

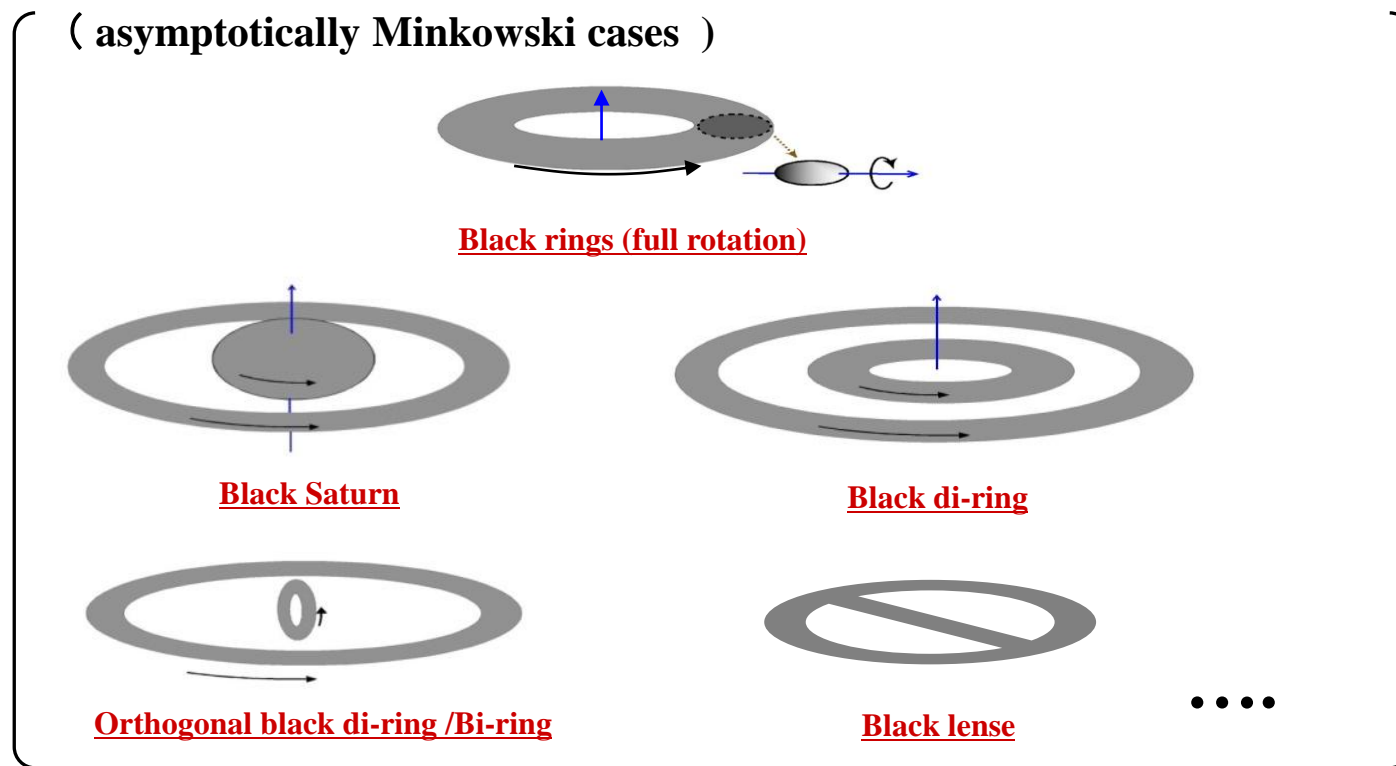
「Existence and some properties of Thermodynamic Black Di-rings」

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**This poster contribution is based on the work in
Phys. Rev. D82, 084009(2010):arXiv:1008.4290v2 [hep-th]**

I. Introduction

- Since the discovery of S^1 rotational black ring by Emparan and Reall several 5-dimensional black hole systems have been obtained using solitonic methods.

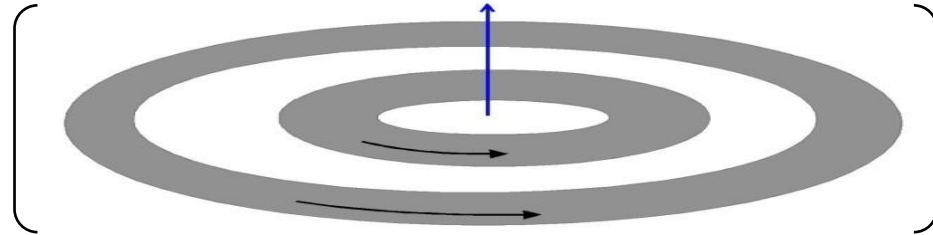


- Previously we succeeded to superimpose two S^1 -rotating black rings in concentric and regular way as simple multi-BH systems. We call the solutions **black di-rings**.

Here we concentrate the study on

black di-rings

(5 dim. concentrically superimposed two S^1 -rotating BRs)



< Two different solution-sets of black di-rings >

(**diring I**)

the Backlund-like transformation.
(Kramer-Neugebauer's -like method)

Iguchi & Mishima(**I&M**): hep-th/0701043
Phys. Rev. D75, 064018 (2007)

(**diring II**)

Inverse Scattering Method (ISM)
(Belinsky-Zakharov technique)

Evslin & Krishnan(**E&K**): hep-th/0706.1231
CQG26:125018(2009)

