & \lambda_1 \end{pmatrix} \begin{pmatrix} \theta_c \\ \Delta\theta^\pi \end{pmatrix} + \xi$$

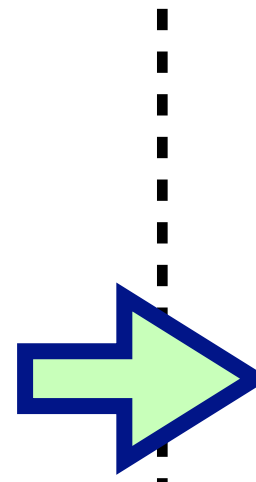
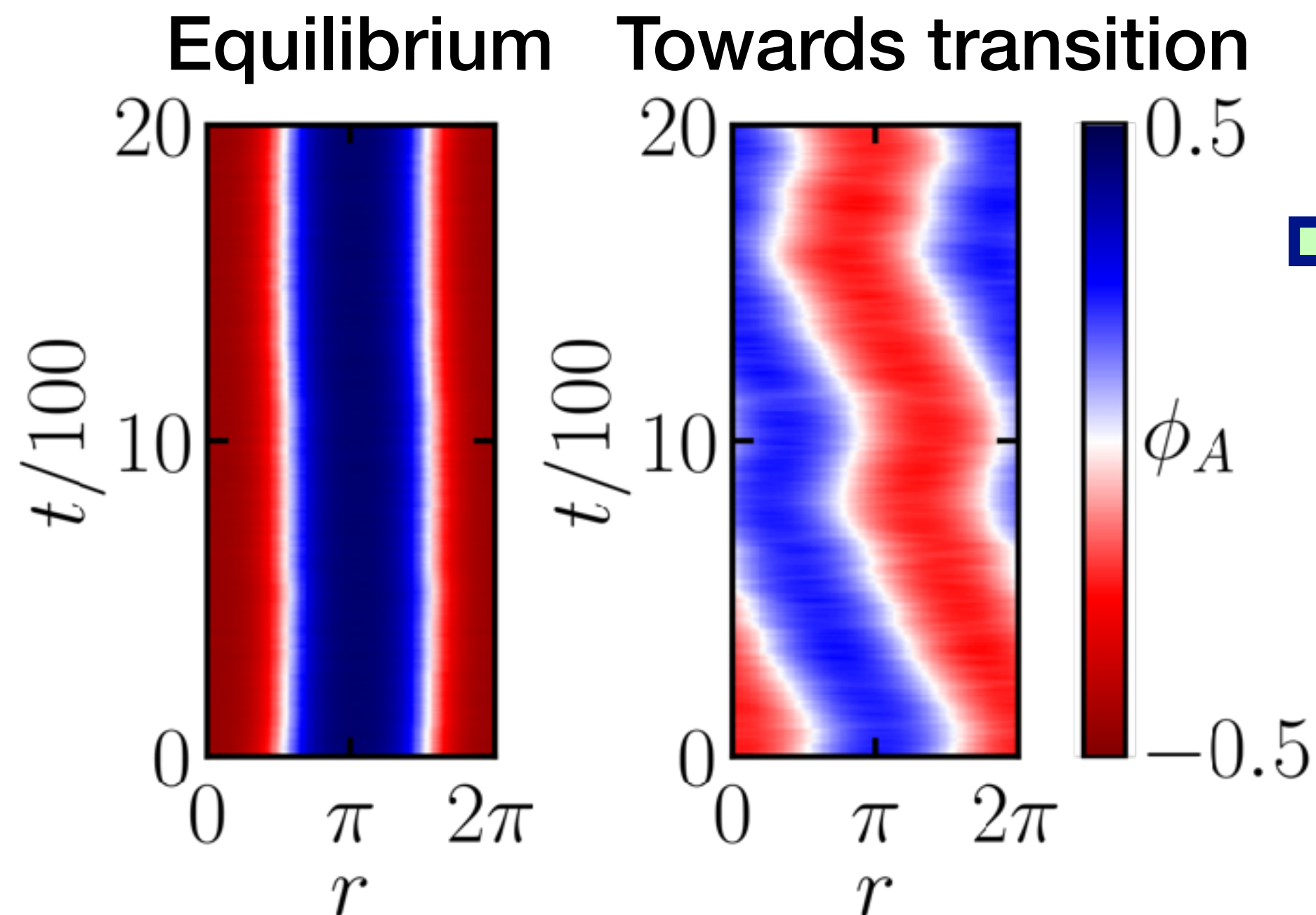
→ Can be mapped onto AOUP model

→ “self-propelled patterns”

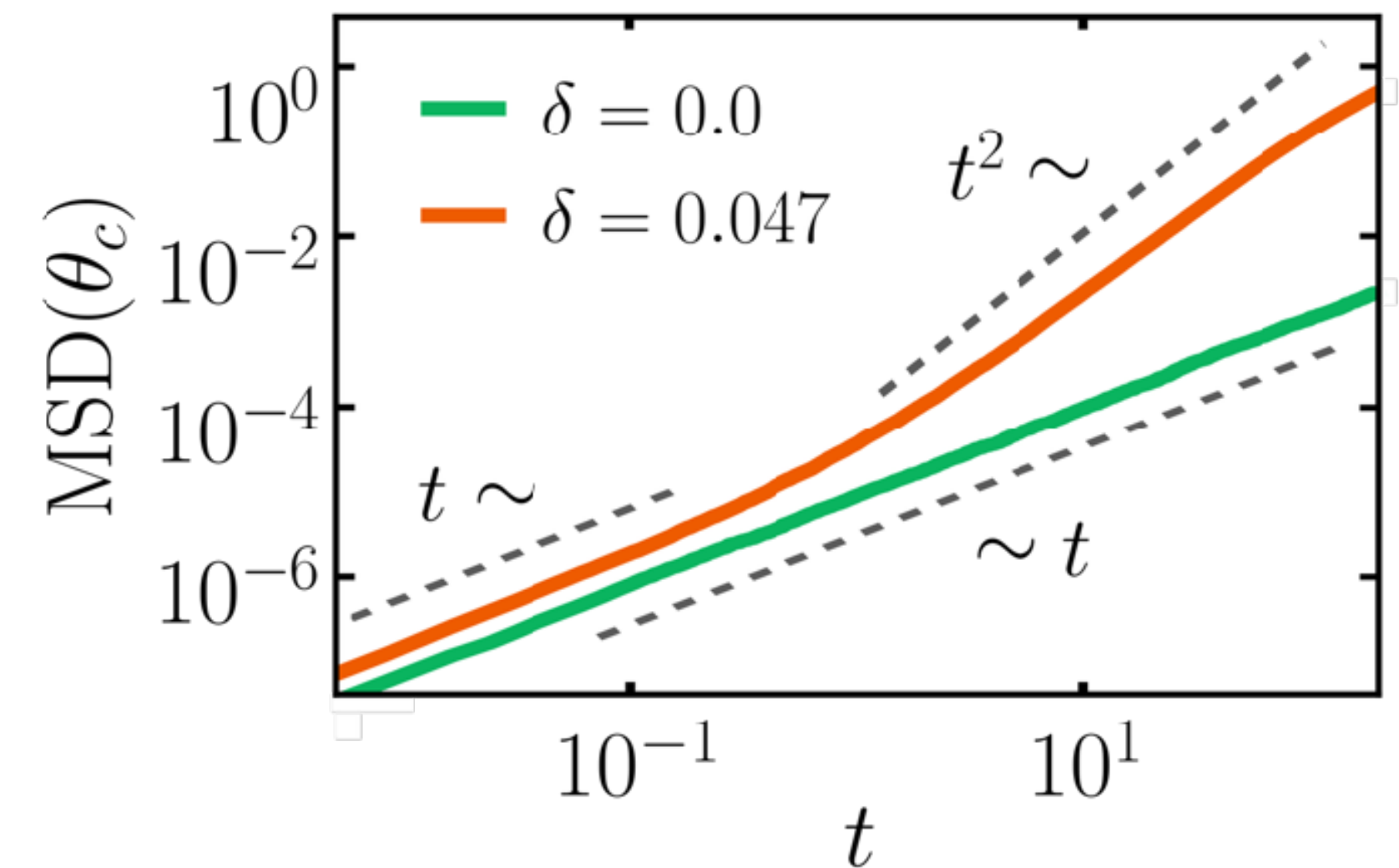
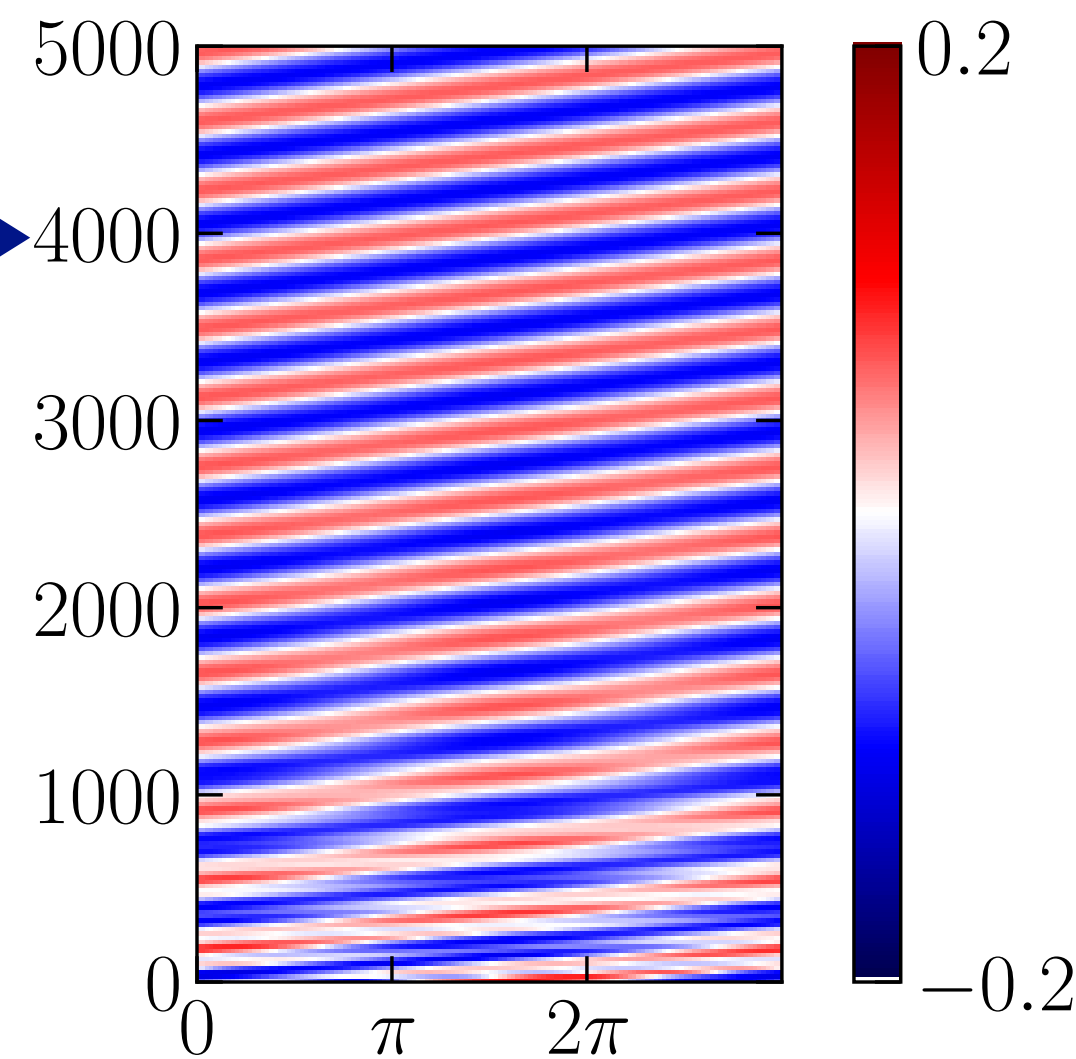


