

New static black holes with nonspherical horizon topology

(work done together with B. Kleihaus and J. Kunz)

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- Motivation
- The solutions
- Further remarks

based on:

B. Kleihaus, J. Kunz and E. Radu:

- $d \geq 5$ “static black holes with $S^2 \times S^{d-4}$ event horizon topology”, e-Print: arXiv:0904.2723
(Phys. Lett. B: 301, 2009)

Motivation:

–gravity in higher dimensions – a lot of interest in the last decade

–the $d = 5$ black ring solution (Emparan and Reall 2001)

$d > 5$ solutions with nonspherical topology of the horizon?

- the phase structure becomes increasingly intricate and diverse
- no closed form useful solutions (apart from the Myers-Perry black holes)
- approximate solutions: the method of asymptotic expansions (Emparan et. al.); however, limited validity...

numerical (-nonperturbative) approach?

our proposal:

–metric ansatz ($d \geq 5$):

$$ds^2 = -f_0(r, z)dt^2 + f_1(r, z)(dr^2 + dz^2) \\ + f_2(r, z)d\psi^2 + f_3(r, z)d\Omega_{d-4}^2$$

("generalized" Weyl coordinates+rod structure)

–the Einstein equations:

$$\nabla^2 f_0 - \frac{1}{2f_0}(\nabla f_0)^2 + \frac{1}{2f_2}(\nabla f_0) \cdot (\nabla f_2) \\ + \frac{(d-4)}{2f_3}(\nabla f_0) \cdot (\nabla f_3) = 0, \\ \nabla^2 f_1 + \dots = 0, \\ \nabla^2 f_2 + \dots = 0, \\ \nabla^2 f_3 + \dots = 0.$$

no obvious structure except for $d = 5$

Known solutions

the $d = 5$ static black ring:

the metric functions are known in closed form (Emparan and Reall 2001):

$$f_0 = \frac{R_2 + \xi_2}{R_1 + \xi_1},$$
$$f_1 = \frac{(R_1 + \xi_1 + R_2 - \xi_2)((1 - c)R_1 + (1 + c)R_2 + 2cR_3)}{8(1 + c)R_1R_2R_3},$$
$$f_2 = \frac{(R_2 - \xi_2)(R_3 + \xi_3)}{R_1 - \xi_1}, \quad f_3 = R_3 - \xi_3,$$

where

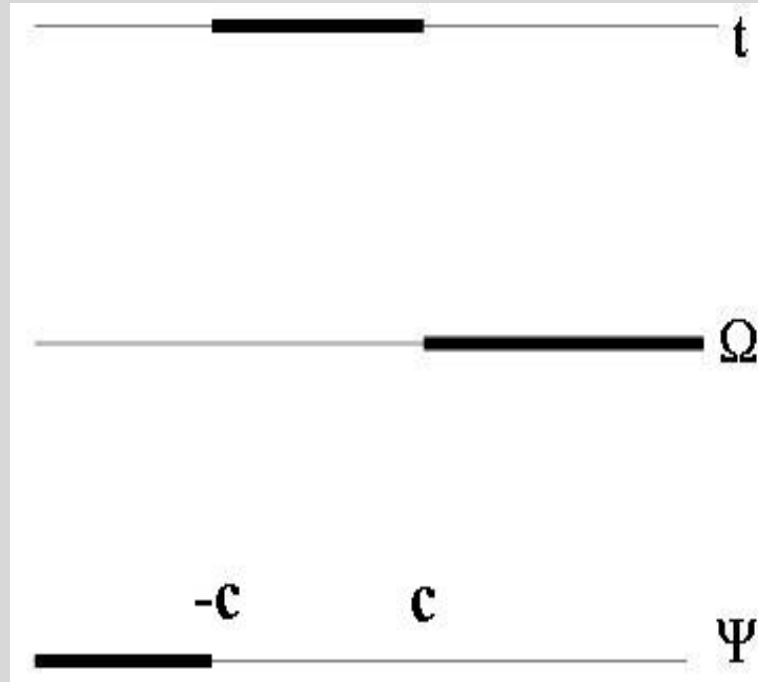
$$\xi_i = z - z_i, \quad R_i = \sqrt{r^2 + \xi_i^2} \quad \text{and} \quad z_1 = -a, \quad z_2 = a, \quad z_3 = b,$$

a and b positive constants, $c = a/b < 1$.

$$M^{(5)}, \quad A_H^{(5)}, \quad T^{(5)} : \text{functions of } (a, b),$$

conical singularity!