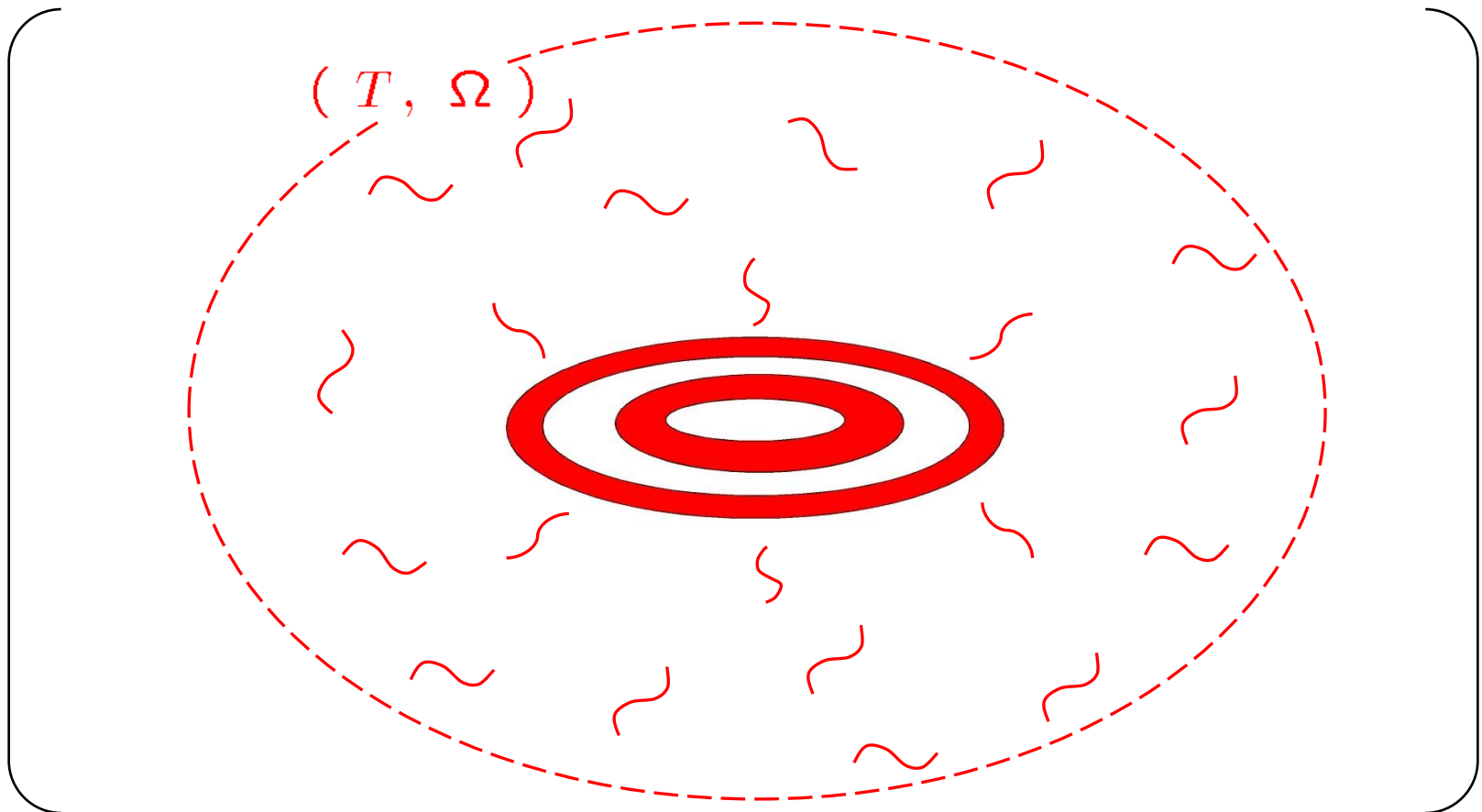
➤ The above two solution-sets may be equivalent, confirmed with the aid of numerical methods. (See Phys. Rev. D82, 084009(2010): arXiv:1008.4290v2[hep-th])

The choice depends on the problem.

< Purpose of this contribution >

We show some physical properties of the di-ring systems,
especially the existence and properties of **thermodynamic regular
black di-ring systems** (globally defined thermal states of regular black di-rings).

TBD



II. Existence of Regular thermodynamic black di-rings

Here we omit the explanation of the solution-generating methods and the expression of the di-ring solutions, and just show the results.

For more details,
 See Phys. Rev. D75, 064018 (2007) hep-th/0701043 (I&M)
 CQG26:125018 (2009) hep-th/0706.1231 (E&K)

First some physical quantities are normalized using the ADM mass M .

$$j^2 = \frac{27\pi}{32GM^3} J^2, \quad a_h^2 = \frac{27}{256\pi} \frac{A_h^2}{(GM)^3},$$

$$\tau_h^2 = \frac{32\pi}{3} (GM) T_h^2, \quad \omega_h^2 = \frac{8}{3\pi} (GM) \Omega_h^2, \quad \dots$$

**j : total angular mom
 a_h : total horizon area
 τ_h : horizon temperature
 ω_h : horizon angular velocity
 ...**

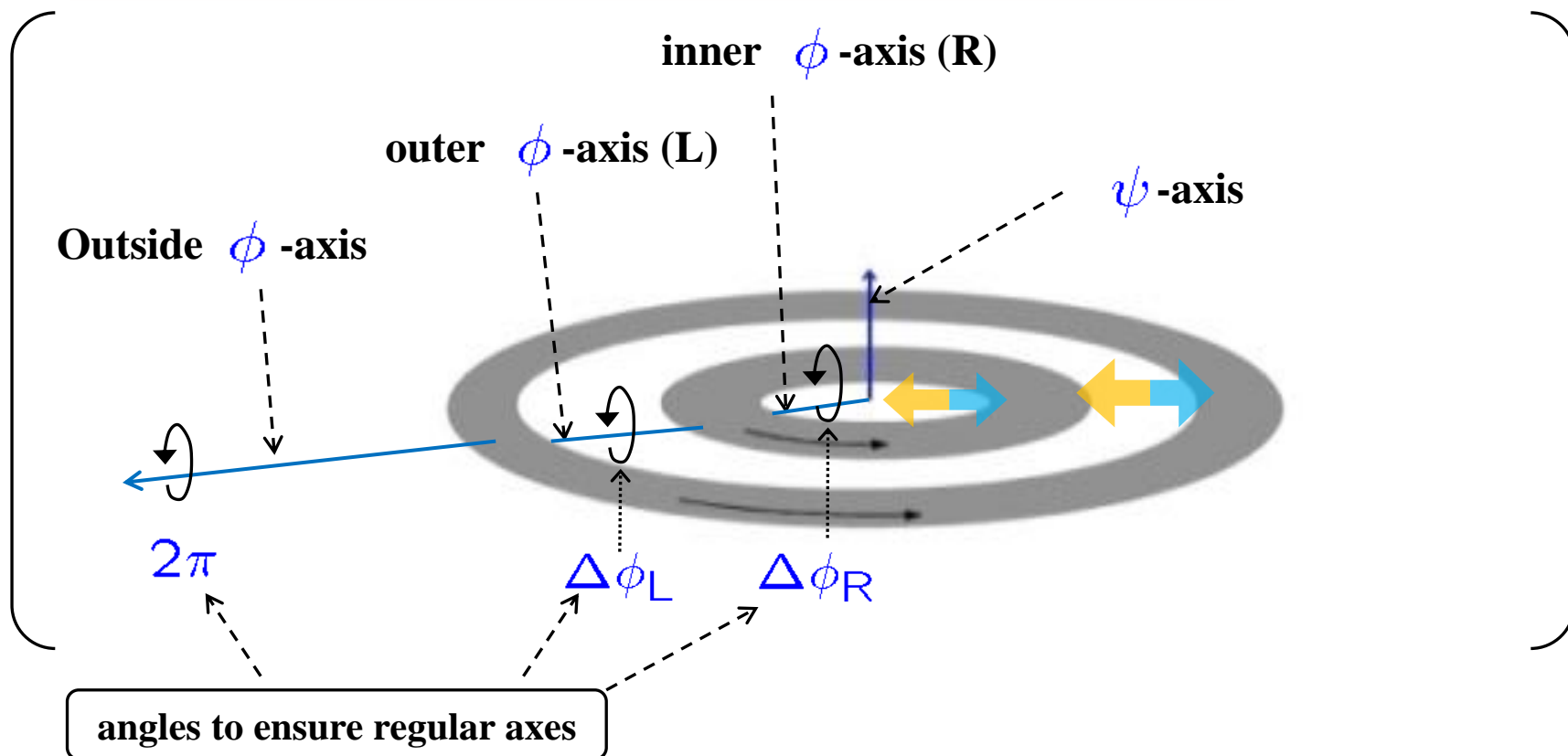
Then the necessary conditions are given in the following.