$$\Delta\theta^\pi \propto [\theta_A^1 - (\theta_B^1 + \pi)]$$

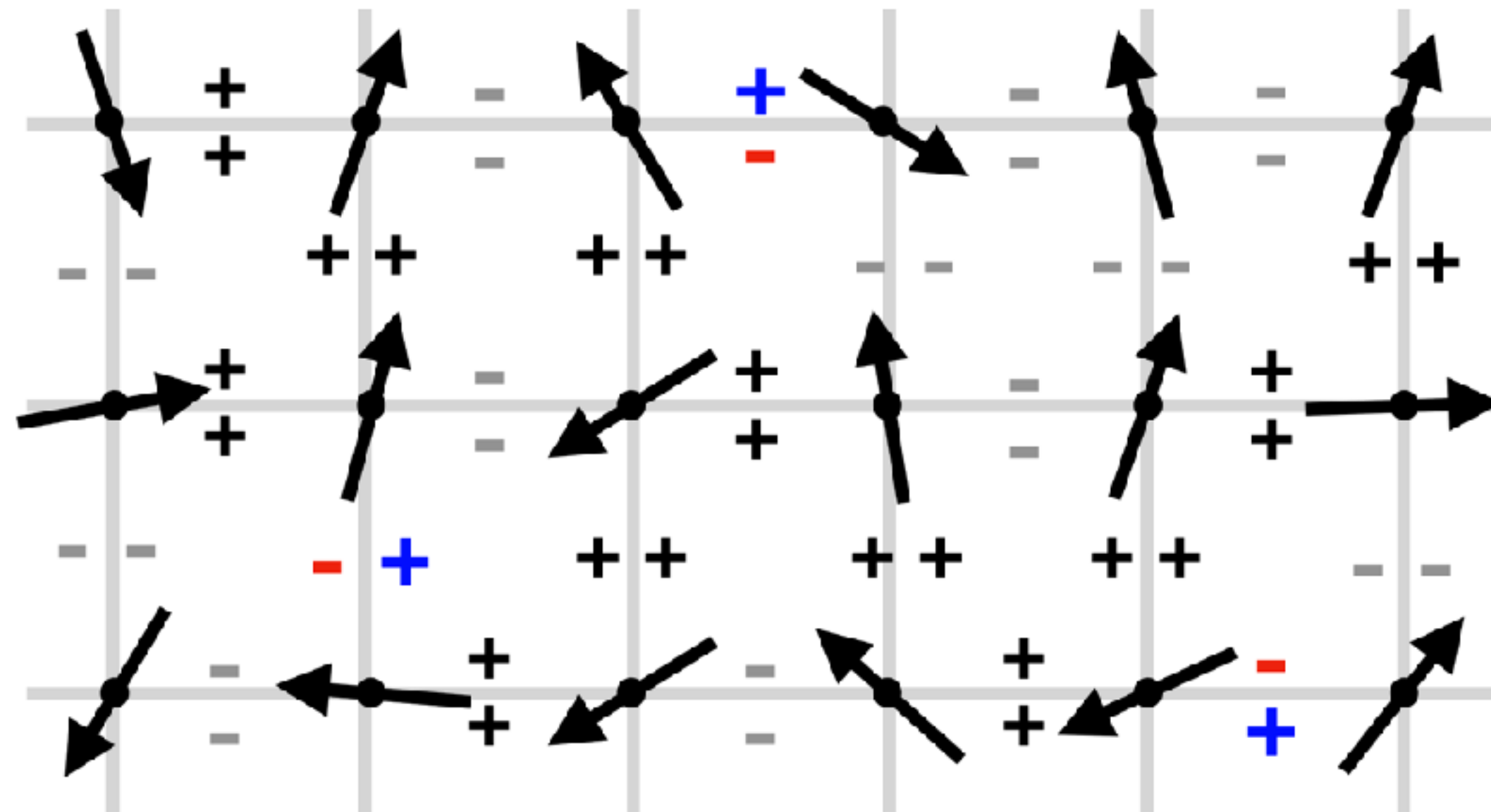
Static demixed



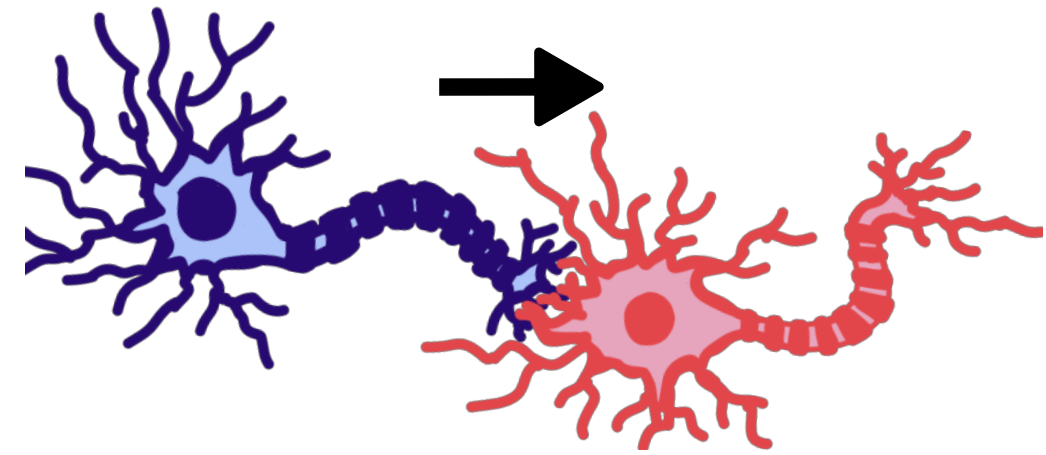
Travelling wave



One-species NR system with disorder

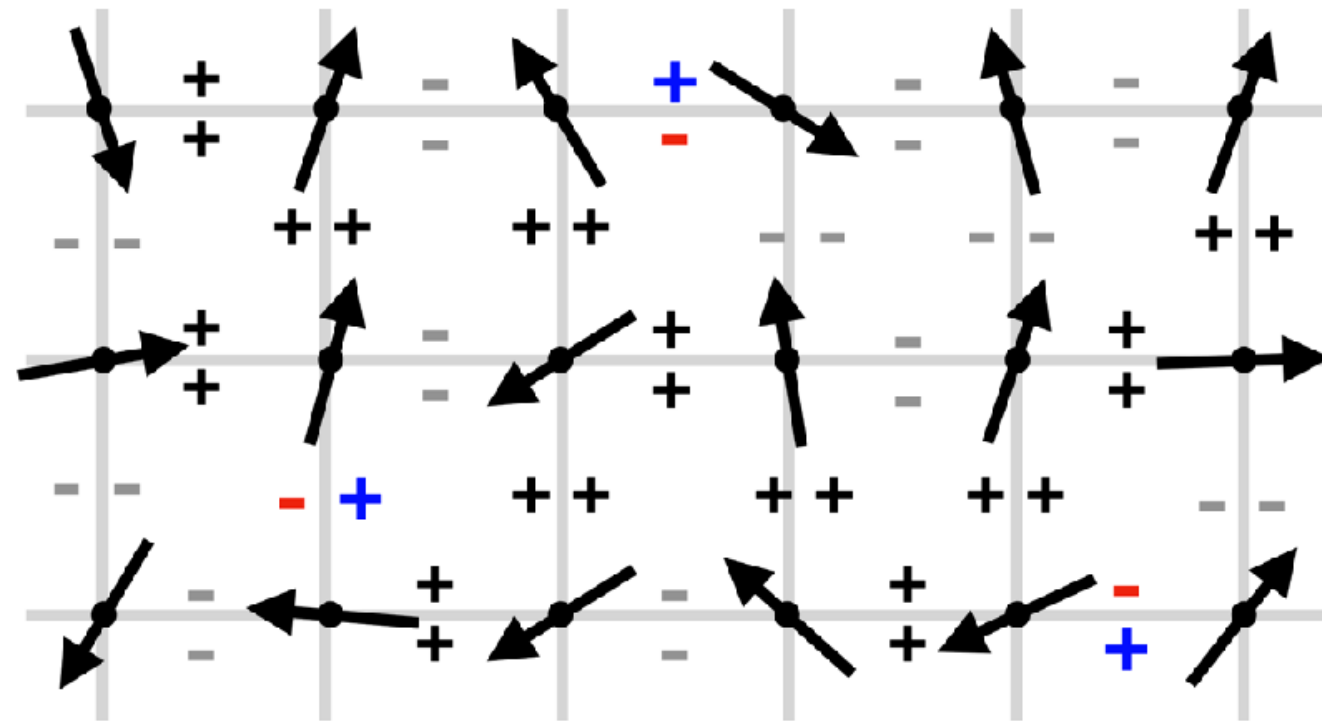


Single species NR XY model with quenched disorder



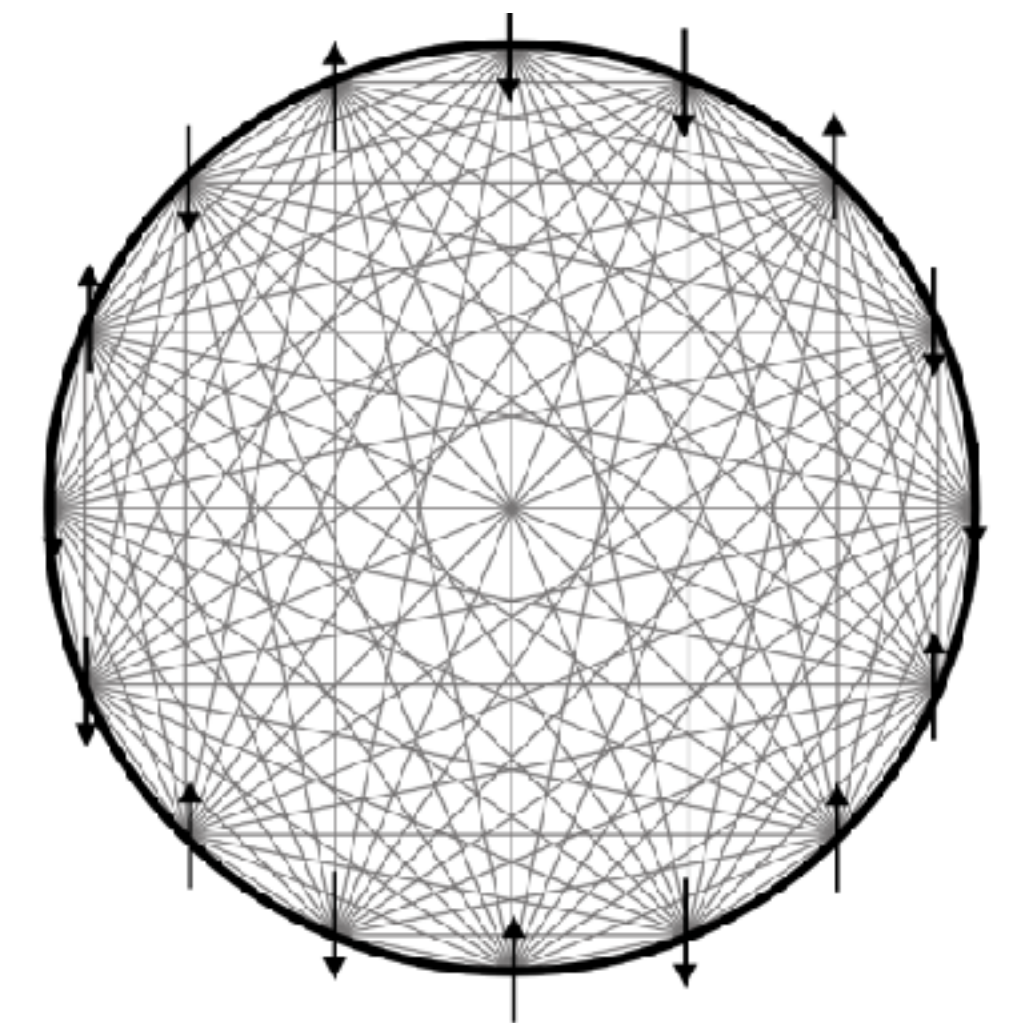
$$\dot{\phi}_i = - \sum_j J_{ij} \sin(\phi_i - \phi_j) + \sqrt{2D} \xi_i$$