$d \geq 5$: the Schwarzschild-Tangerlini solution: S^{d-2} topology of the event horizon



$$f_0 = \left(\frac{v^{(d-3)/2} - 1}{v^{(d-3)/2} + 1} \right)^2, \quad f_1 = c \frac{(v^{(d-3)/2} + 1)^{4/(d-3)}}{4v \left(z^2 \frac{(v^2-1)^2}{(v^2+1)^2} + r^2 \frac{(v^2+1)^2}{(v^2-1)^2} \right)},$$

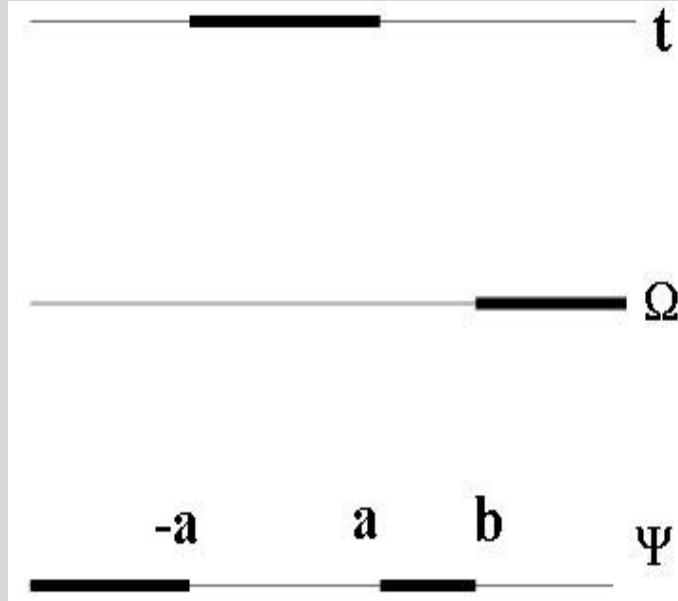
$$f_2 = \dots, \quad f_3 = \dots, \quad \text{with } v = v(r, z)$$

complicated expressions (see arXiv:0904.2723)

New solutions:

Black holes with $S^2 \times S^{d-4}$ topology of the event horizon

–rod structure:



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$d=5$: *static black ring!*

numerical solutions only!

- computation done with the new metric functions F_i
- f_i^0 –the static black ring solution: impose the rod structure

$$f_i = f_i^0 F_i$$

- boundary conditions:

$$\partial_r F_i|_{r=0} = 0, \quad \text{for } -\infty < z \leq b,$$

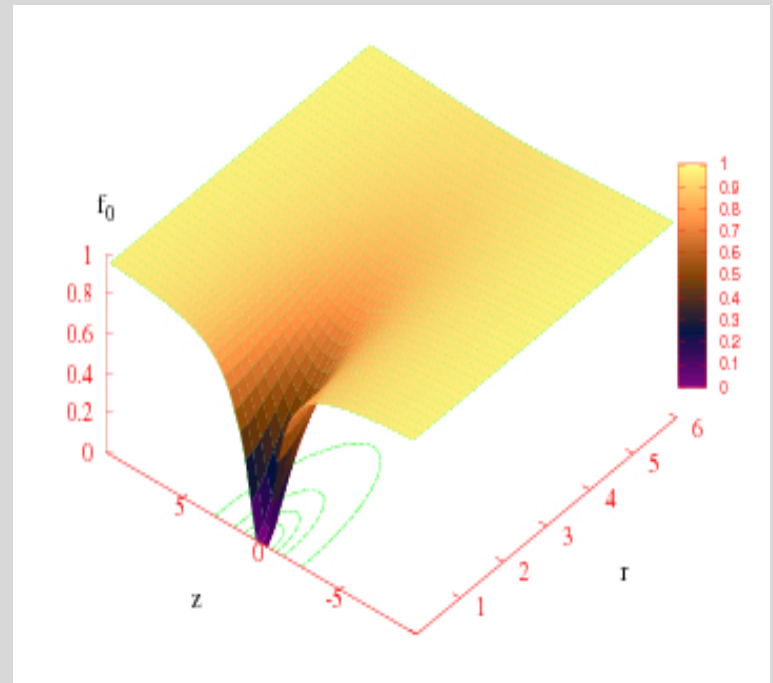
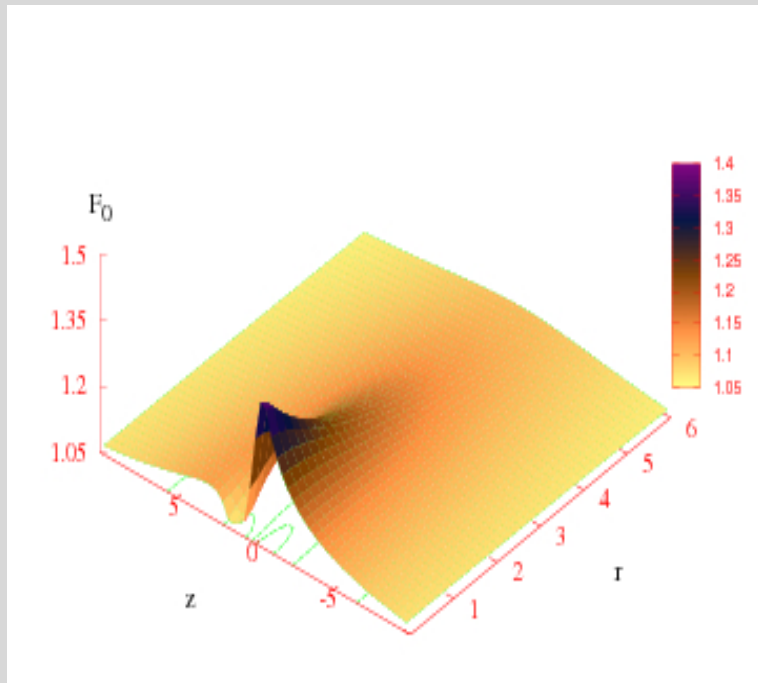
$$\partial_r F_0|_{r=0} = 0, \quad \partial_r F_1|_{r=0} = 0, \quad \partial_r F_2|_{r=0} = 0, \quad F_1|_{r=0} = F_3|_{r=0}$$

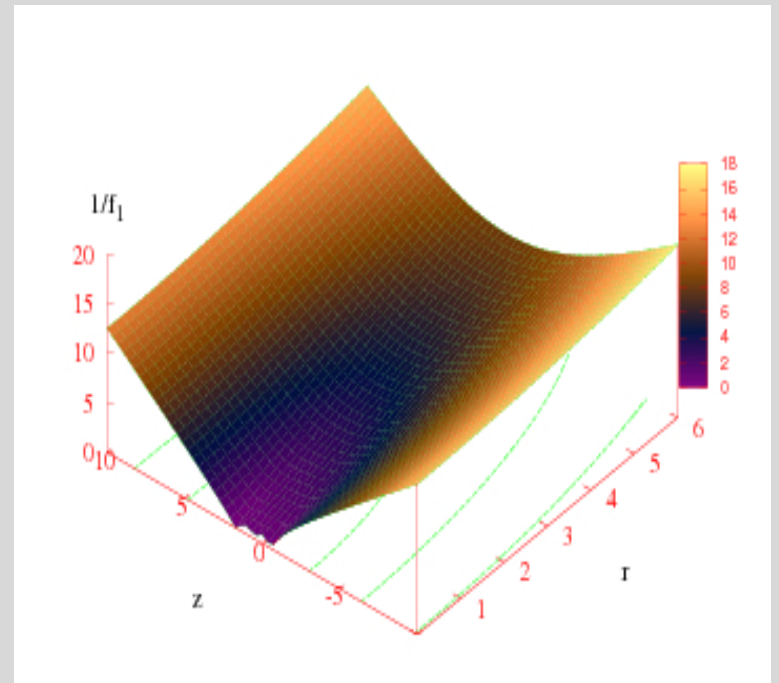
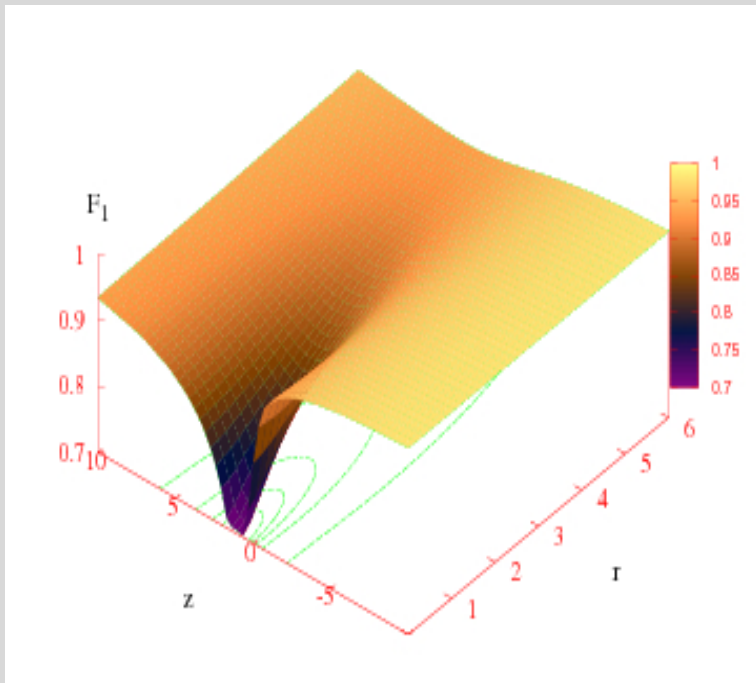
and $F_i = 1$ as $r \rightarrow \infty$ or $z \rightarrow \pm\infty$.