1. Regularity of black di-rings

The condition of conical singularity-free :

(= the balance condition of gravitational attraction and centrifugal force)
[← yellow arrow] [→ blue arrow]

$$(1) \Delta\phi_L = 2\pi, \quad (2) \Delta\phi_R = 2\pi$$

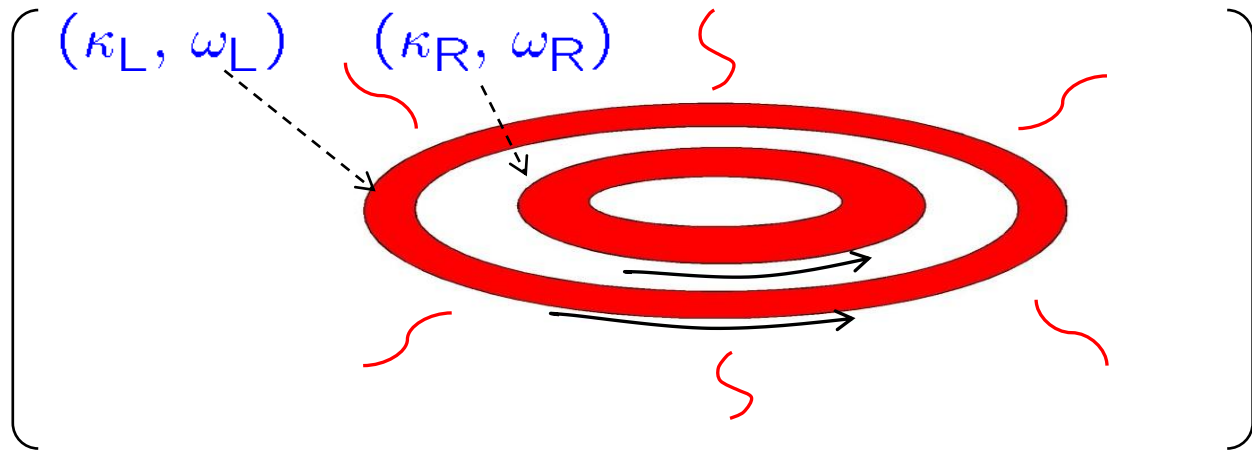


2. Global thermal equilibrium of black dirings

Surface gravity

In the thermal system, a temperature ($T \sim \kappa$) and an angular velocity (ω) must be globally defined as intensive variables,

$$(3) \kappa = \kappa_L = \kappa_R, \quad (4) \omega = \omega_L = \omega_R.$$



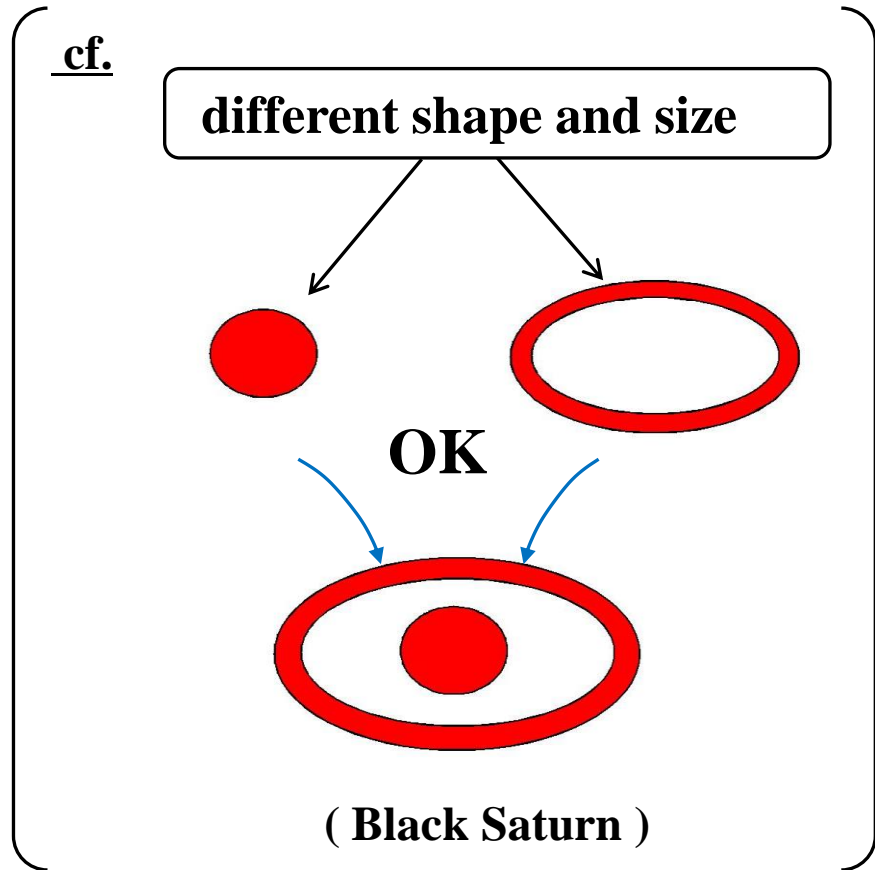
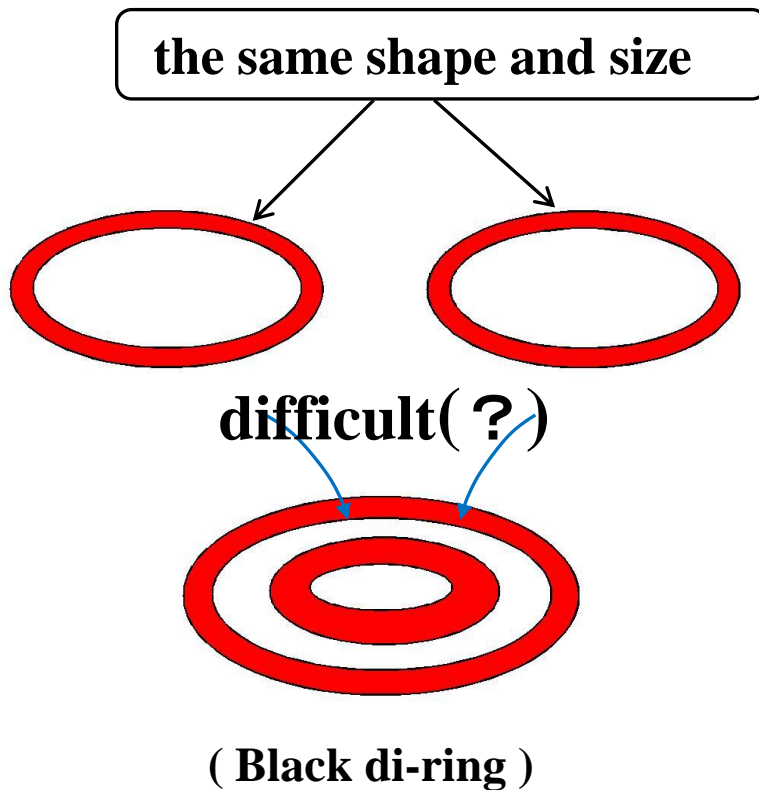
The first important problem is to determine