$\{J_{ij}\} \sim$ some distribution (quenched disorder) and can be asymmetric



- $|J_{ij}| = 1$
- $\mathbb{E}(J_{ij}) = 0 \rightarrow$ **Frustrated system**
- With probability $(1 - p)$, $J_{ij} = J_{ji}$
- With probability p , $J_{ij} = -J_{ji} \rightarrow$ **NR system**

- Fully connected graph

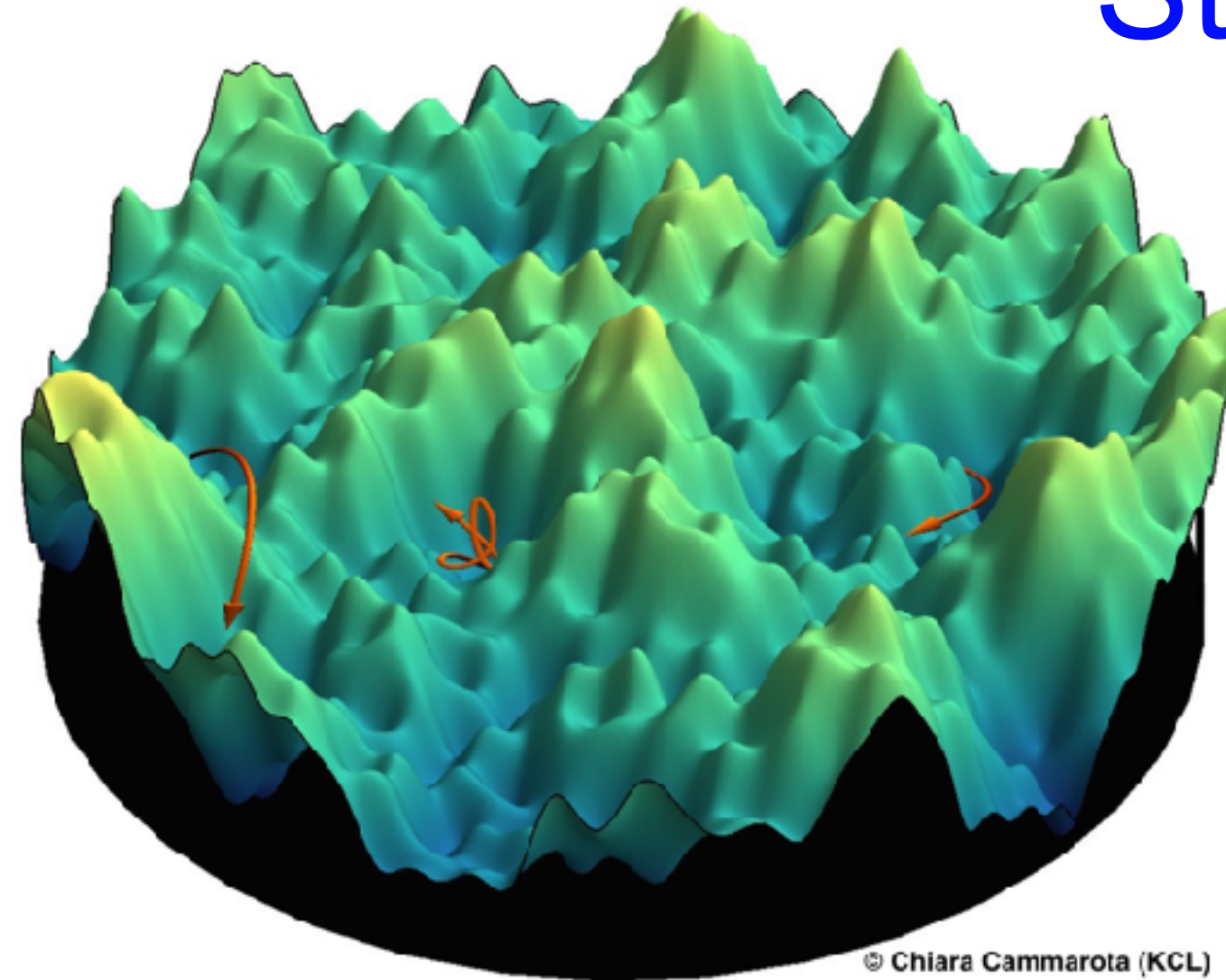
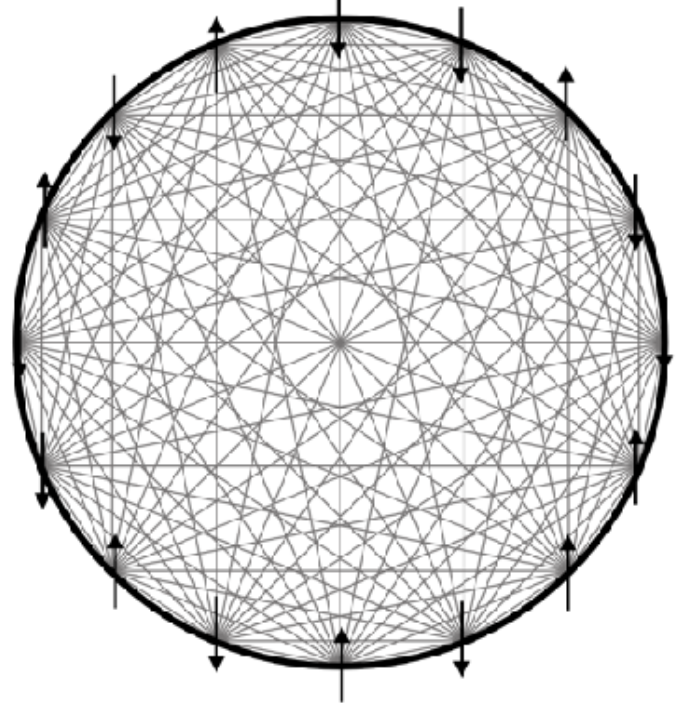