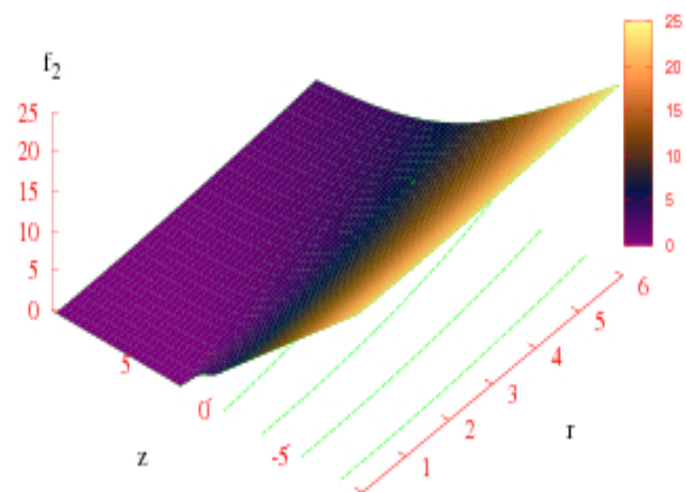
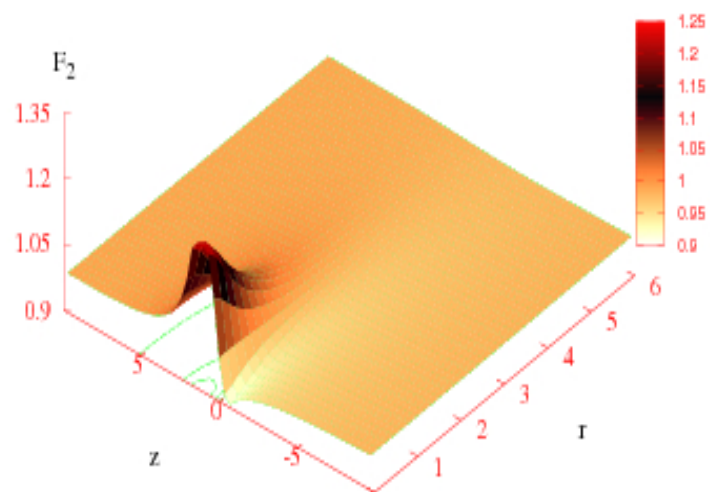
- solve the elliptic partial differential equations for F_i