whether the thermodynamic regular black diring systems exist or not.

< **Existence is not so trivial** >

Some plausible suggestion ... (Elvang, Emparan and Virmani JHEP09 (2008) 003)

If two black rings have the same temperature and angular velocity, ...



If the answer is yes, some mechanism is necessary ... ?

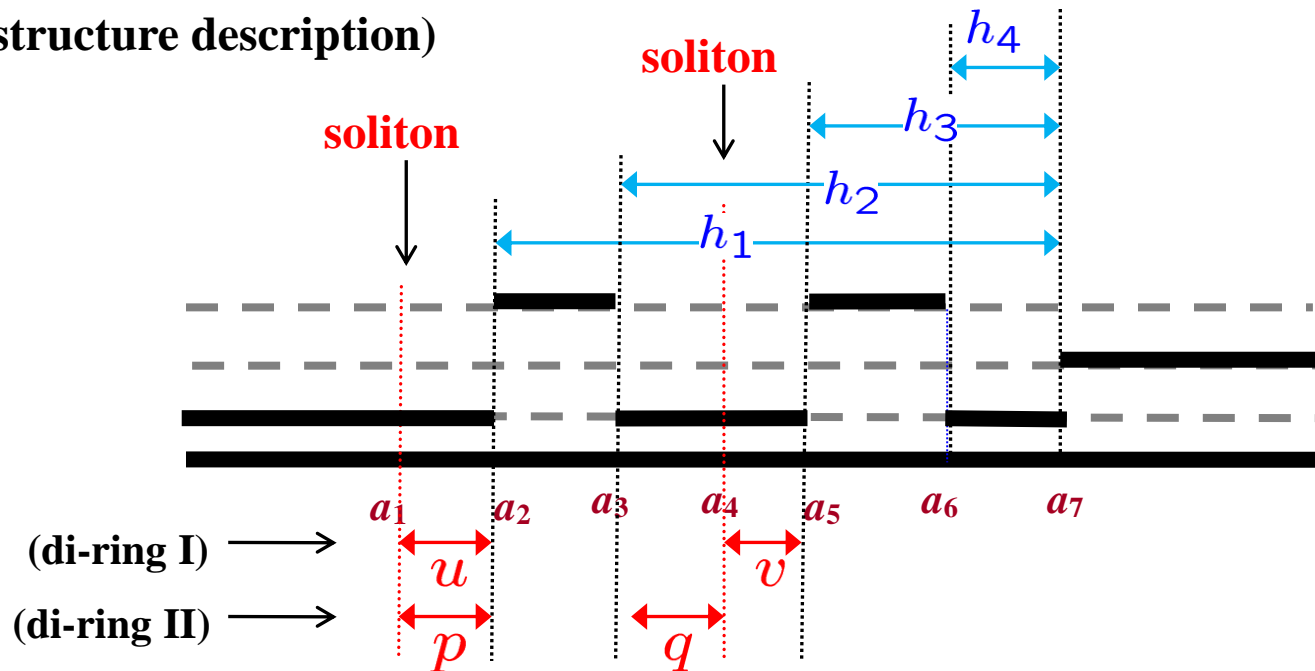
< Remarkable reduction of the constraints for TBD >

After describing the above four constraints from (1) to (4) of thermodynamic regular BD by appropriate six independent moduli-parameters, solve the constraints:

The following six moduli parameters are used.

$$\left[\begin{array}{lll} \underline{h_1} = a_{72}, & \underline{h_3} = a_{75}, & \underline{p} = a_{21} \text{ (or } u = a_{21}), \\ \underline{h_2} = a_{73}, & \underline{h_4} = a_{76}, & \underline{q} = a_{54} \text{ (or } v = a_{43}) \end{array} \right]$$

(Rod structure description)





$$(I) \ \& \ (II) \left\{ \begin{array}{l} \bullet \ 0 = (4h_1^2 - 7h_1h_4 + 4h_4^2)h_3^3 - (8h_1^3 - 8h_1^2h_4 - h_1h_4^2 + 4h_4^3)h_3^2 \\ \quad \quad \quad + (5h_1^4 - 2h_1^3h_4 - 2h_1^2h_4^2 + h_1h_4^3 + h_4^4)h_3 - h_1^5 \\ \bullet \ h_2 = h_1 - h_3 + h_4 \quad (\sim d_1 = d_3) \end{array} \right. \left. \begin{array}{l} \\ \\ \end{array} \right\} \quad \text{(cubic equation for } h_3 \text{)}$$



$$(I) \left\{ \begin{array}{l} \bullet \ u - v = \frac{(h_1 - h_3)(h_2^2 + 2h_1h_3 - h_3^2)}{h_3h_4}, \end{array} \right. \quad \bullet \ (u + h_1)(h_2 - v) = \frac{h_1^2h_2^2}{h_3h_4} \left. \vphantom{\frac{(h_1 - h_3)(h_2^2 + 2h_1h_3 - h_3^2)}{h_3h_4}}} \right\}$$