Spin Glass/Hopfield models under asymmetry:

H. Gutfreund, J D Reger¹ and A. P. Young J Phys A (1988),
J. Hopfield PNAS (1982), G. Parisi J Phys A (1986),
A. Crisanti and H. Sompolinsky PRA (1987), L. F. Cugliandolo, et al. PRL (1997) ...

Starting point, XY spin glass (p=0)

- $|J_{ij}| = 1$
- $\mathbb{E}(J_{ij}) = 0$
- $J_{ij} = J_{ji}$

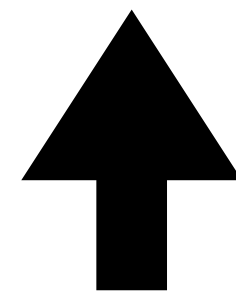


Spin Glass

Shallow and degenerate energy landscape with
high number of marginally stable directions

→ **highly susceptible to perturbations**

→ *Glass destroyed by asymmetry,
leading to chaos*



Difference to SK, p-spin
model... (usually Ising-based)
U(1) symmetry ($\lambda_0 = 0$)

Spin Glass/Hopfield models under asymmetry:

H. Gutfreund, J D Reger¹ and A. P. Young J Phys A (1988),
J. Hopfield PNAS (1982), G. Parisi J Phys A (1986),
A. Crisanti and H. Sompolinsky PRA (1987), L. F. Cugliandolo, et al. PRL (1997) ...