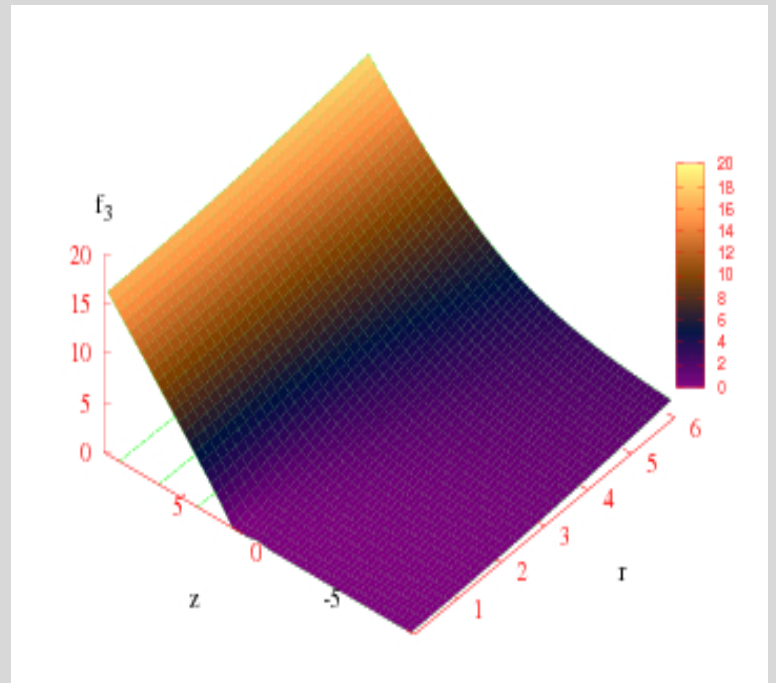
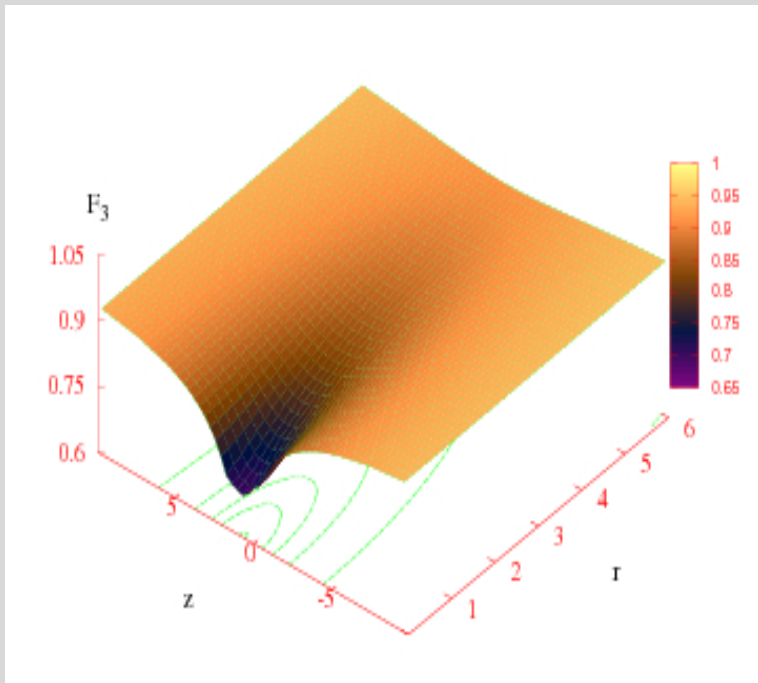
main idea : increase the value of d (dimension of spacetime)