OR

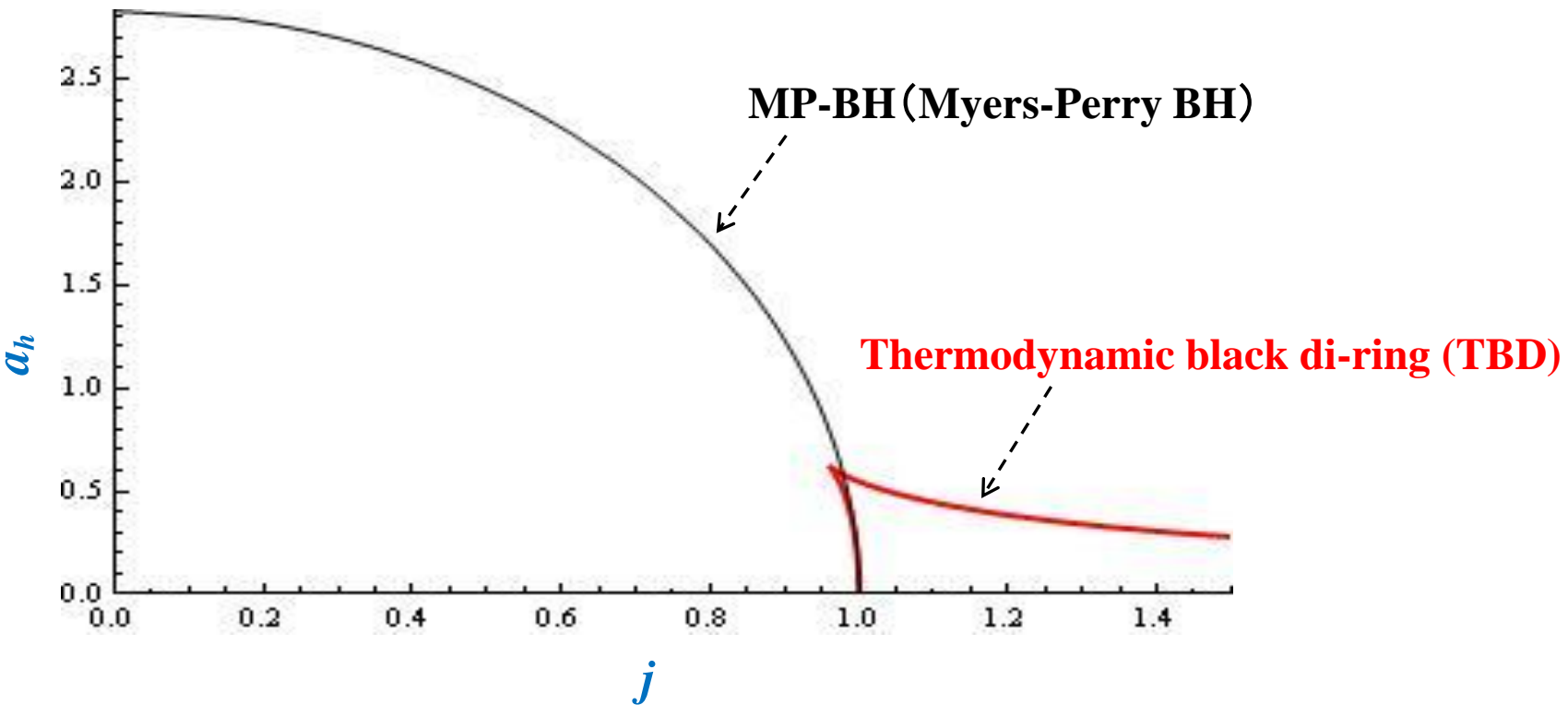
$$(II) \left\{ \begin{array}{l} \bullet \ p + q = \frac{(h_1 - h_2)(h_3^2 + 2h_1h_2 - h_2^2)}{h_2h_4}, \end{array} \right. \quad \bullet \ (p + h_1)(q + h_2) = \frac{h_1^2h_3^2}{h_2h_4} \left. \vphantom{\frac{(h_1 - h_2)(h_3^2 + 2h_1h_2 - h_2^2)}{h_2h_4}}} \right\}$$

Procedure to determine moduli-parameters of TBD

- (i) First h_1 is fixed to be 1 using the arbitrariness of global scaling.
- (ii) h_4 given, h_3 is determined by solving the above cubic equation. Other parameters h_2 and (u, v) or (p, q) are determined by using the other equations.
- (iii) Then we choose the set which satisfies $1 > h_2 > h_3 > h_4 > 0$ as physical one.