Chaotic
State

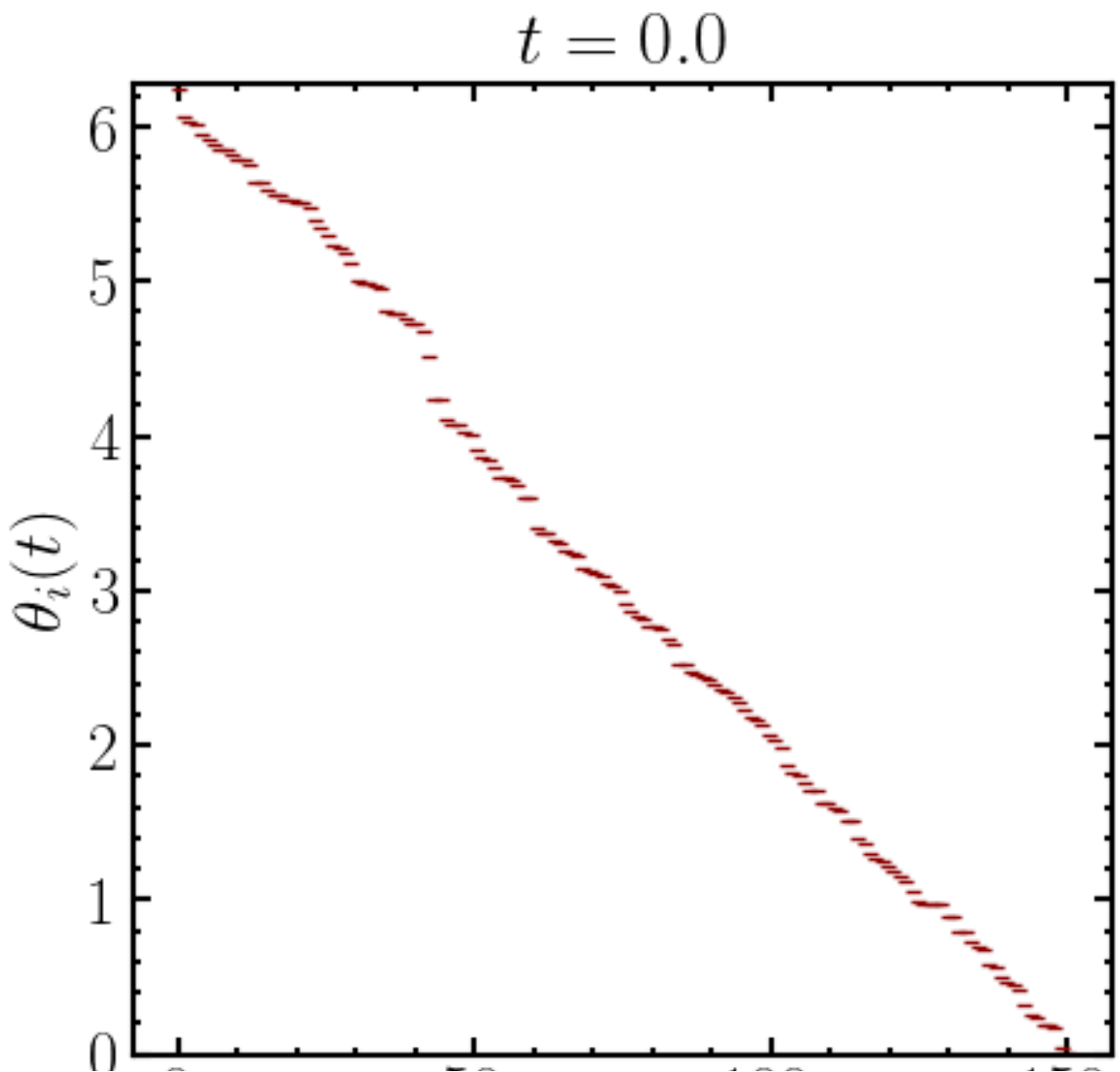
$p=0.3$

Rotating Disordered State

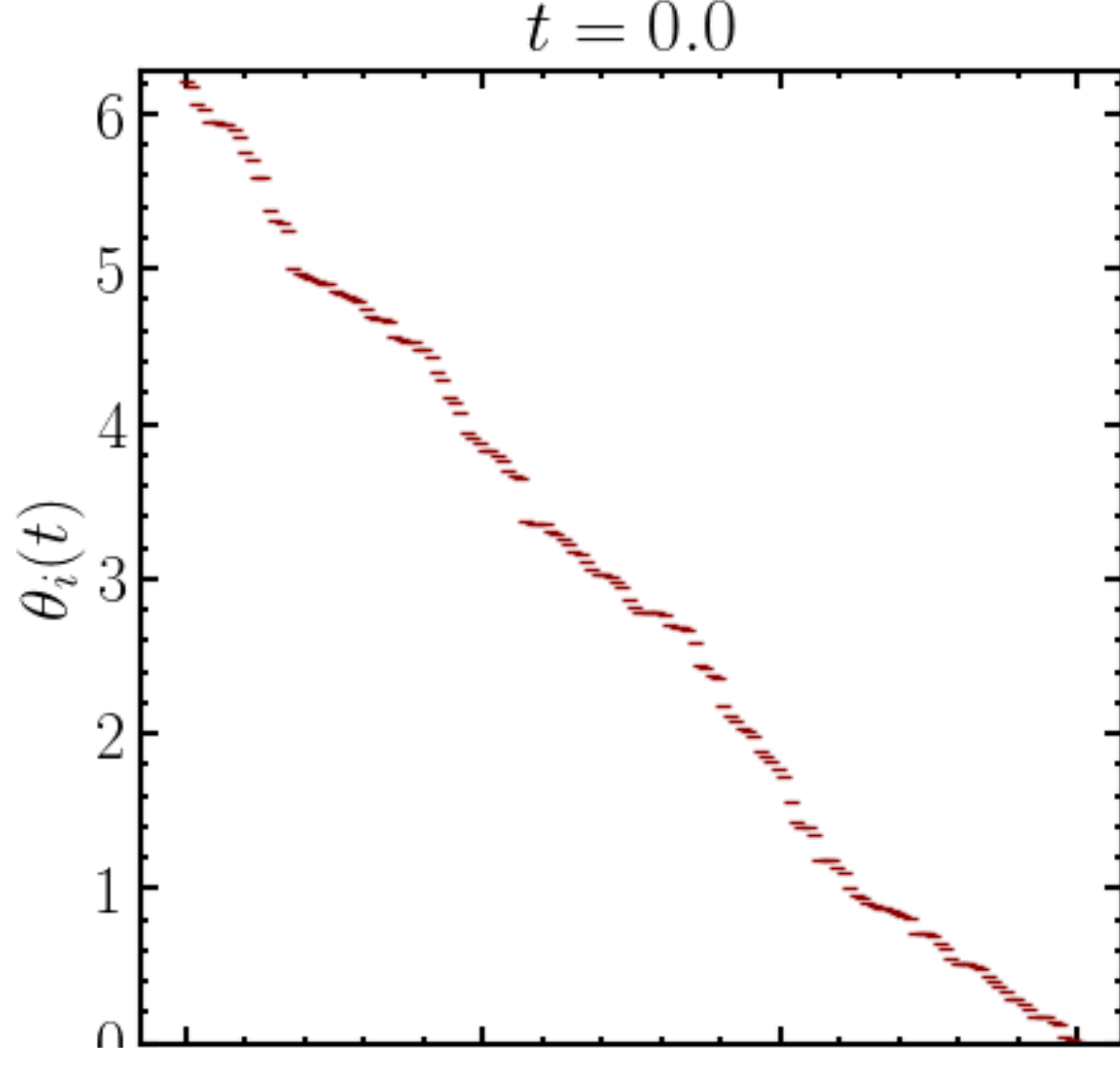
$p=0.05$

Rotating
Disordered
State

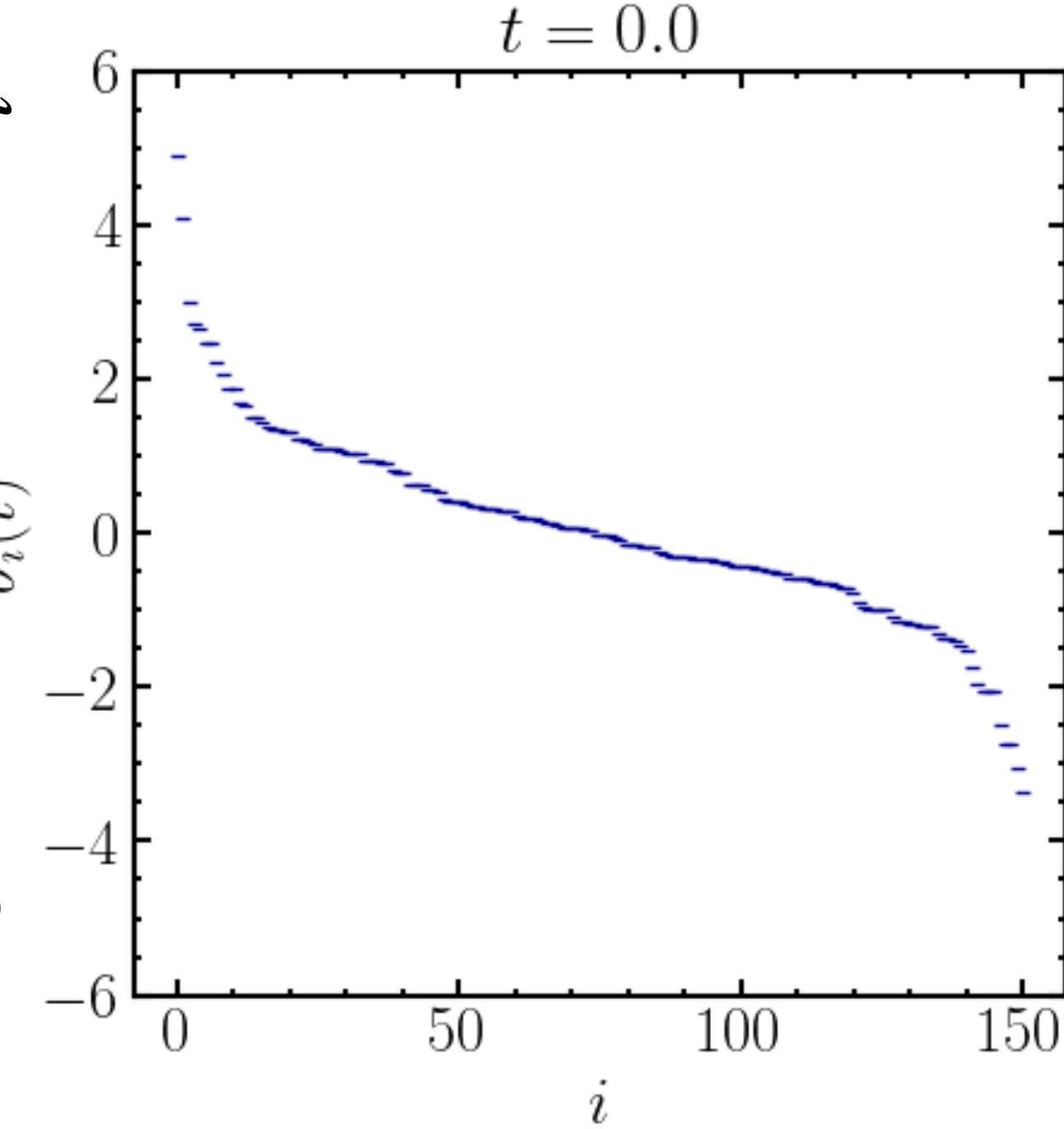