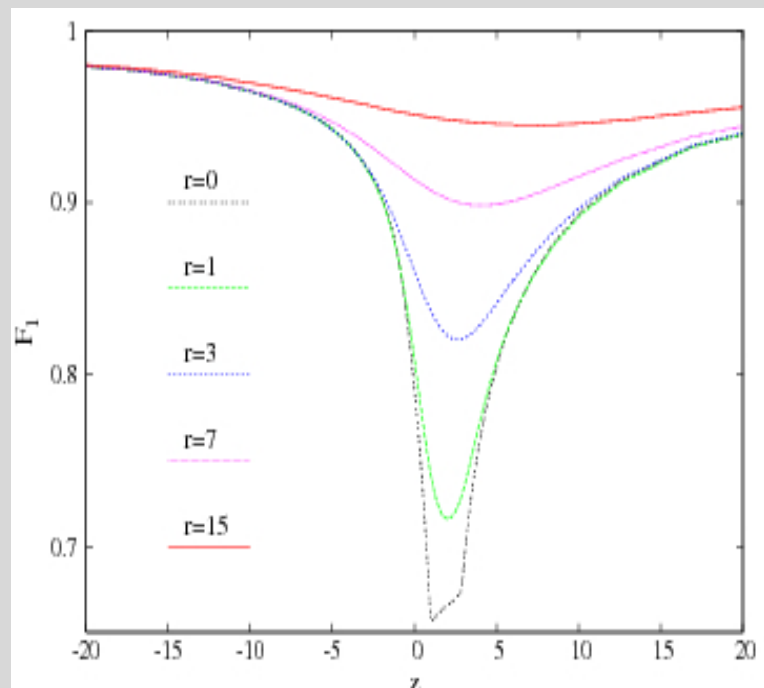
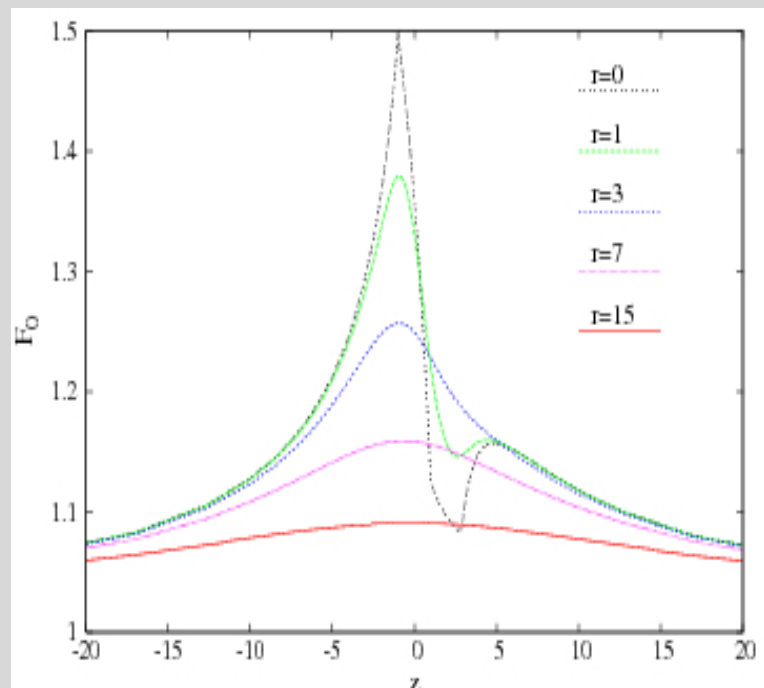
- $d = 6, 7$ solutions with good accuracy
- $d > 7$?