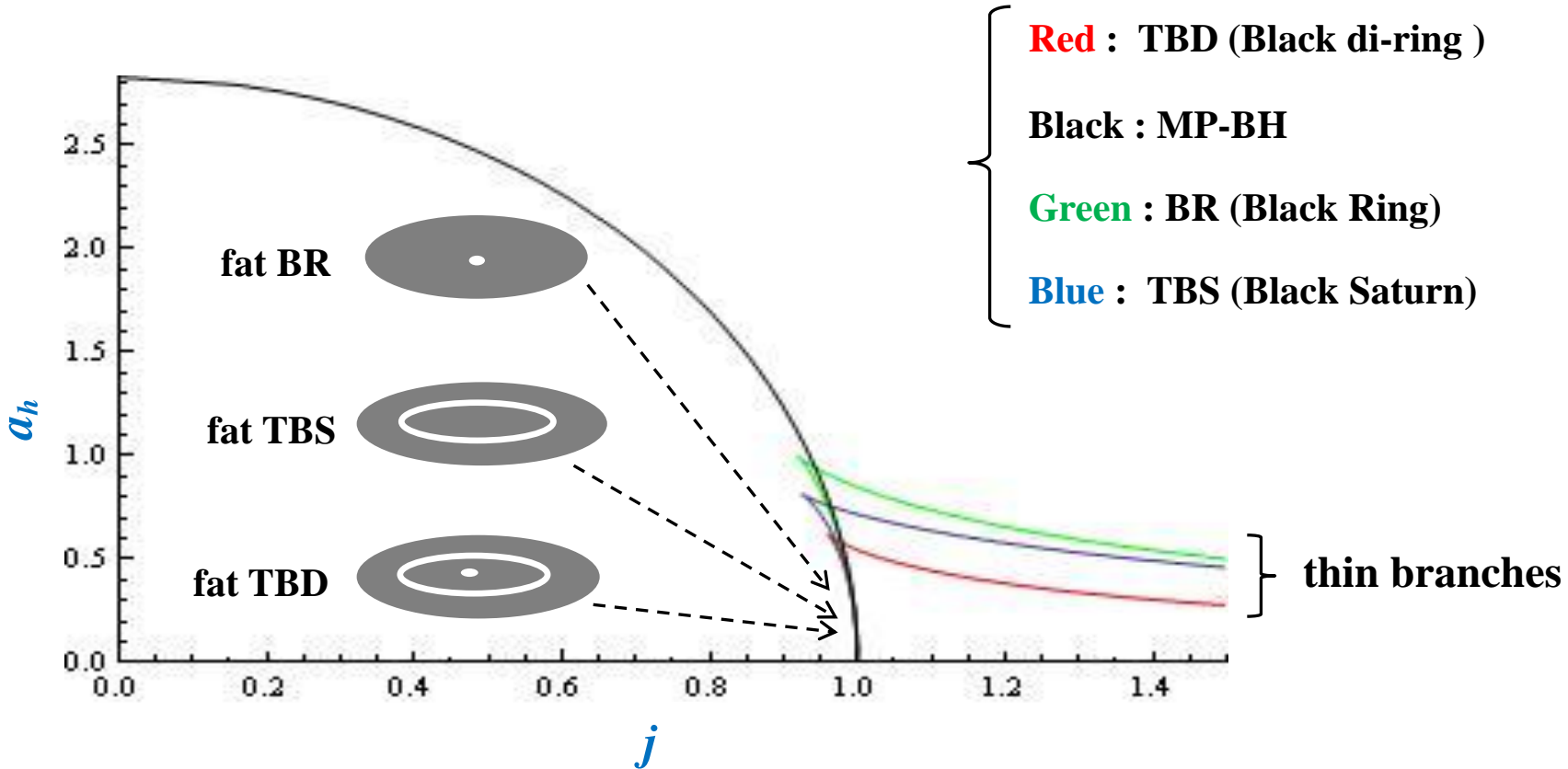
< existence of TBD >

Phase diagram of total areas (a_h) against angular momentum (j)



➤ **Thermodynamic black di-ring systems exist !**

< Comparing thermodynamic black di-ring with other thermodynamic objects >



■ The phase of thermodynamic di-ring has a similar behavior to the BR or TBS system.

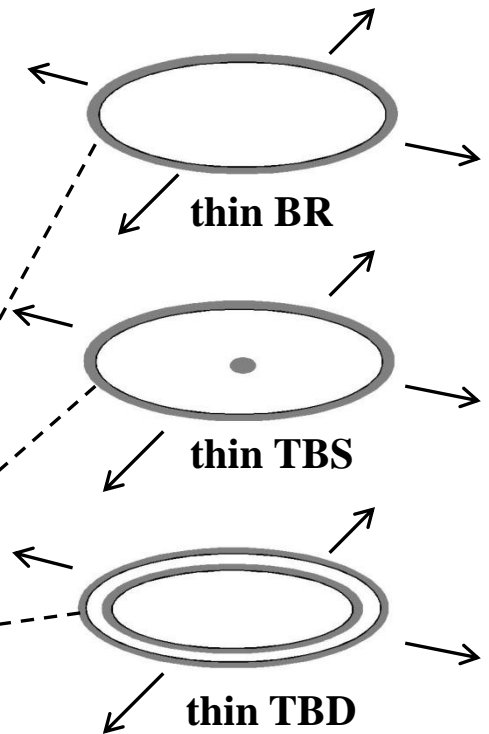
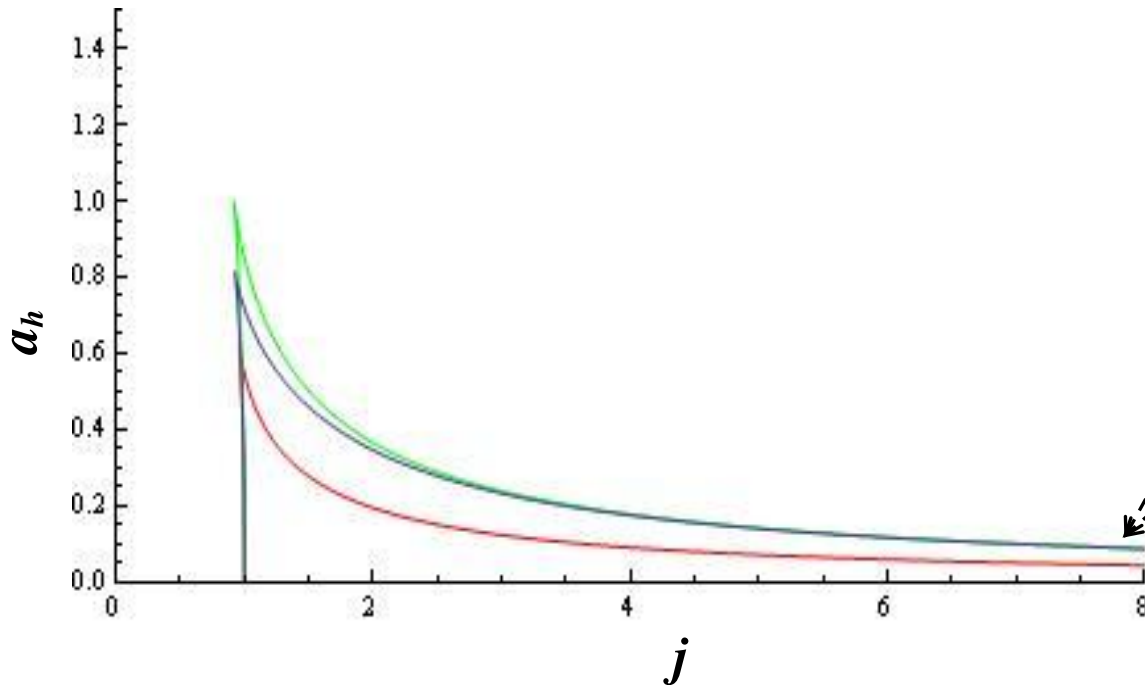


a 'fat ring' branch and a 'thin ring' branch appear.

III. Some peculiarities of thermodynamic black di-rings

To clarify the property of TBD, we contrast TBD with TBS .

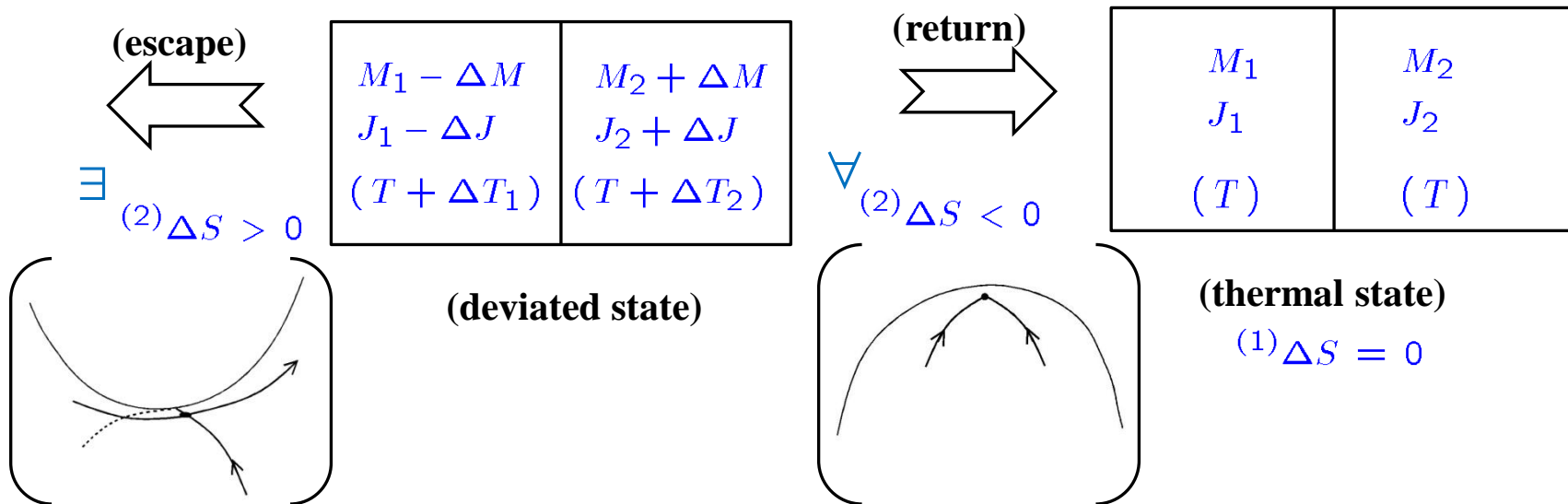
1. Behavior of total areas in the thin branch (large j)