Phases θ_i



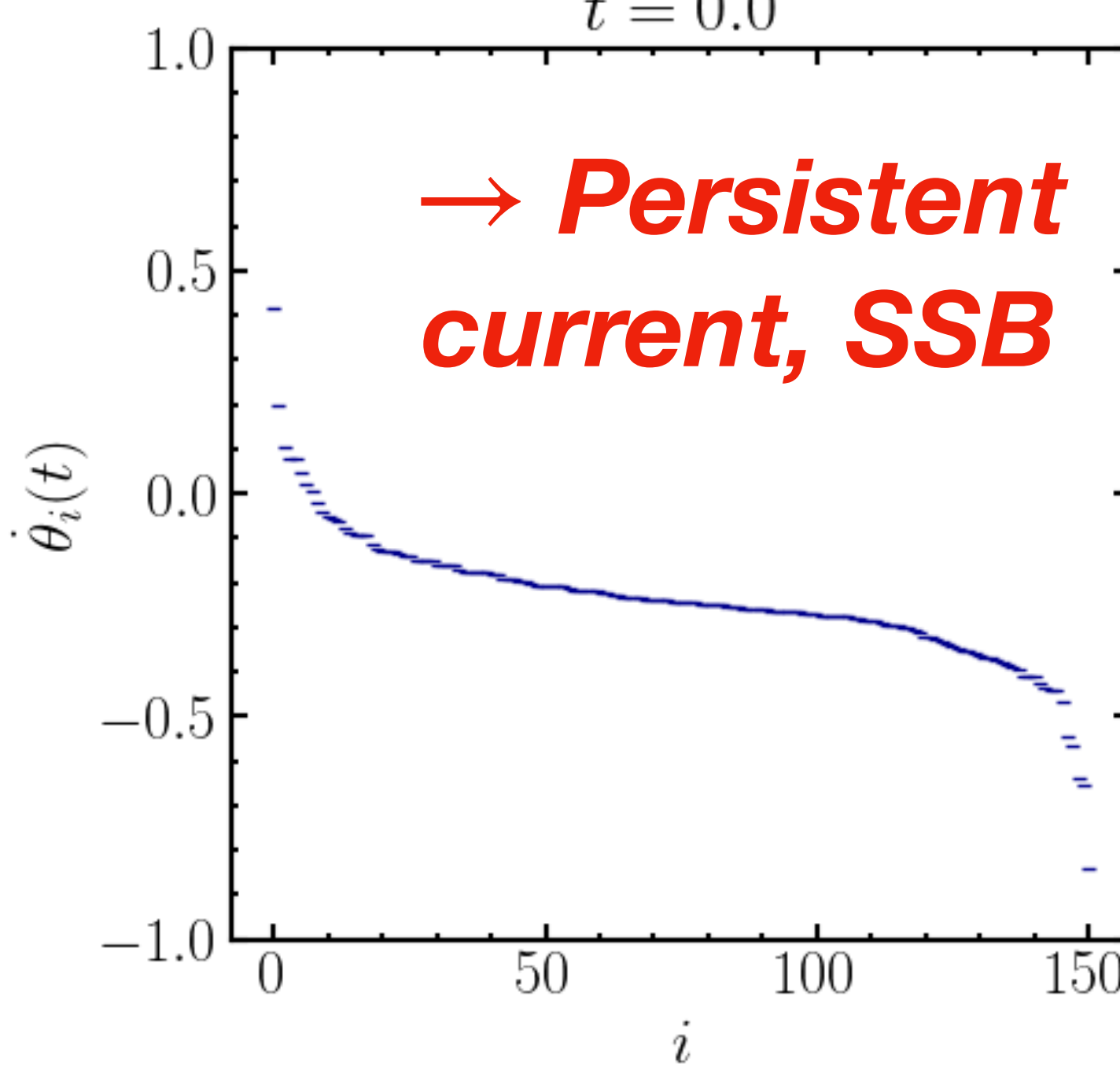
Phases θ_i



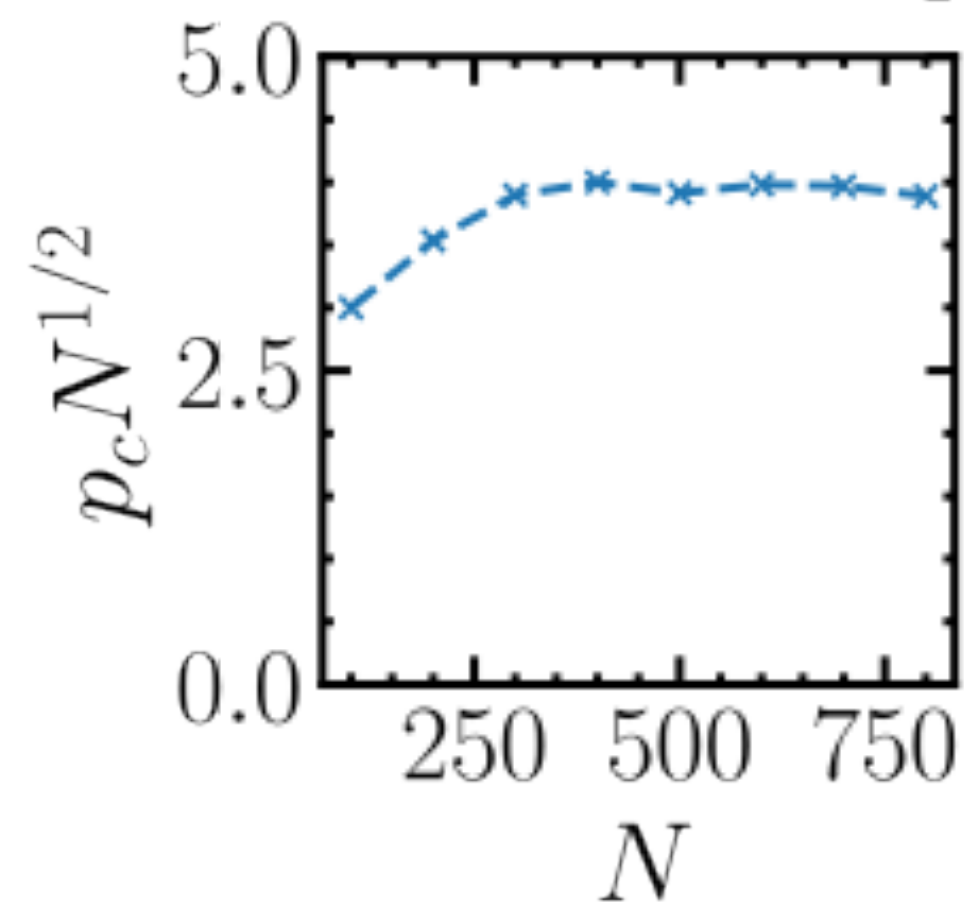
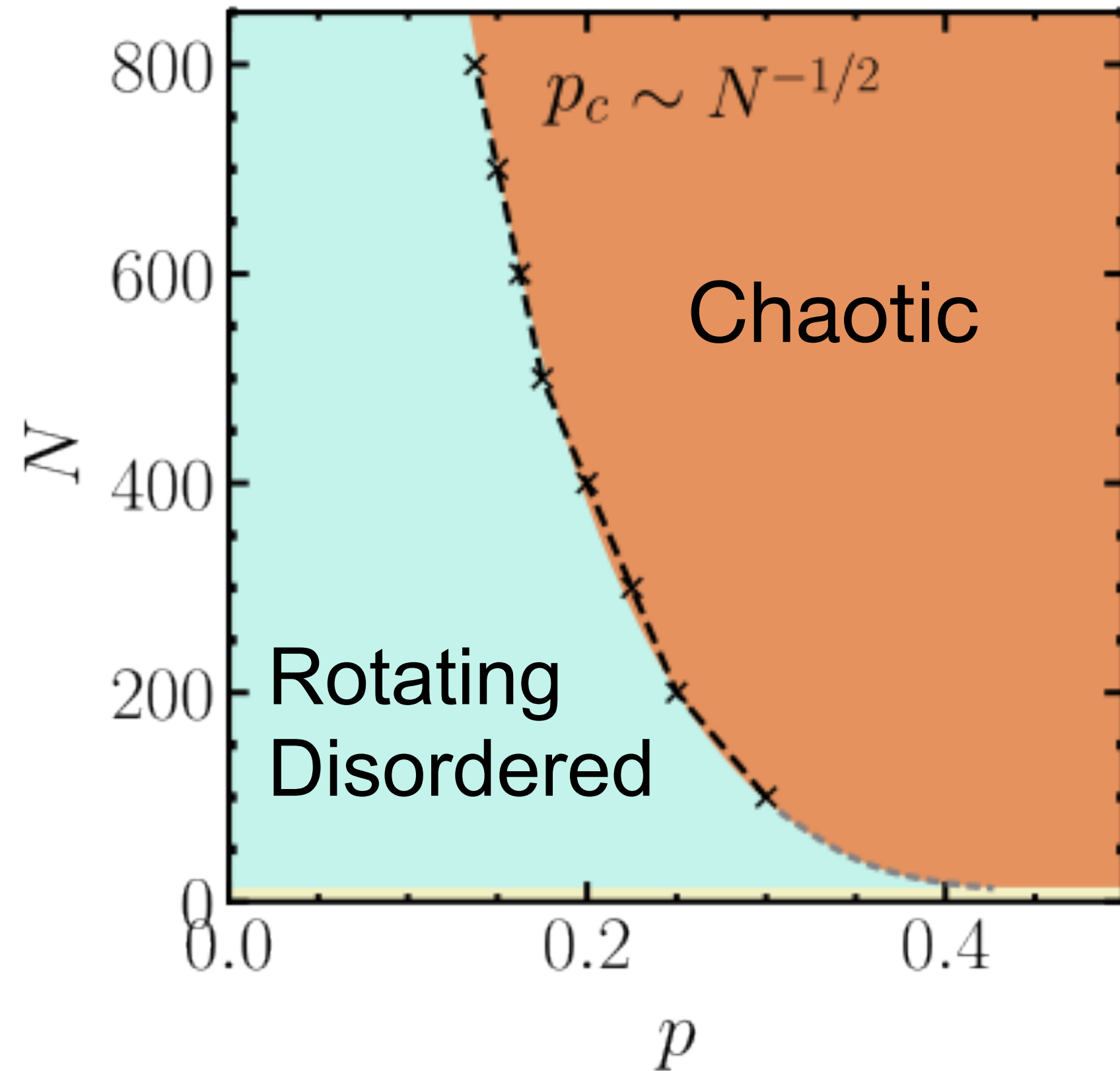
Angular velocities $\dot{\theta}_i$