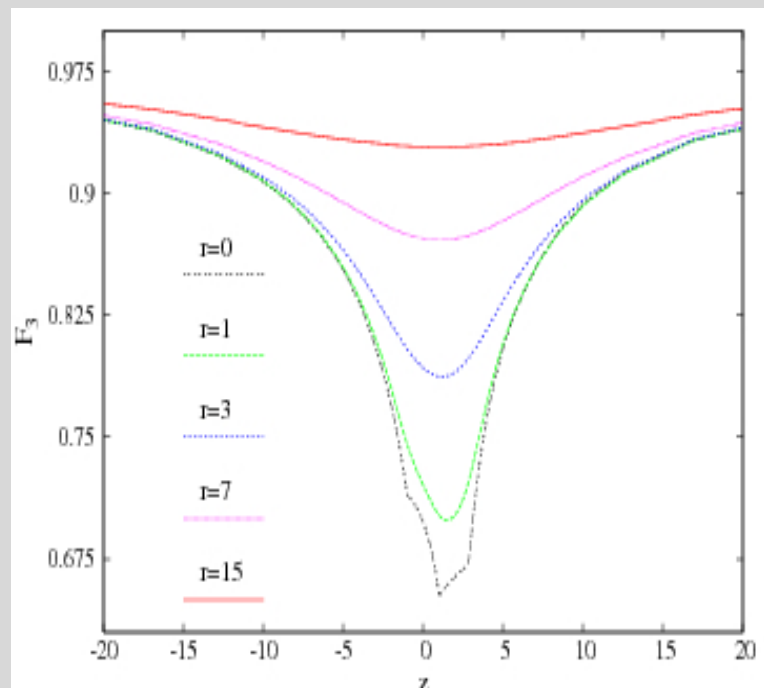
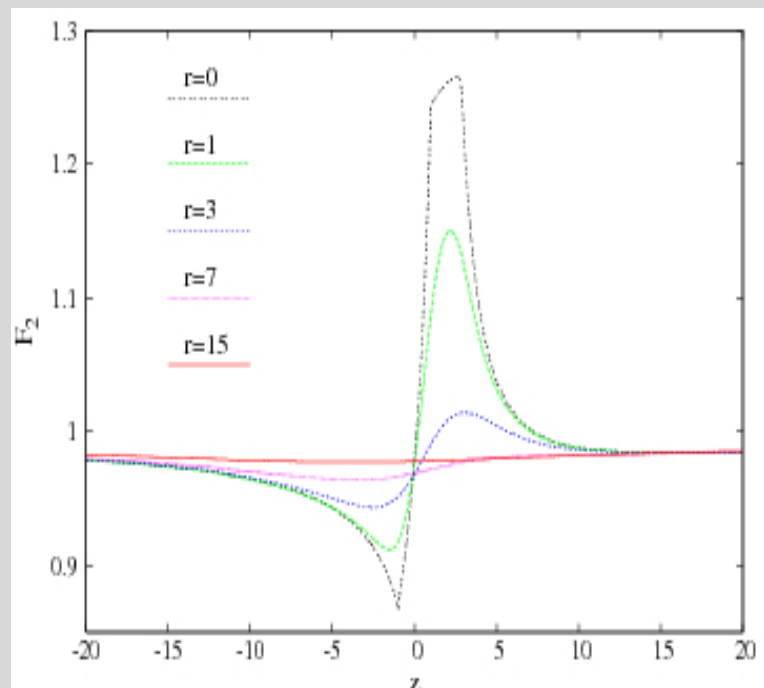












properties of the solutions:

The horizon metric:

$$d\sigma^2 = f_1(0, z)dz^2 + f_2(0, z)d\psi^2 + f_3(0, z)d\Omega_{d-4}^2,$$

The event horizon area:

$$\begin{aligned} A_H &= \Delta\psi V_{d-4} \int_{-a}^a dz \sqrt{f_1 f_2 f_3^{d-4}} \\ &= 2\Delta\psi V_{d-4} \frac{2^{(d-4)/2} a}{\sqrt{a+b}} \int_{-a}^a dz (b-z)^{(d-5)/2} \sqrt{F_1 F_2 F_3^{d-4}} \end{aligned}$$

The Hawking temperature:

$$T = \frac{1}{4\pi a} \sqrt{\frac{a+b}{2}} \sqrt{\frac{F_0}{F_1}},$$

Mass: (from the asymptotics of f_0)

$$f_0 \sim 1 - \frac{16\pi GM}{(d-2)V_{d-2}(r^2+z^2)^{(d-3)/4}} + \dots$$

the Smarr law: **(test of numerics!)**

$$M = \frac{d-2}{4G(d-3)} T A_H .$$

However, all solutions have a conical singularity!