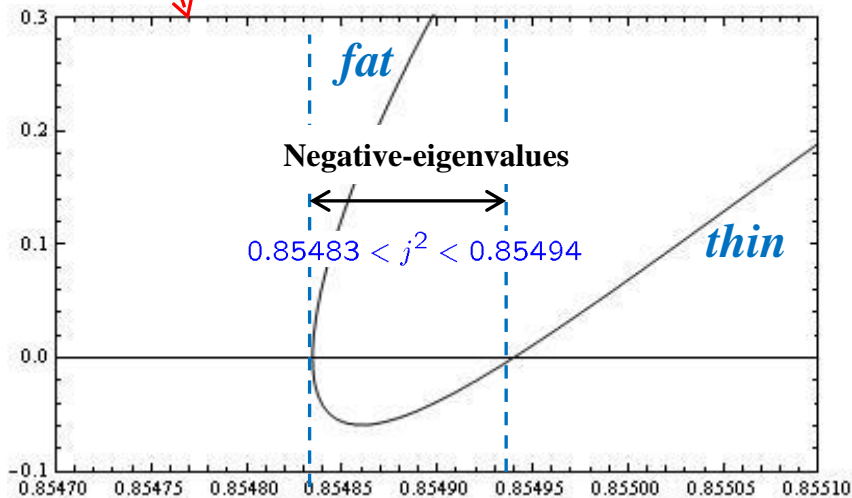
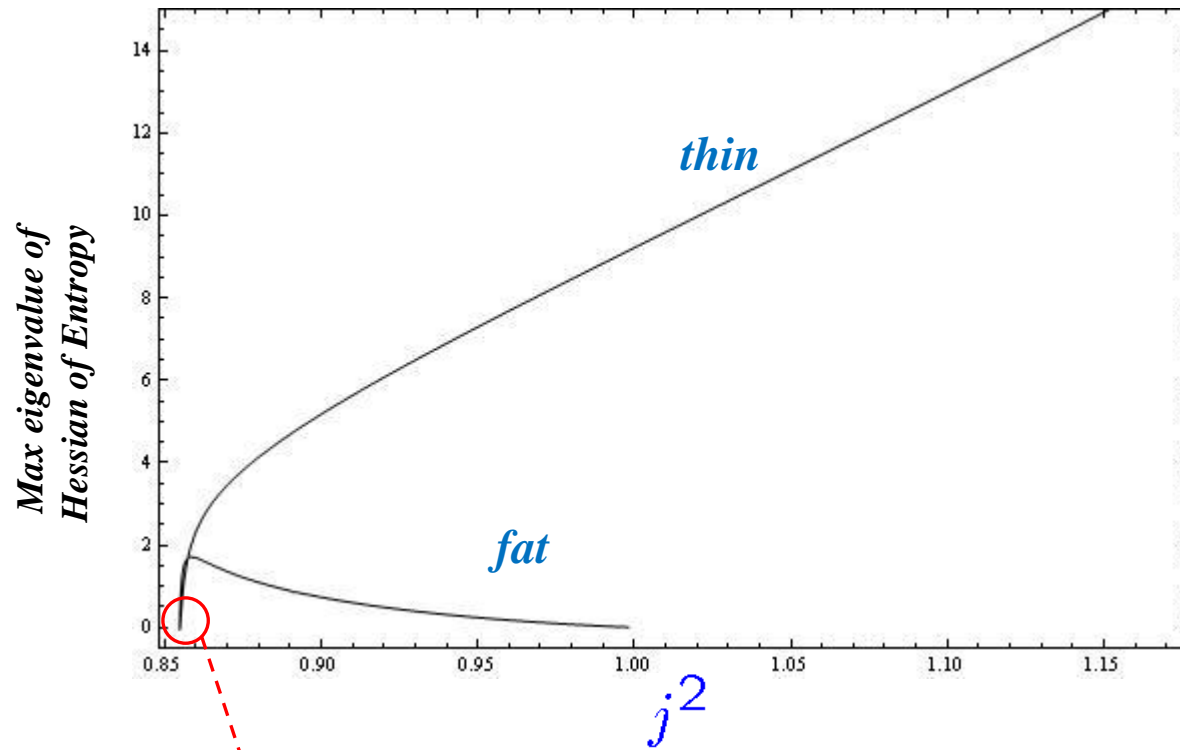
- As j increases, TBS immediately approaches BR, while TBD does not approach the BR relatively.
- As j increases, in the TBS case the central black hole is just left, in the TBD case the inner ring always runs after the outer ring.

2. Meta-stability of the thermodynamic di-ring system

- Next , as another peculiar property of the di-ring , we consider a certain kind of thermodynamic local stability following the discussion introduced by Evslin & Krishnan (JHEP 09(2008)003).
- To do this, under the condition of fixed mass and angular momentum , we shall search for local maxima of the corresponding entropy function, which is a function of appropriate moduli-parameters deduced from rod structure.
- If maximal eigenvalue of Hessian of the entropy function becomes negative, we can say that meta-stability occurs at this point.