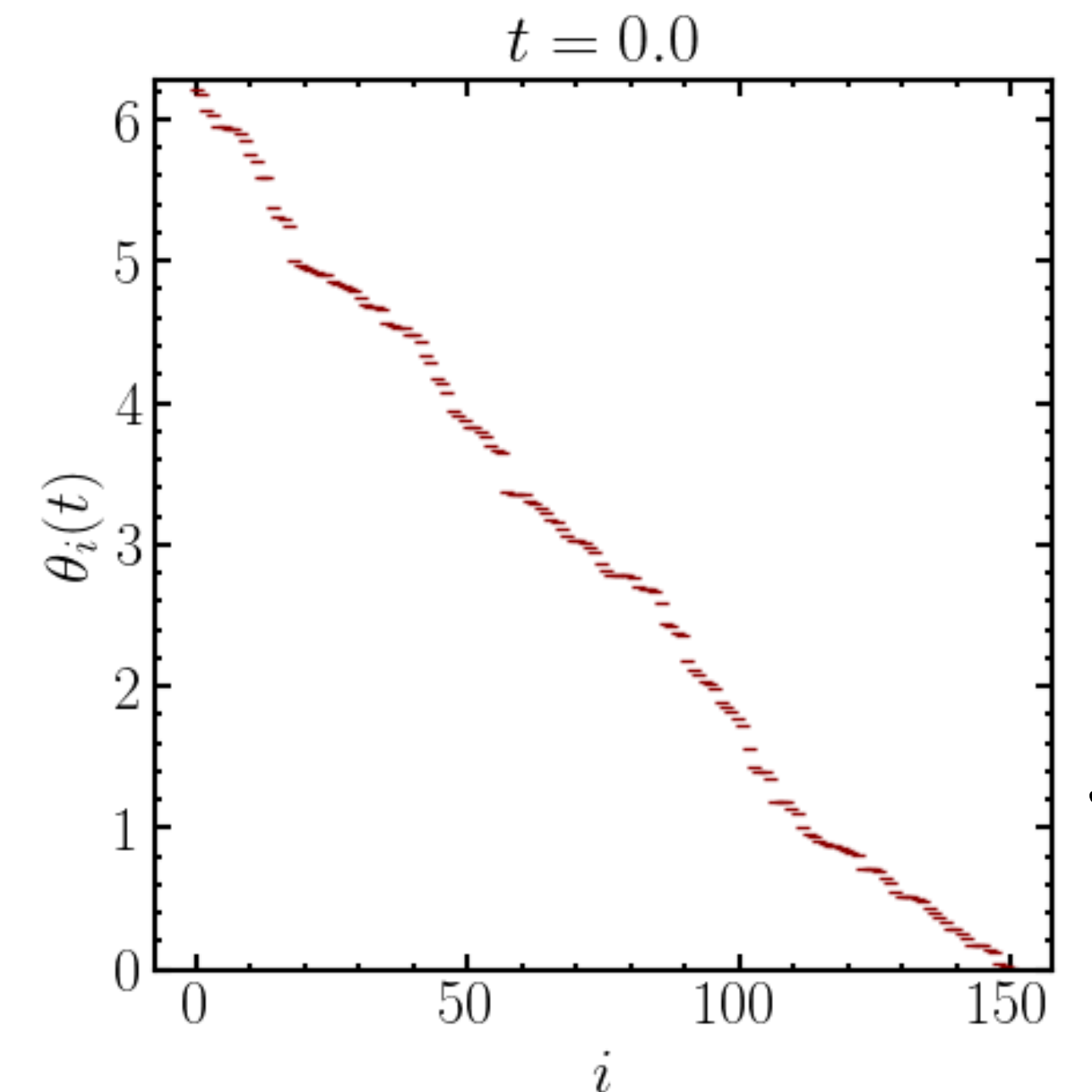
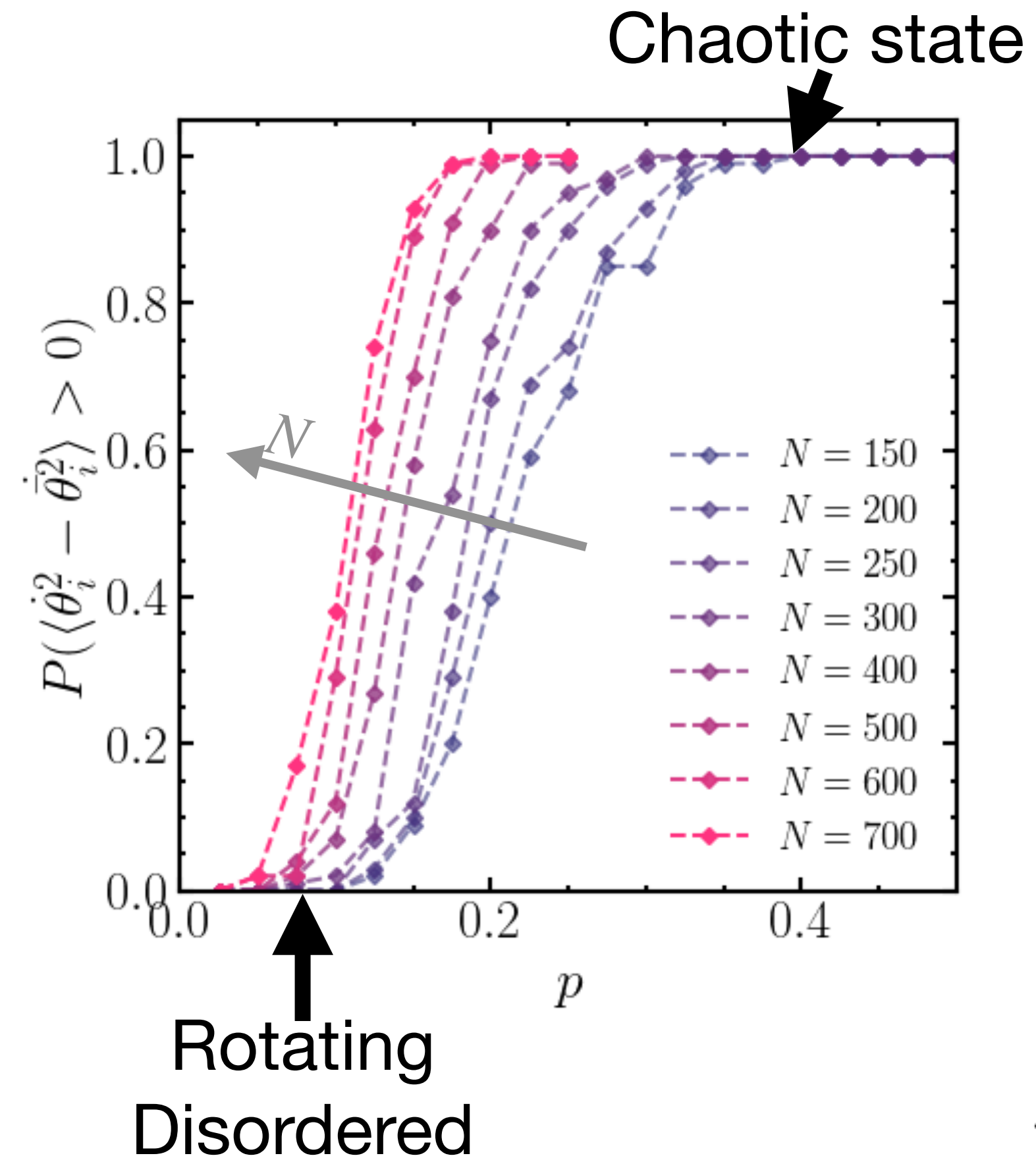
Angular velocities $\dot{\theta}_i$



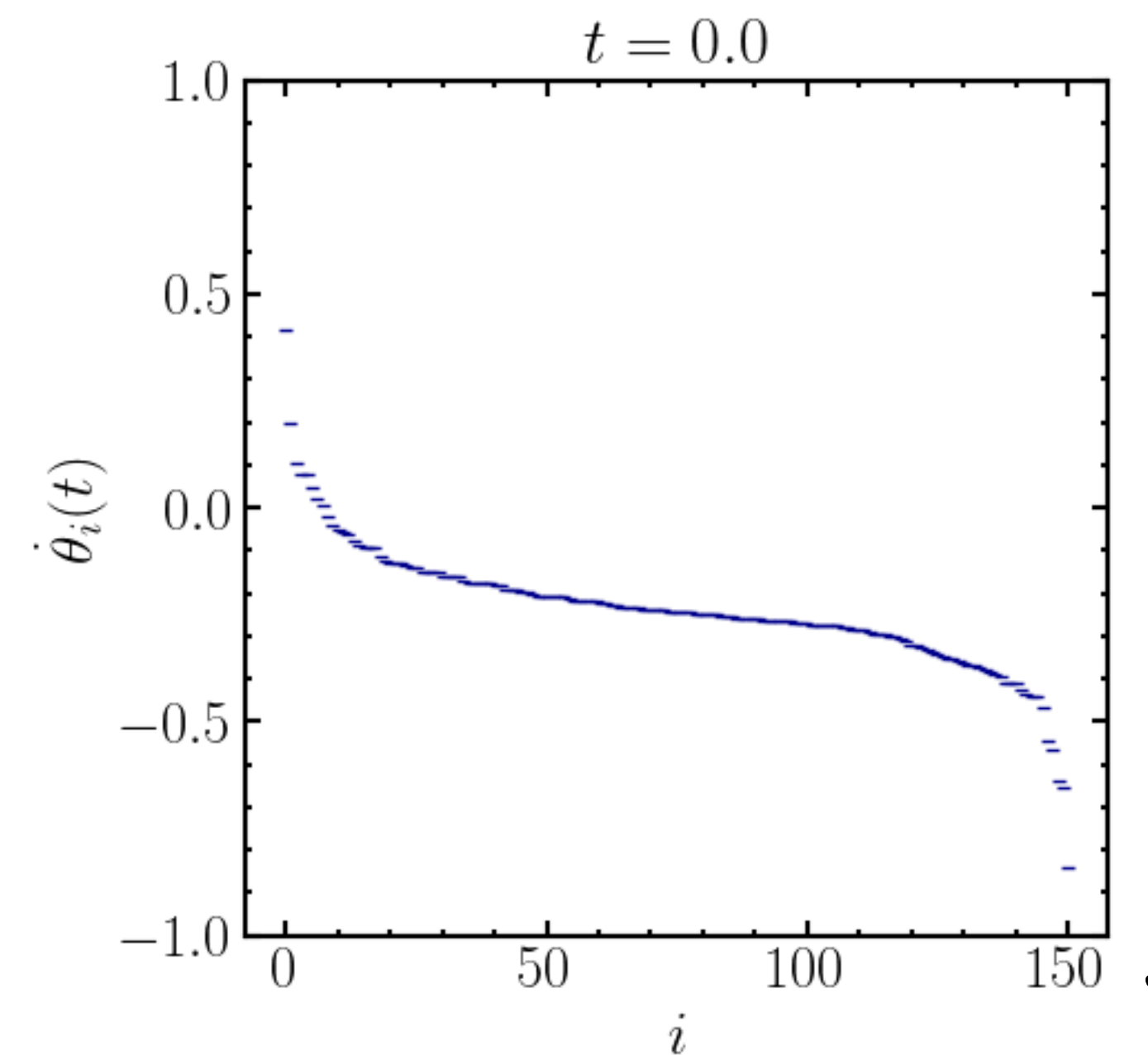
Rotating Disordered State



Disappears for $N \rightarrow \infty$
But slowly $p_c \sim N^{-1/2}$



Phases θ_i



Angular velocities $\dot{\theta}_i$

Thank you to my collaborators!

- **Fluctuating NR field theories:**

Thomas Suchanek (U Leipzig)
Klaus Kroy (U Leipzig)



UNIVERSITÄT
LEIPZIG

Thomas Suchanek



Suchanek, Kroy, & SL, PRL (2023)

Suchanek, Kroy, & SL, PRE 108, 064610 (2023)

Suchanek, Kroy, & SL, PRE 108, 064123 (2023)