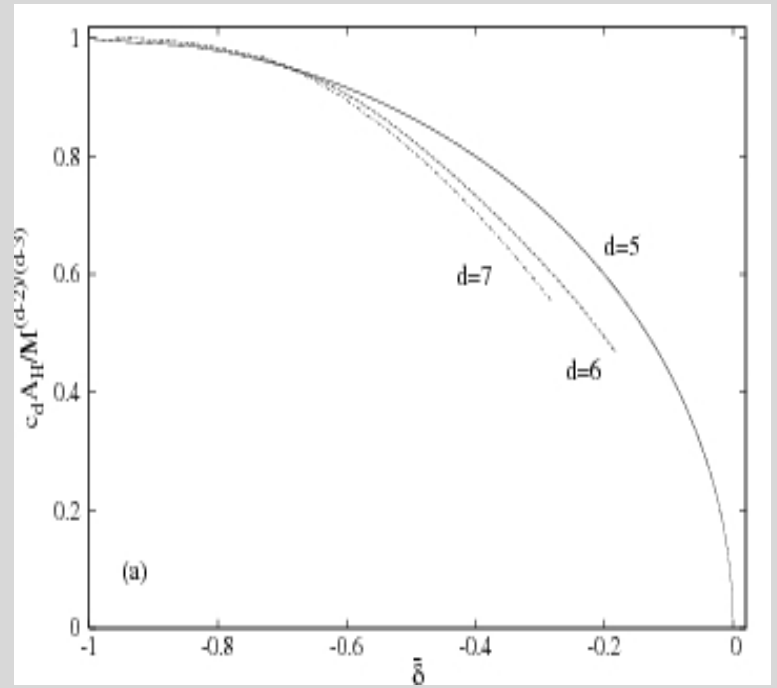
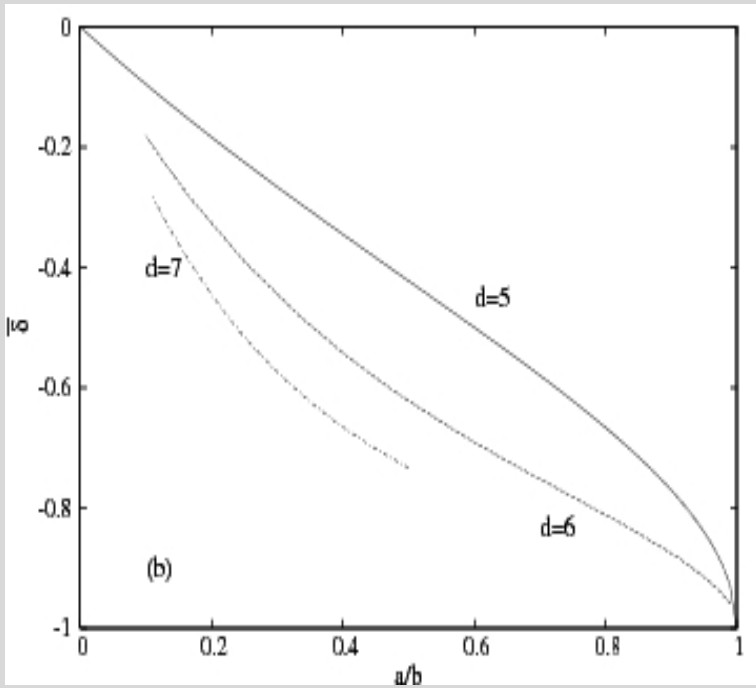
(the same pathology of the $d = 5$ static black ring)

–the value of the conical excess of ψ for $a < z < b$ is

$$\delta = 2\pi \left(1 - \sqrt{\frac{b+a}{b-a}} \sqrt{\frac{F_2}{F_1}} \right).$$

–the 'relative angular excess' ($-1 \leq \bar{\delta} \leq 0$):

$$\bar{\delta} = \frac{\delta/(2\pi)}{1 - \delta/(2\pi)},$$



(left): The relative angular excess $\bar{\delta}$ is shown as a function of the ratio between the two length scales a/b .

(right): The scale free ratio $A_H / M^{(d-2)/(d-3)}$ is shown as a function of the relative angular excess $\bar{\delta}$. The value of the parameter c_d there is $c_d = (d-2)/(16\pi)^{(d-2)/(d-3)} V_{d-2}^{1/(d-3)}$ and has been chosen such that the point $(1, -1)$ on the plot corresponds to the Schwarzschild-Tangerlini black hole.

Further remarks/Conclusions

- Q: Charged static solutions?
- A: Yes! (*however, the same pathology...*)

– electrically charged solutions generated from vacuum via an Harrison transformation

$$I = \frac{1}{16\pi G} \int d^d x \sqrt{-g} \left(R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{4} e^{-2\bar{a}\phi} F^2 \right)$$

– a second Harrison transformation \implies regular solutions in a magnetic Melvin Universe (details in e-Print: arXiv:0904.2723)

situation similar to the $d = 5$ static black rings...

- Q: rotating solutions? $d > 5$ black rings? (within this approach)
- A: very likely. (however, difficult numerical problem (3D eqs., unclear metric ansatz etc) – not impossible...)

- Q: asymptotically (A)dS solutions with this method? (e.g. $d = 5$ AdS black rings)
- A: unlikely

generalizations (currently under study):

- Schwarzschild-Tangerlini + $S^2 \times S^{d-4}$ solution
- new solutions with bubbles