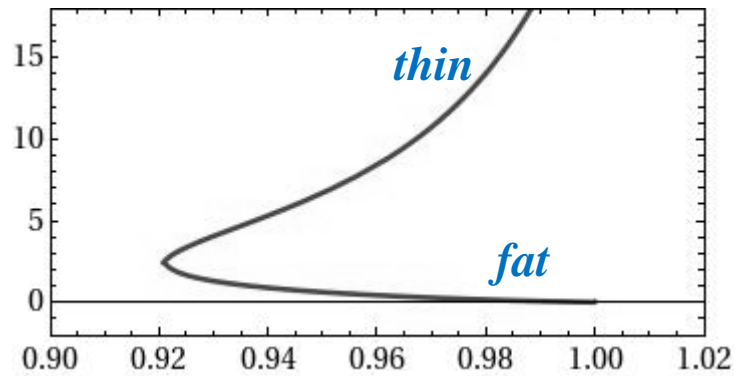
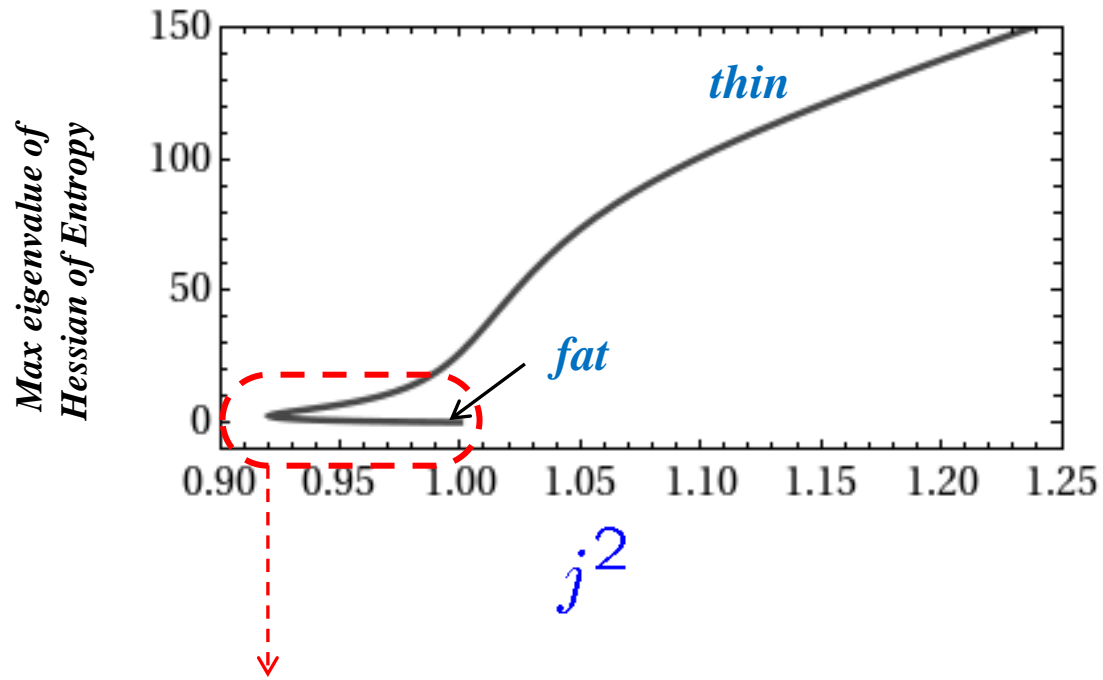


(1) The behavior of the thermodynamic black Saturn as a demonstration of our method



The same meta-stability as the result of Evslin and Krishnan is re-established!

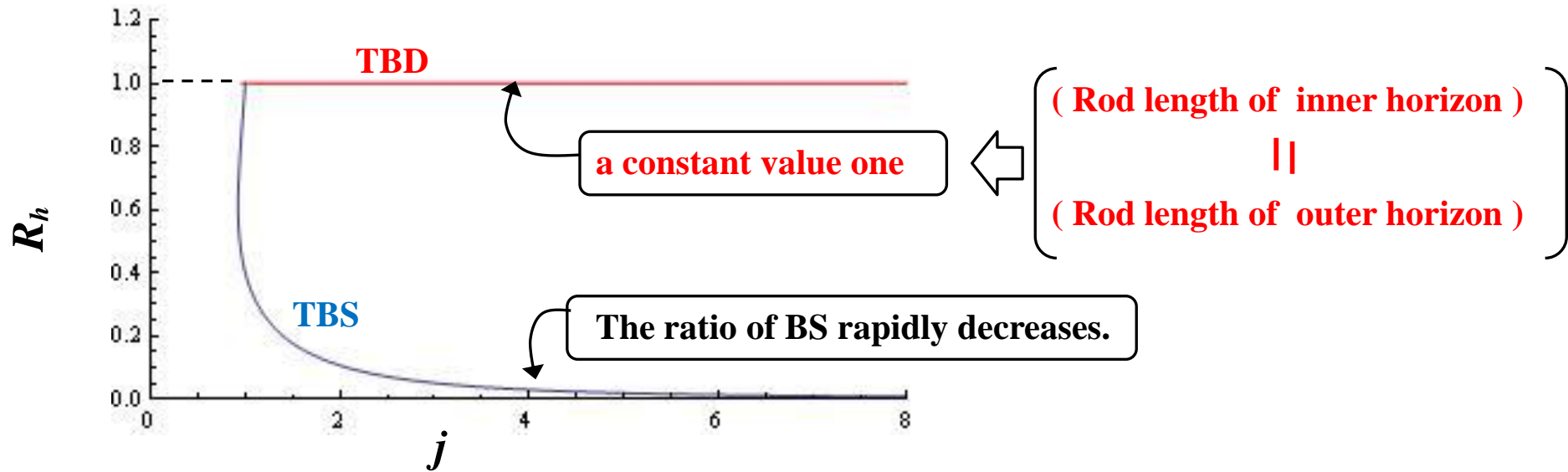
(2) The behavior of the thermodynamic black di-rings



no meta-stability appears.

3. Ratio of horizon areas of the inner BR(BH) and the outer BR

$$R_h = \frac{\text{area of inner BH or BR}}{\text{area of outer BR}}$$



(TBS)

- The influence of the inner BH quickly vanishes so that the BS behaves like a single BR.



(TBD)

- The interaction between two rings do not become so weak or their nonlinear intimate relation remains ... ?

III. Summary

- **The thermodynamic systems of regular di-rings exist.**
 - **The phase of thermodynamic di-ring has a ‘fat ring’ branch and a ‘thin ring’ branch.**
 - **In the thin branch, the behavior of the phase of black di-ring is very different from that of black Saturn.**
 - **No meta-stable state seems to be realized in the thermodynamic black di-ring, compared with the black Saturn.**
 - **Entropies (i.e. horizon areas) of the inner ring and the outer ring are always equal. Some hidden mathematical or physical structure might be expected.**
- ✓ **See [Phys. Rev. D82, 084009\(2010\): arXiv:1008.4290v2\[hep-th\]](#) for further results and discussion.**
 - ✓ **Some similar results are also obtained by Emparan and Figueras. ([JHEP11\(2010\)022: arXiv:1008.3234\[hep-th\]](#))**