Klaus Kroy



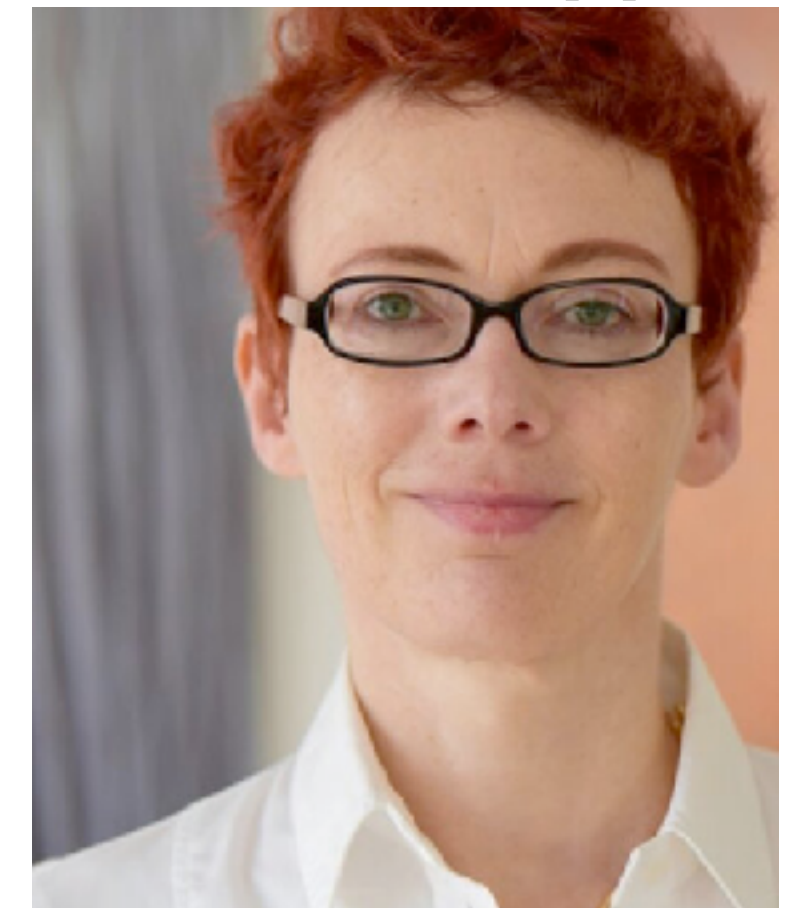
- **Vision-cone NR XY Model:**

Thomas Martynec (then, TU Berlin)
Sabine Klapp (TU Berlin)



SL, Klapp & Martynec, PRL (2023)

Sabine Klapp



- **Disordered NR YX Model:**

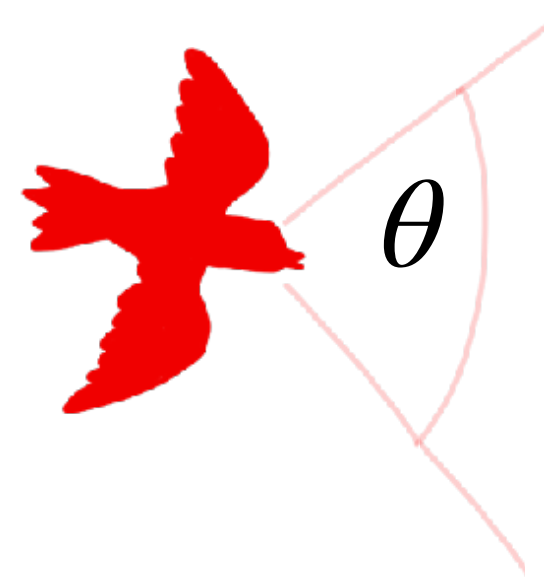
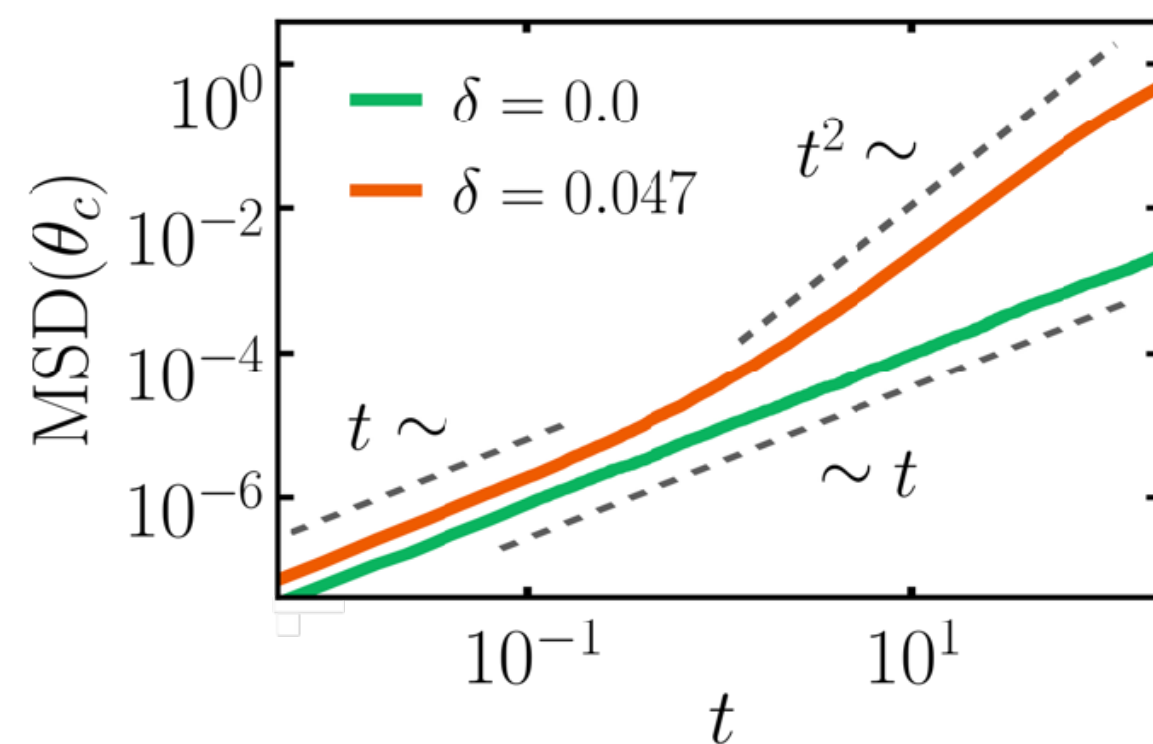
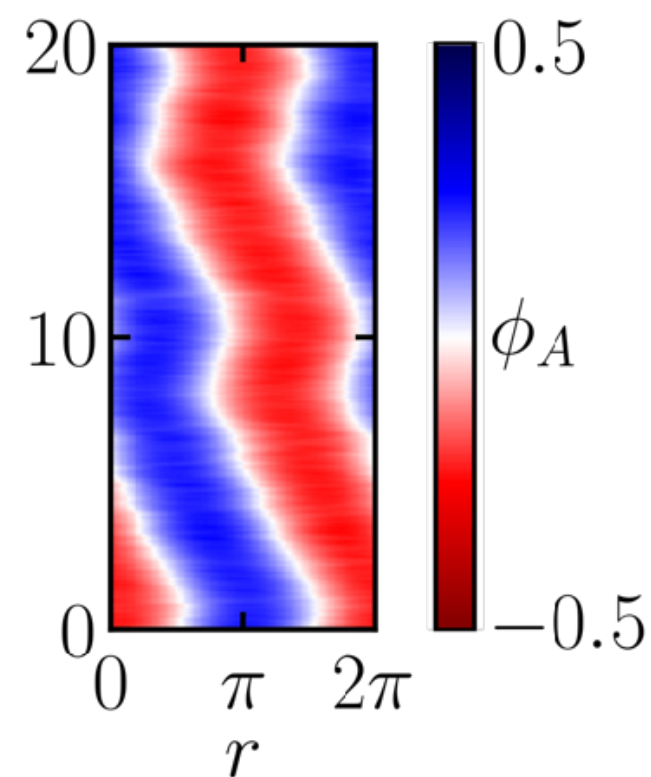
Thomas Suchanek (U Leipzig)
Mike Cates (U Cambridge)
Peter Sollich (U Göttingen)



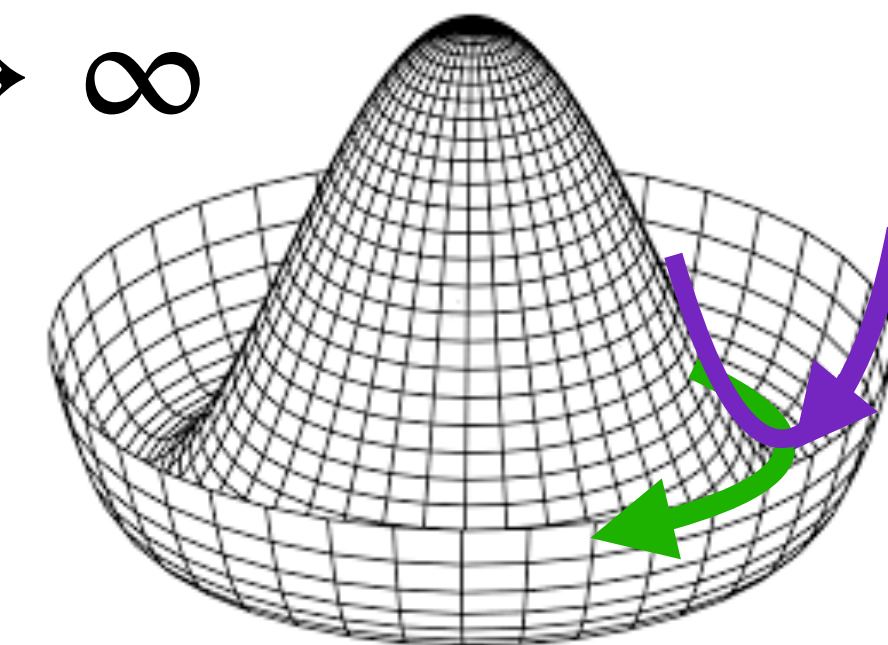
UNIVERSITY OF
CAMBRIDGE



[in preparation]



$$\dot{S} \sim \chi \rightarrow \infty$$



Thank you!

- NR Field theories: Fluctuations close to Critical Exceptional Point+Oscillatory instability
- NRCH: $\dot{S} \sim \chi \rightarrow \infty$ driven by ballistic dynamics of interfaces (coupling to Goldstone mode)
- Vision-cones lead long-range order due to effective pinning to lattice
- Disordered NR system with U(1): Chaos and Rotating disordered state

